

Impact of Uncertainties in Hadron Production on Air-Shower Predictions

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Introduction :

- ✚ Spectrum
- ✚ Observables

Phenomenology of Air-Showers :

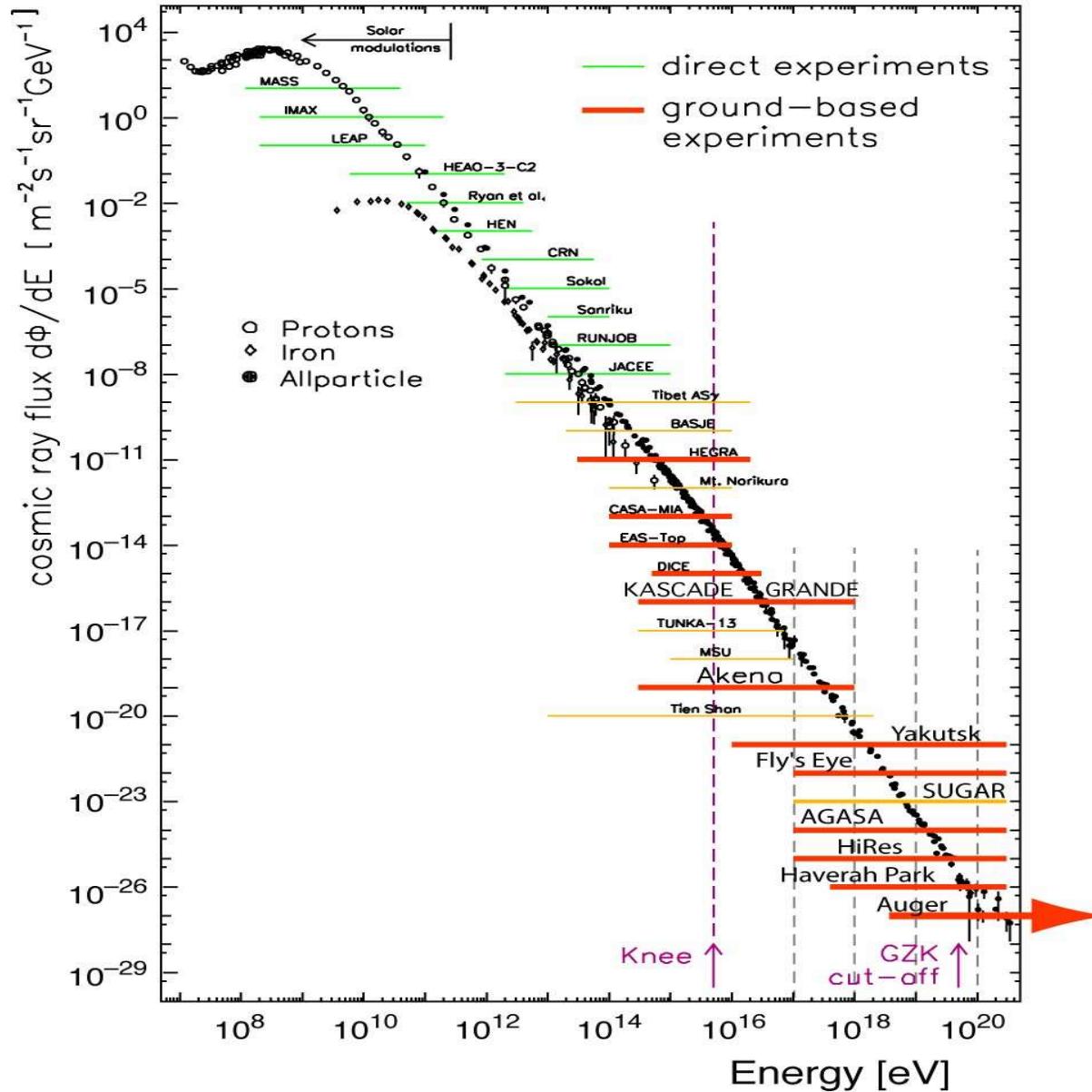
- ✚ em. shower
- ✚ muons
- ✚ from hadr. to em.

Applications :

- ✚ $\langle X_{\max} \rangle$
- ✚ LDF
- ✚ N_μ vs N_e

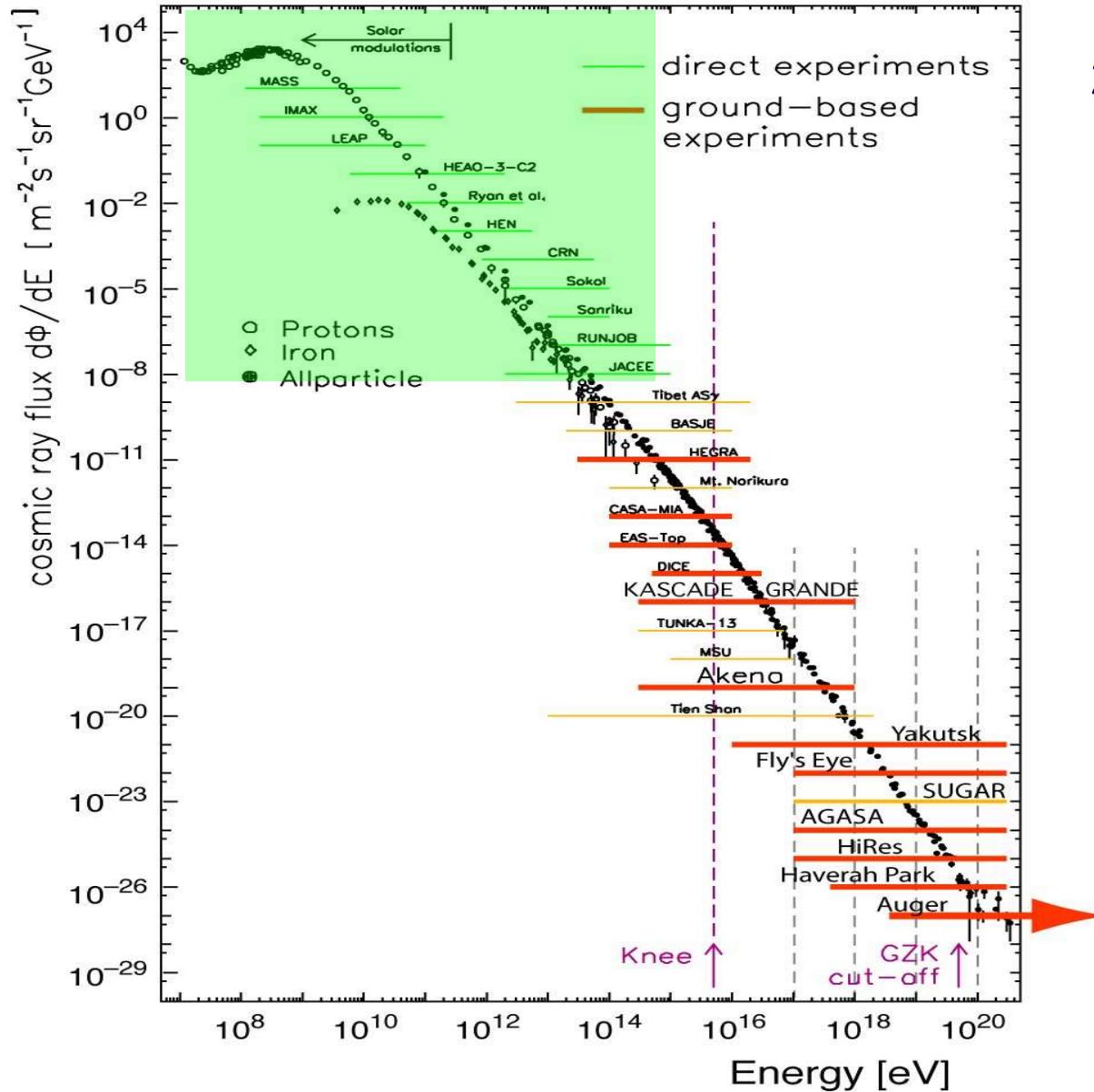
September 9th 2005

Cosmic Ray Spectrum



Cosmic ray spectrum explored by
2 different ways :

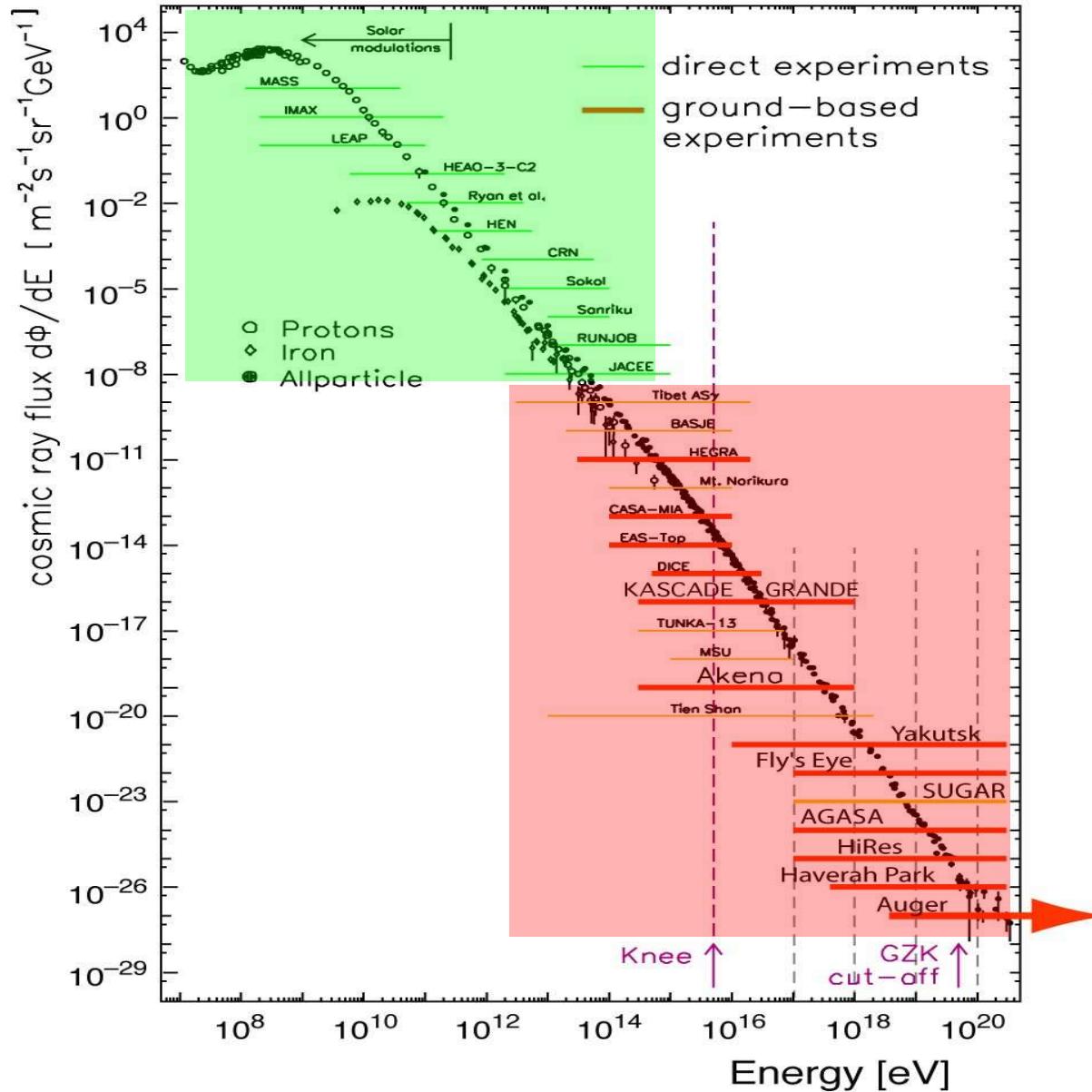
Cosmic Ray Spectrum



Cosmic ray spectrum explored by 2 different ways :

$E < 10^{14}$ eV : large flux
+ direct detection

Cosmic Ray Spectrum

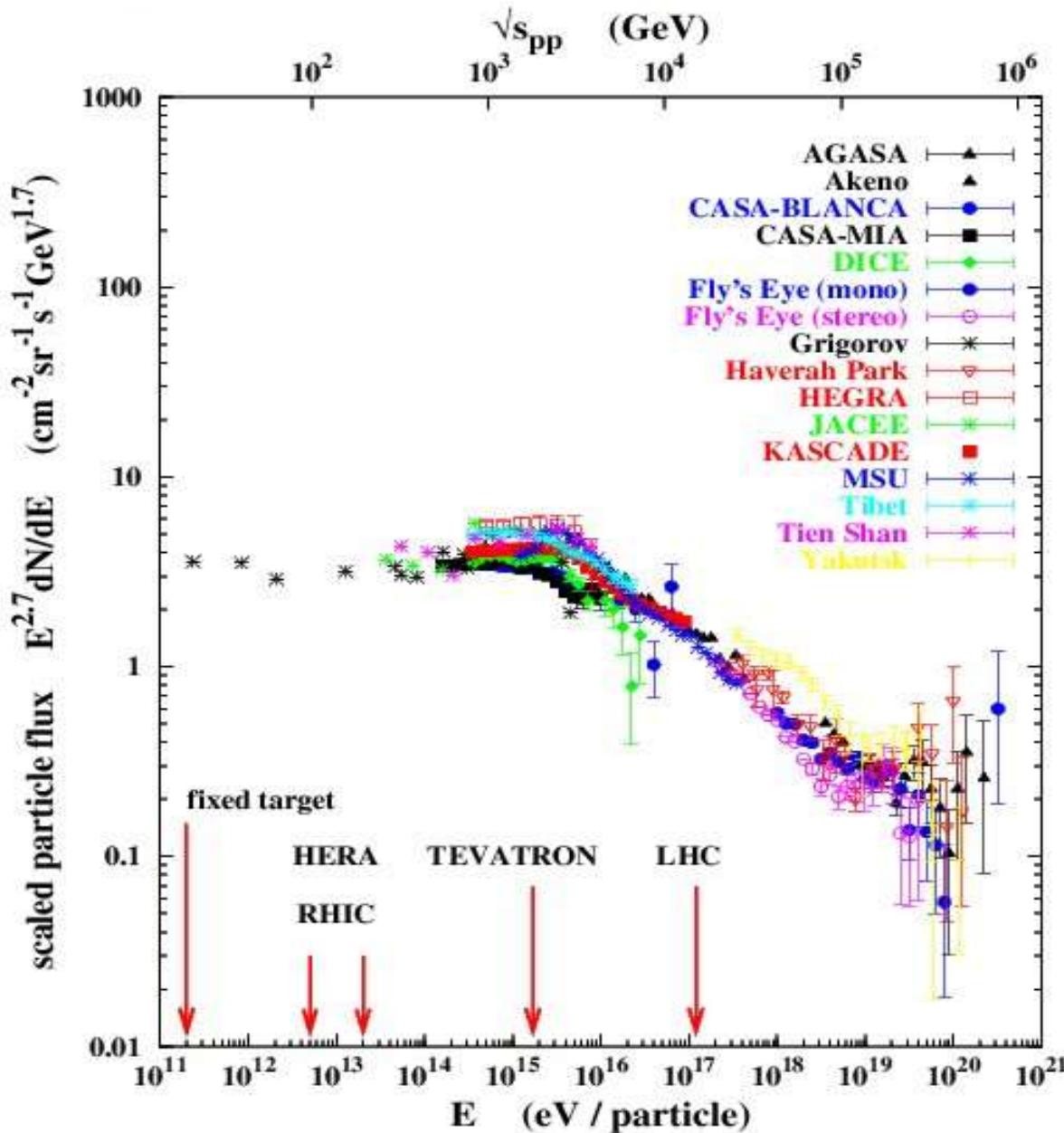


Cosmic ray spectrum explored by 2 different ways :

$E < 10^{14}$ eV : large flux
+ direct detection

$E > 10^{14}$ eV : low flux
+ indirect detection : air-shower reached ground

Cosmic Ray Spectrum



Cosmic ray spectrum explored by 2 different ways :

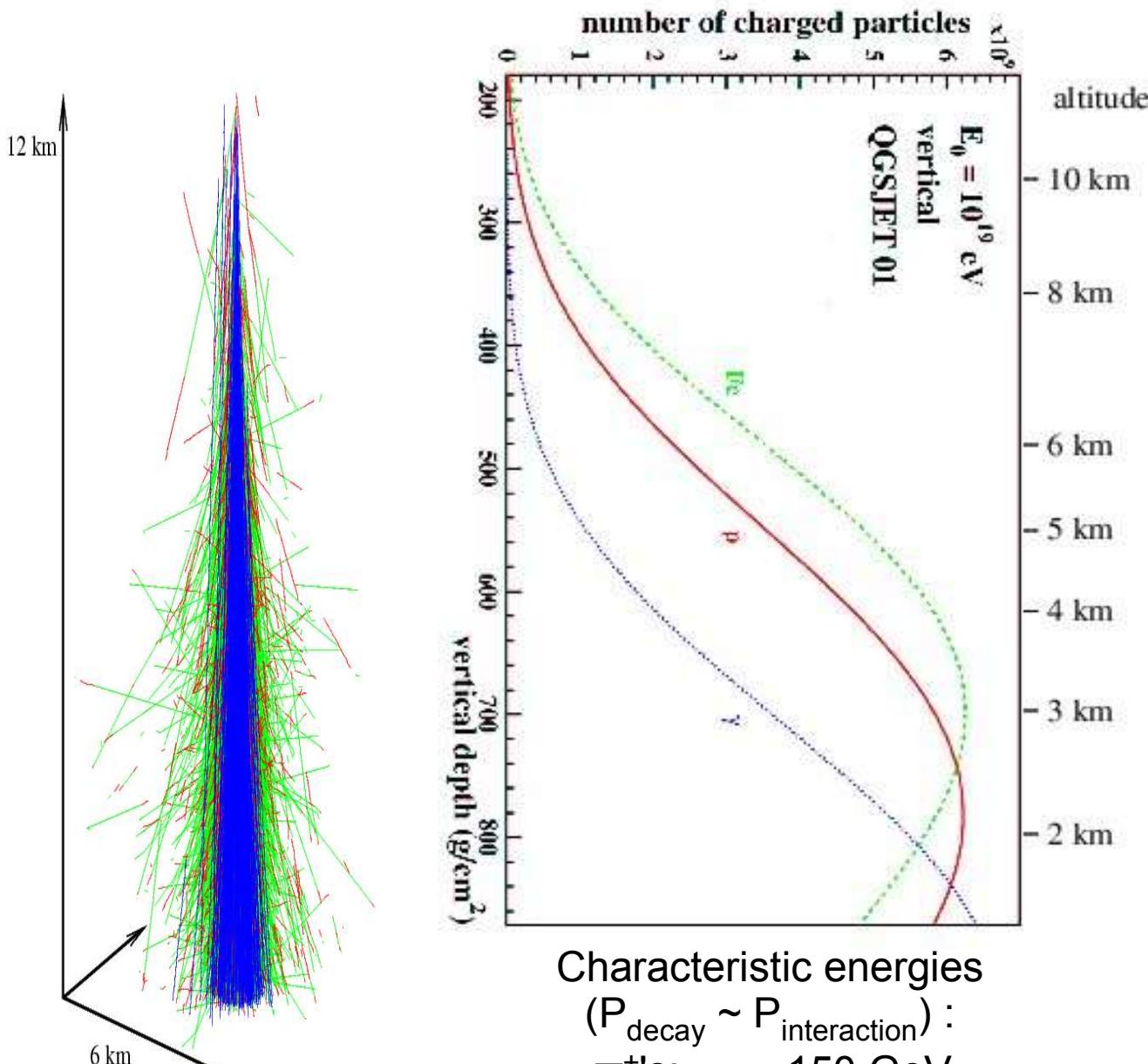
$E < 10^{14}$ eV : large flux
+ direct detection

$E > 10^{14}$ eV : low flux
+ indirect detection : air-shower reached ground

Air-shower simulations needed to interpret high energy data.

Energy above any collider experiment !

Observables (1)



Longitudinal development
of air-showers :

number of particles (or energy
deposited by charged part.)
projected on the axis

Main characteristics :

Characteristic energies

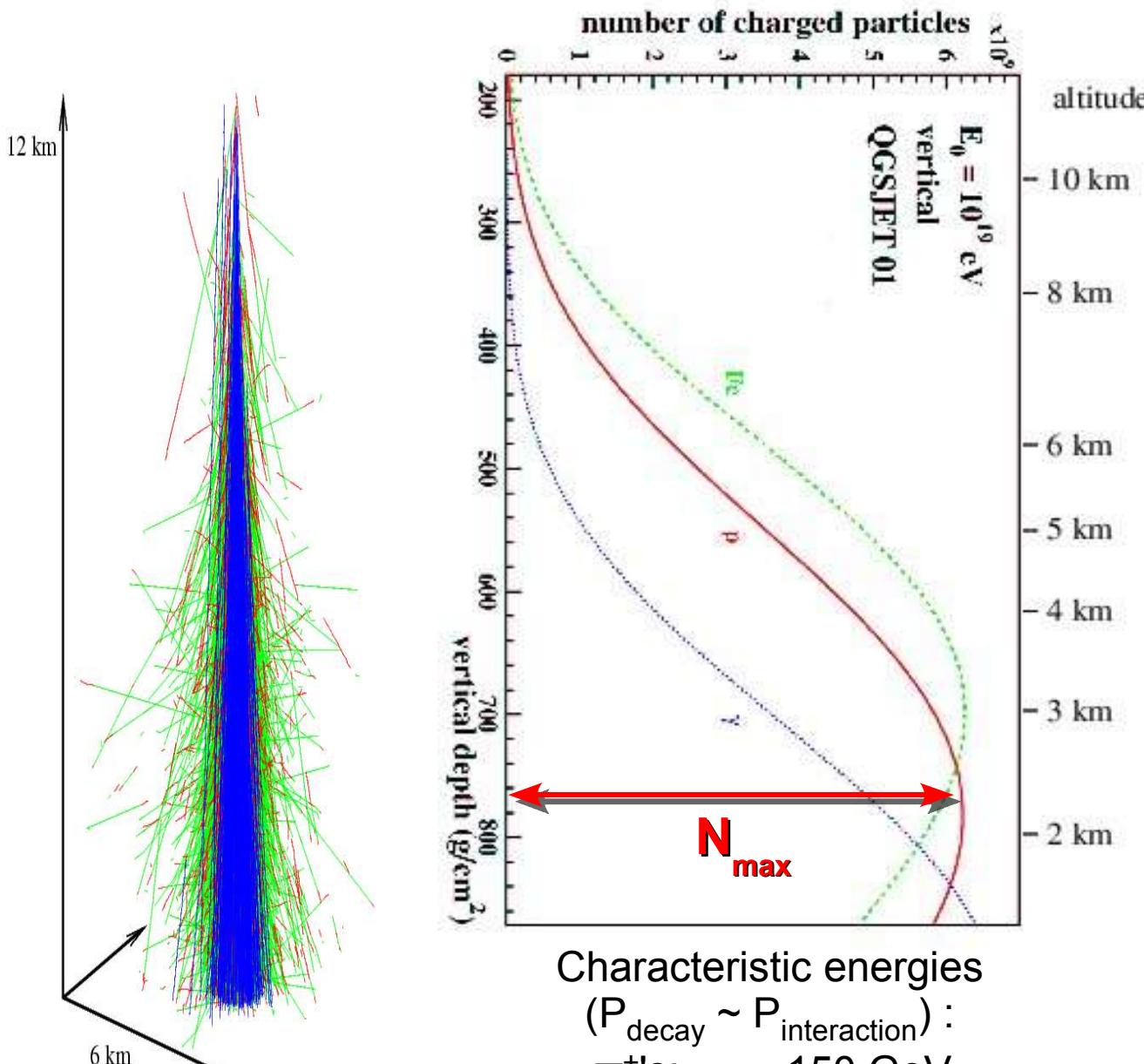
(P_{decay} ~ P_{interaction}) :

π^\pm 's: ~150 GeV

kaons: ~600 GeV

π^0 's: ~10⁹ GeV

Observables (1)



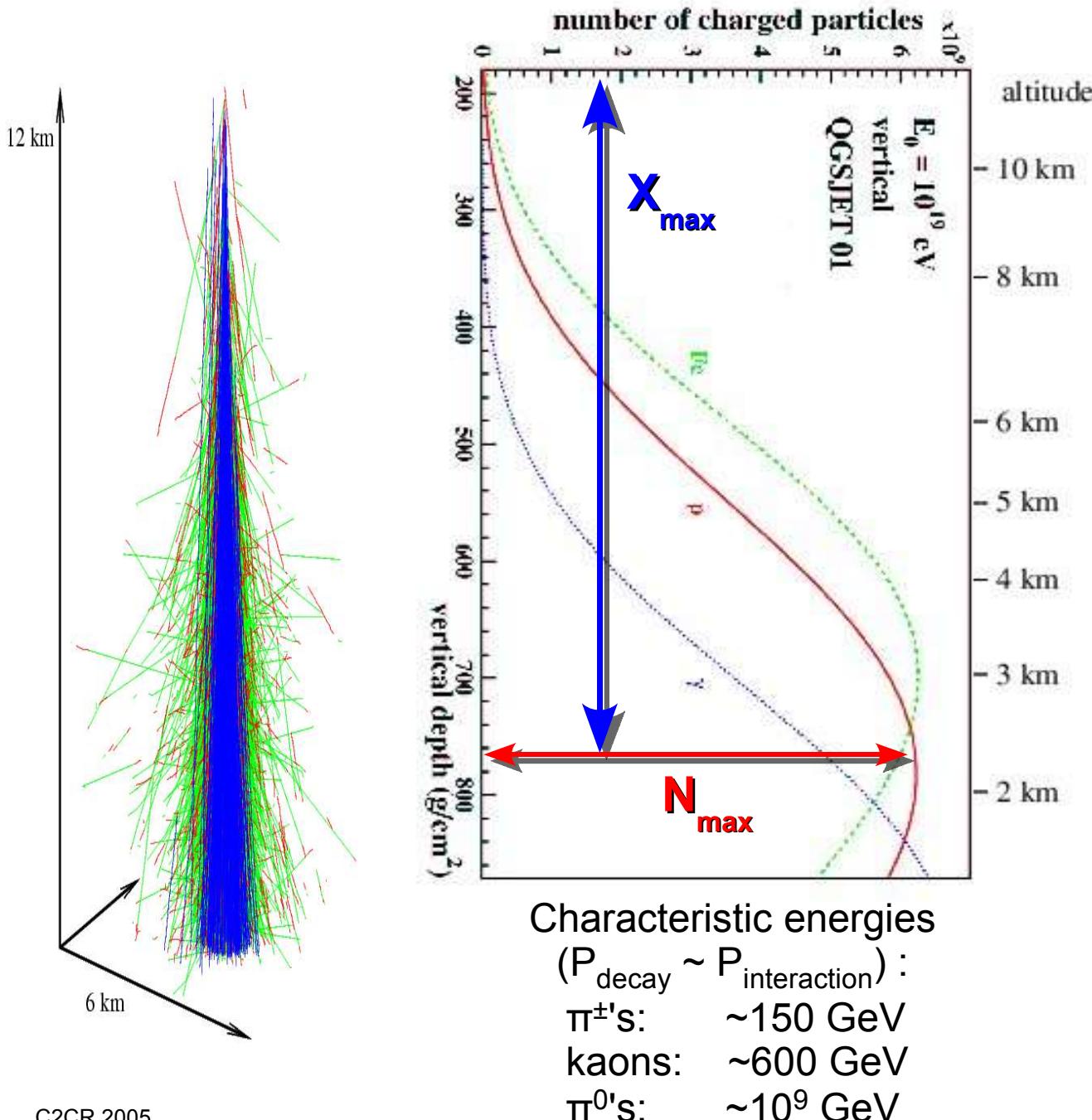
Longitudinal development of air-showers :

number of particles (or energy deposited by charged part.) projected on the axis

Main characteristics :

- Shower maximum N_{\max}
Mass independent,
energy dependent,
small fluctuations

Observables (1)



Longitudinal development of air-showers :

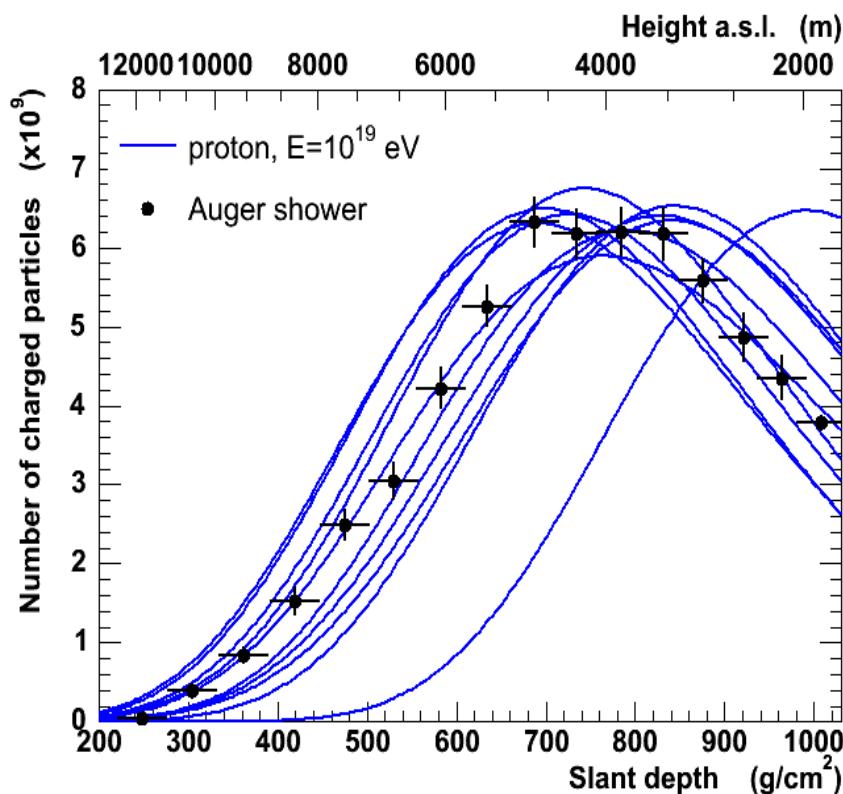
number of particles (or energy deposited by charged part.) projected on the axis

Main characteristics :

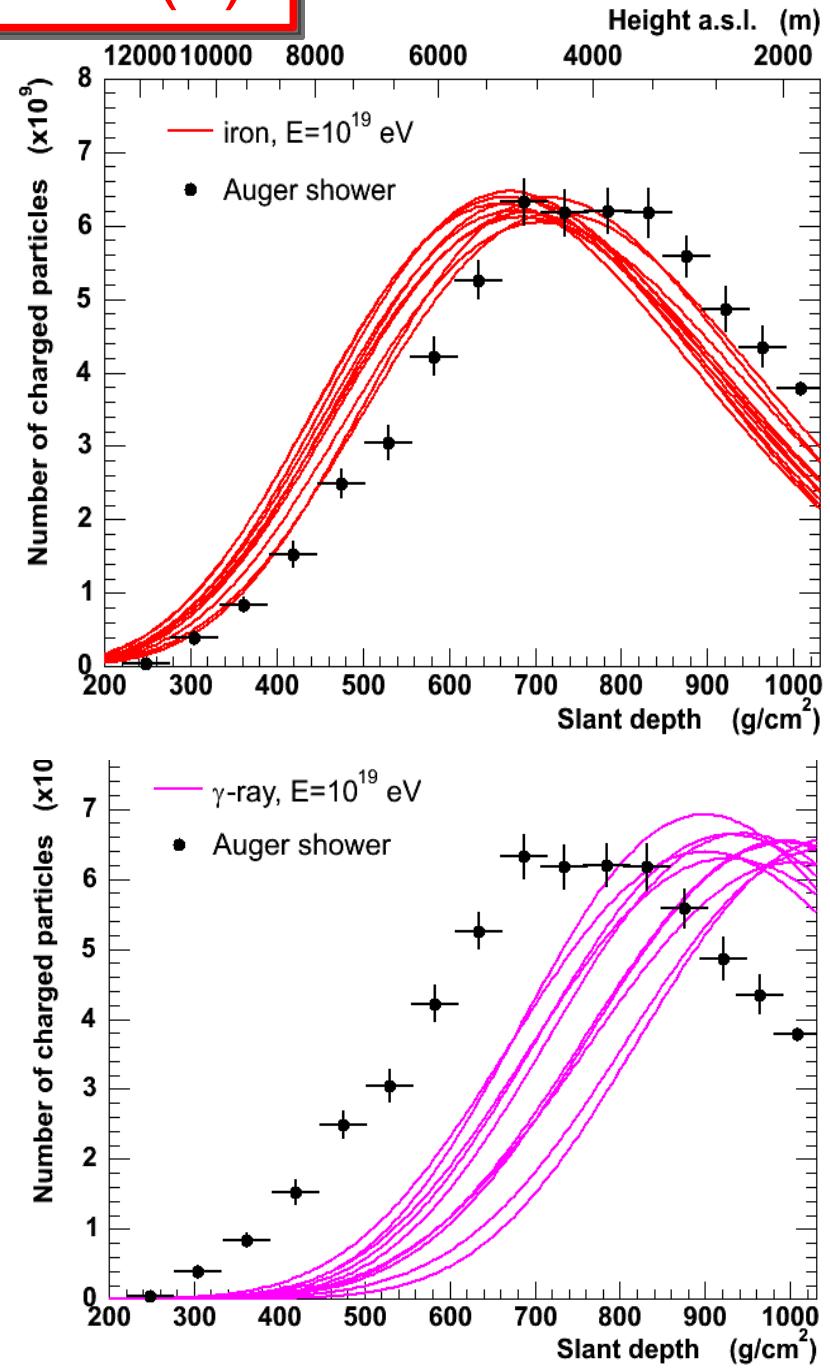
- ➡ Shower maximum N_{\max}
Mass independent,
energy dependent,
small fluctuations
- ➡ Depth of shower maximum X_{\max}
Mass and energy dependent,
large fluctuations

Observables (2)

Detailed MC simulation :
10 showers zenith angle 35° ,
QGSJET

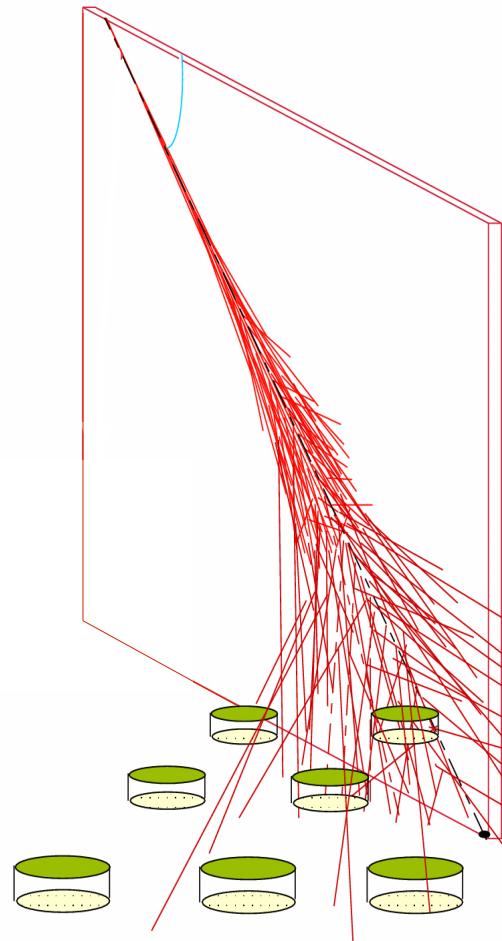


Shower-by-shower fluctuations
are very important and mass
dependent.

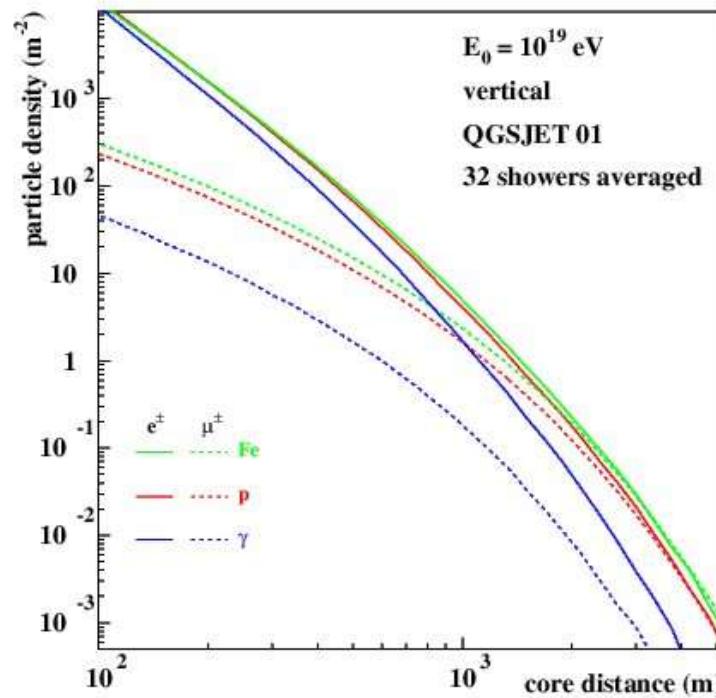


Observables (3)

Particles reaching ground :
most of them at the impact point, but some lateral spread



- ✚ Lateral Distribution Function (LDF) used to reconstruct the total number of part.

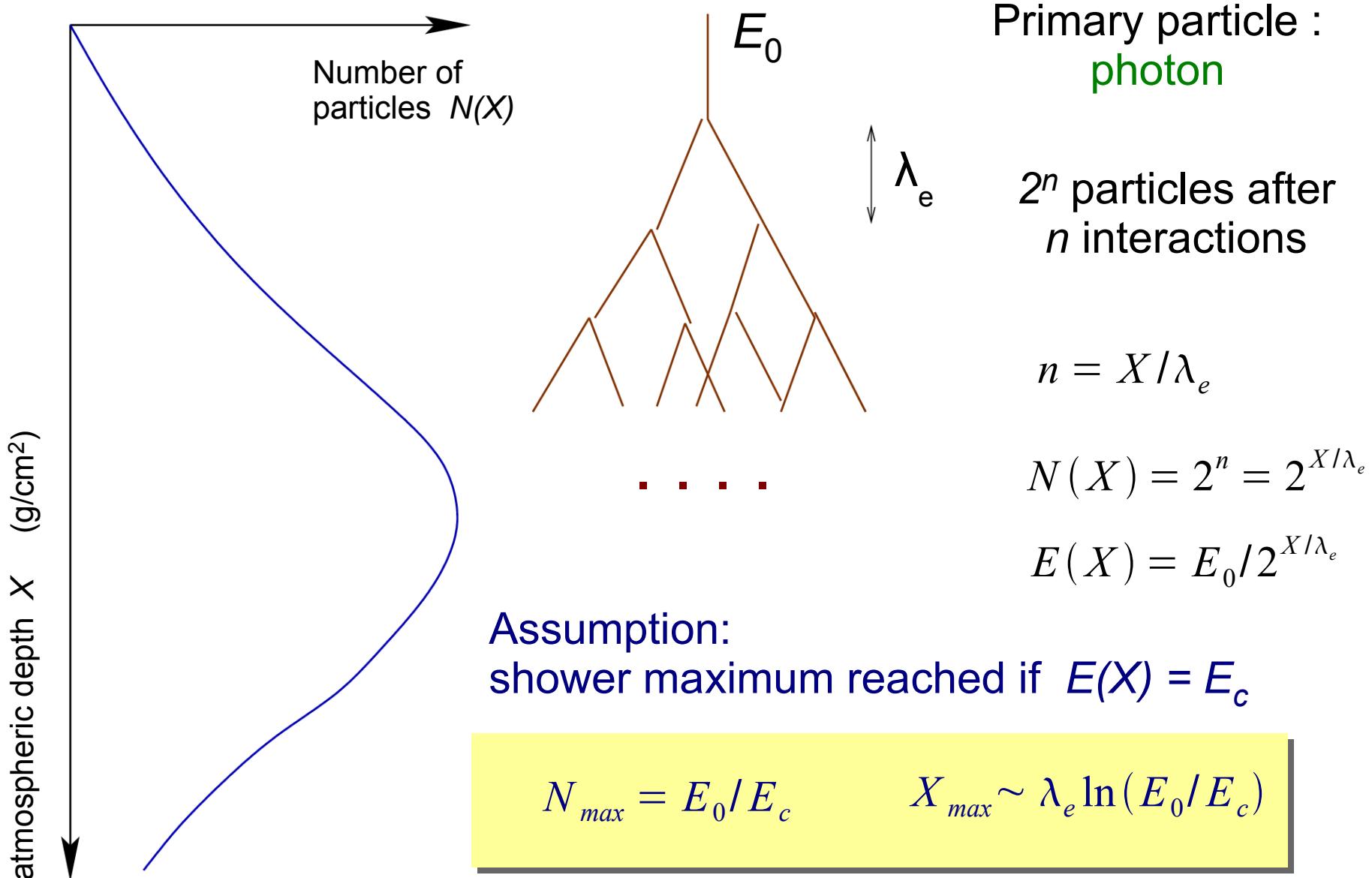


Energy from LDF

Different mass dependence of
muon's LDF and electron's
LDF

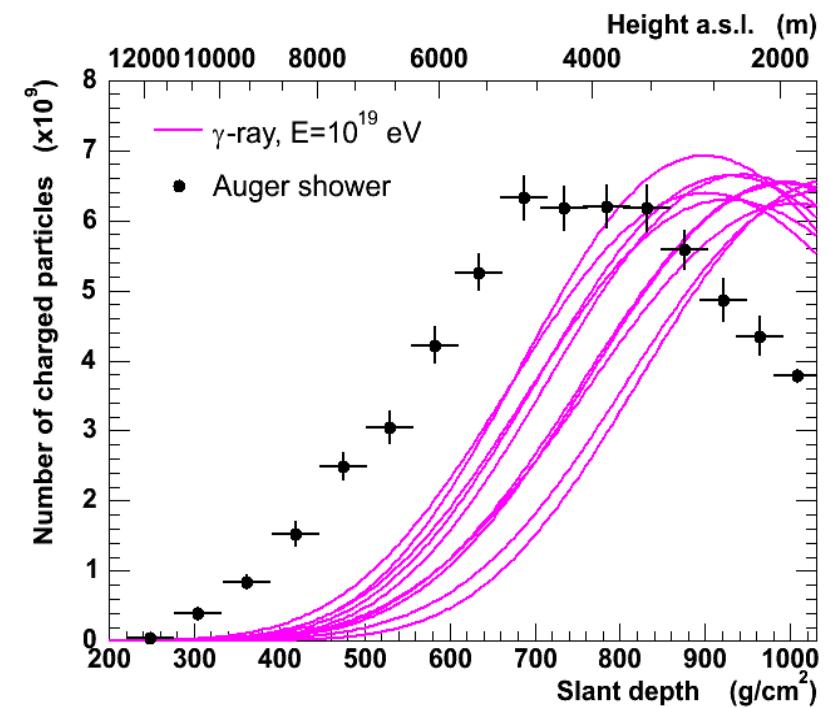
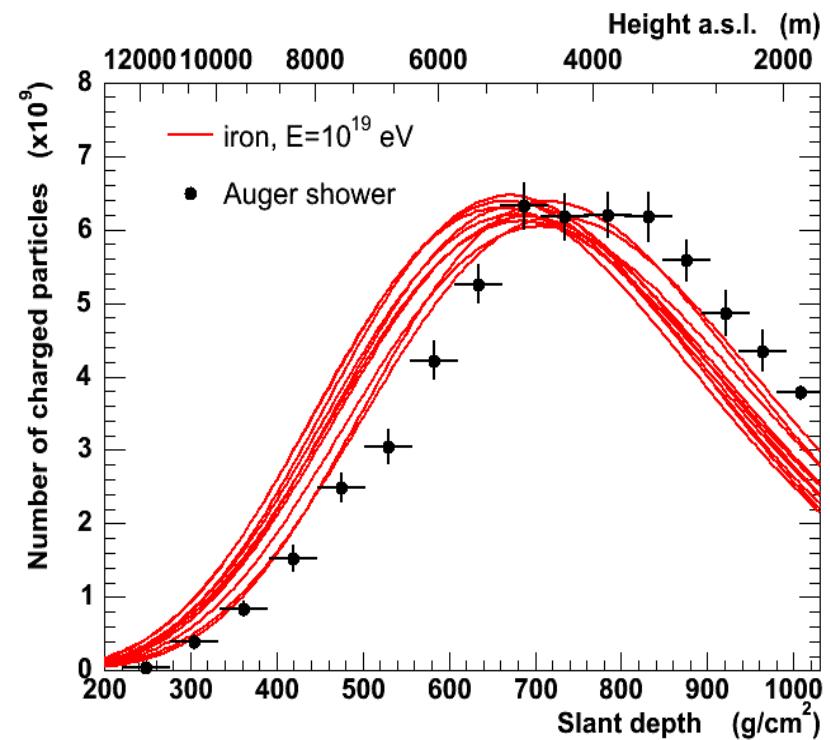
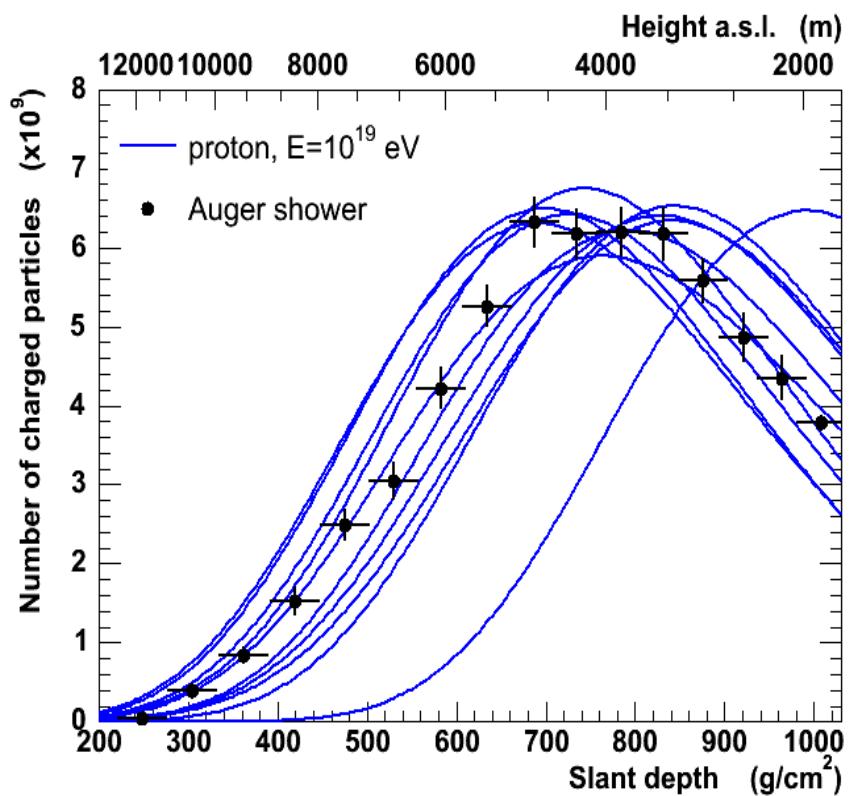
Mass from correlation
 N_μ vs N_e

Heitler's Model of EM Showers

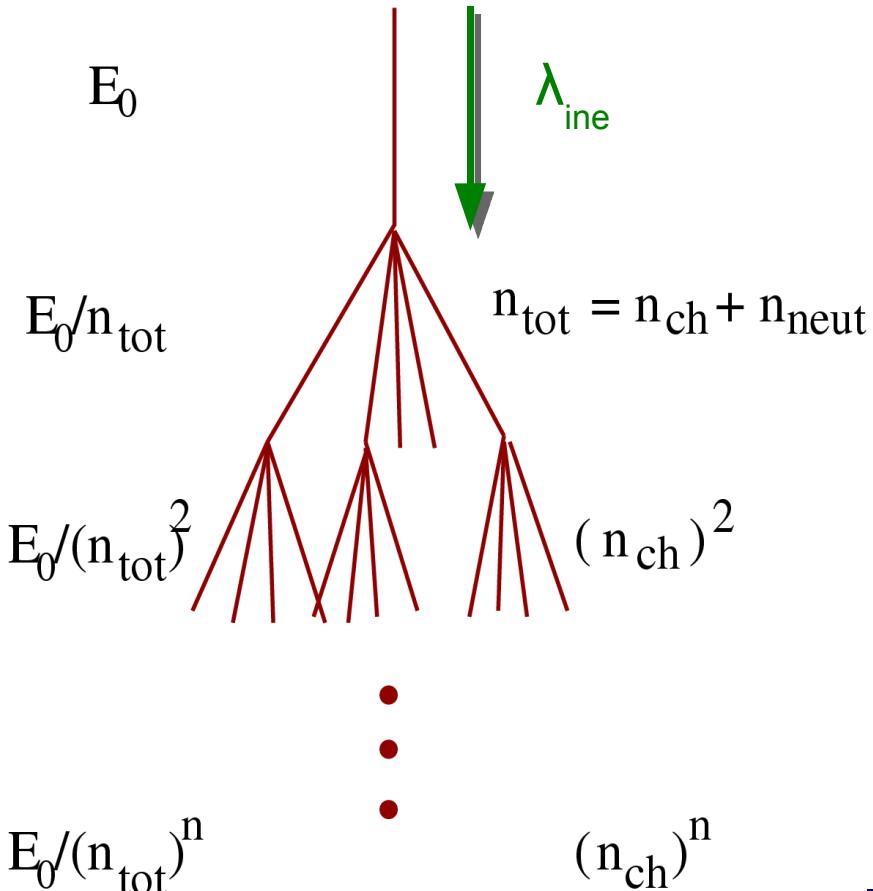


N_{\max}

Detailed MC simulation :
10 showers zenith angle 35° ,
QGSJET



Muon Component



$$X_{max} \sim \lambda_e \ln\left(E_0/(2n_{tot})\right) + \lambda_{ine}$$

Primary particle:
proton

$$n_{ch} \sim \pi^{+/-}$$

$$n_{neut} \sim \pi^0 :$$

decay immediately $\rightarrow 2\gamma$

Cascade ends with
decay at energy E_{dec}

$$E(X) = E_0/(n_{tot})^n = E_{dec}$$

at this n : $N_\mu = (n_{ch})^n$

$$N_\mu = \left(\frac{E_0}{E_{dec}}\right)^\alpha, \quad \alpha = \frac{\ln n_{ch}}{\ln n_{tot}} \approx 0.82 \dots 0.95$$

Superposition Model

Proton shower characteristics:

$$N_{max} = E_0/E_c$$

$$X_{max} = \lambda_e \ln(E_0/n_{tot}) + \lambda_{ine}$$

$$N_\mu = \left(\frac{E_0}{E_{dec}} \right)^\alpha ; \alpha = \frac{\ln n_{ch}}{\ln n_{tot}}$$

Hadronic interaction model

→ interaction cross section

→ mean multiplicity of secondary particles

→ ratio of neutral to charged pion / kaon multiplicities

Atmosphere as target and calorimeter

→ critical energy

→ density : typical pion / kaon decay energy

Assumption:

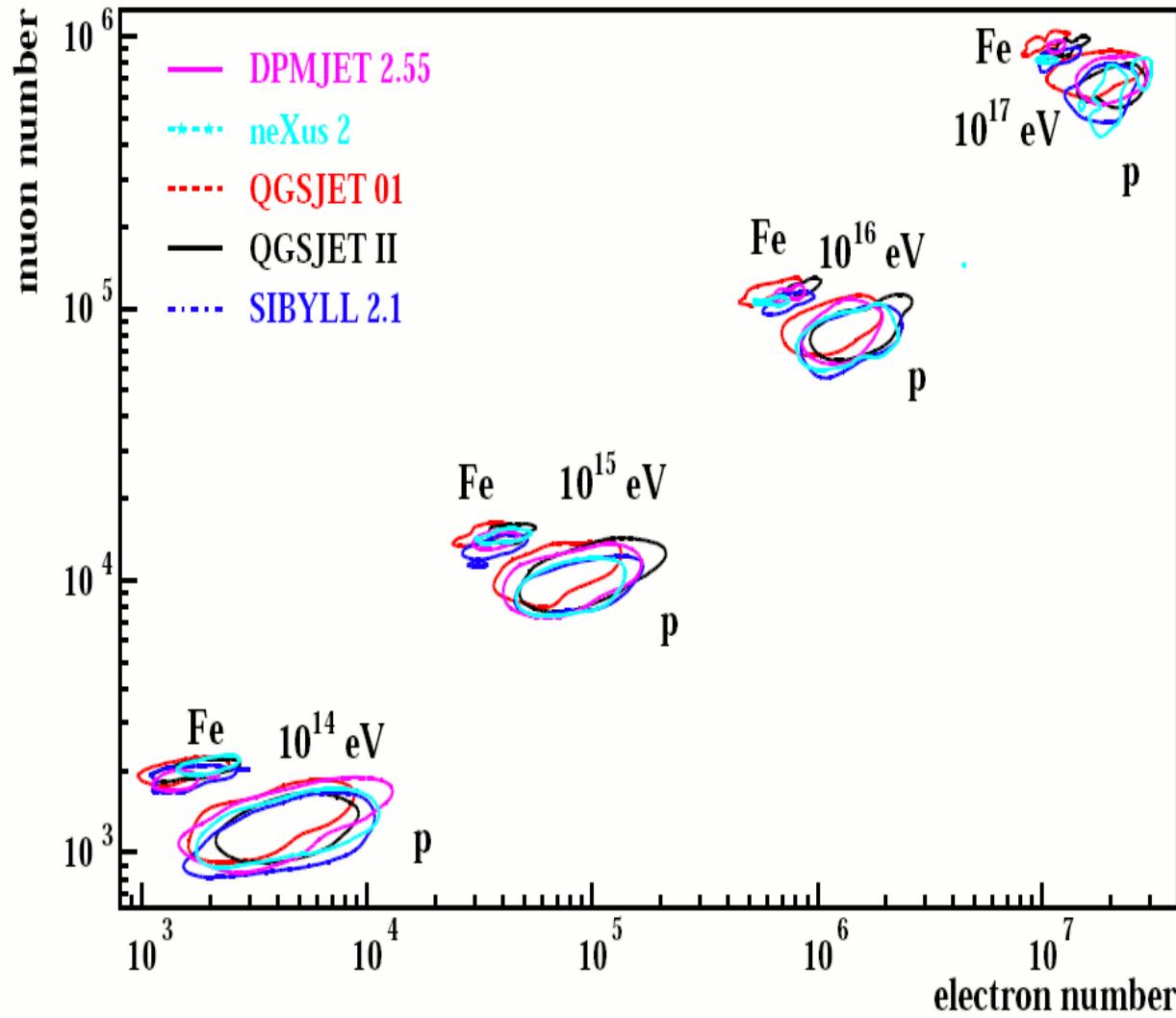
nucleus of mass A and energy E_0 acts like A independent nucleons with energy $E_n = E_0/A$

$$N_{max}^A = A E_n/E_c = E_0/E_c$$

$$X_{max}^A(E_0) = X_{max}^p(E_0/A)$$

$$N_\mu^A = A \left(\frac{E_0/A}{E_{dec}} \right)^\alpha = A^{1-\alpha} N_\mu$$

N_μ and N_e at ground



- N_e ~ increases linearly with energy
- Different behaviour with A for N_e and N_μ :
 - N_e decreases with A ($N_e \neq N_{max}$) (X_{max} decrease → older shower)
 - N_μ increases with A ($N_\mu \sim \text{const. after maximum}$)

good mass estimator ...

... within a given model !

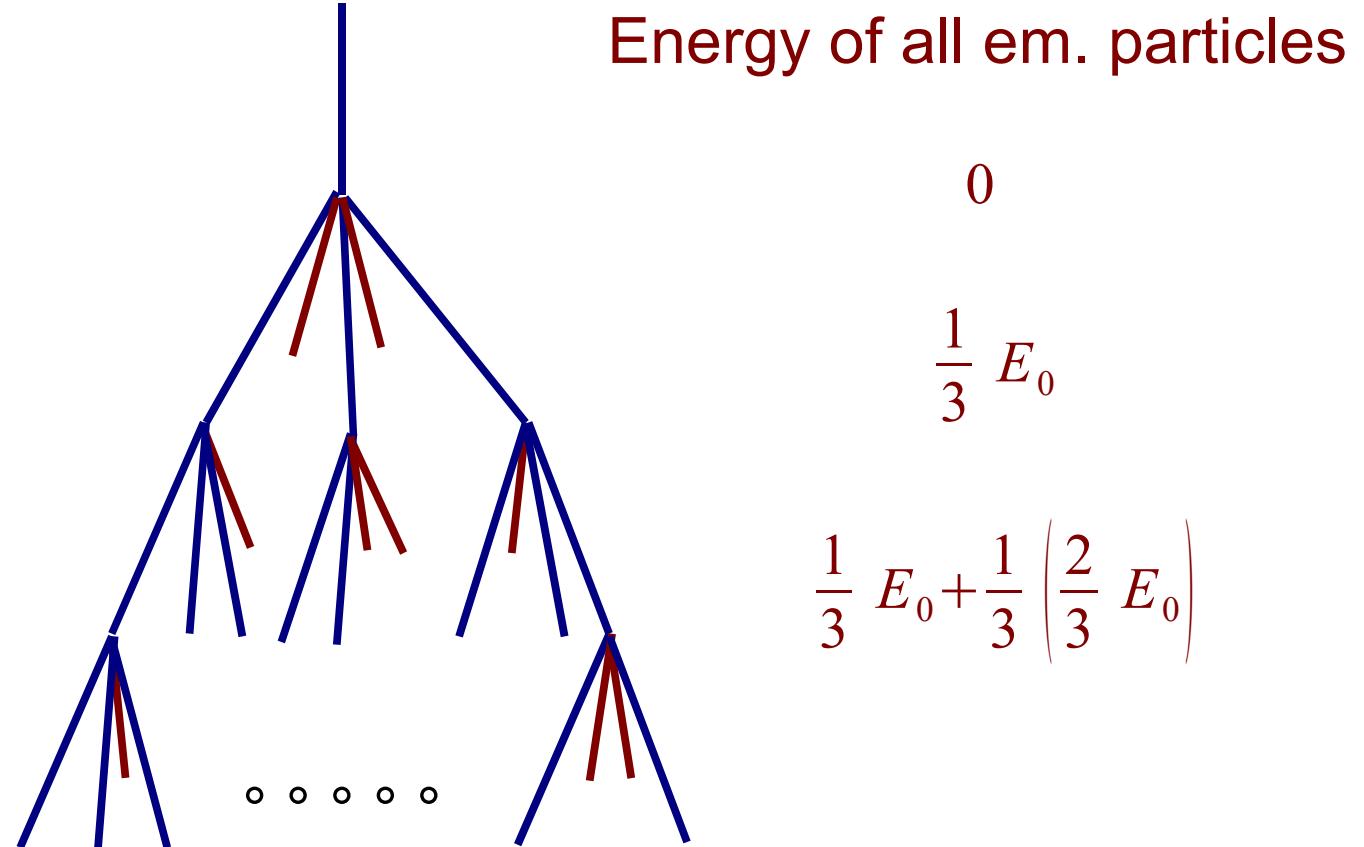
Energy Transfer

Energy of all hadrons

$$\begin{aligned} E_0 \\ \frac{2}{3} E_0 \\ \frac{2}{3} \left(\frac{2}{3} E_0 \right) \end{aligned}$$

After n generations

$$E_{\text{had}} = \left(\frac{2}{3} \right)^n E_0$$



Energy in em. $\sim 90\%$

$(n=5, E_{\text{had}} \sim 12\%)$
 $(n=6, E_{\text{had}} \sim 8\%)$

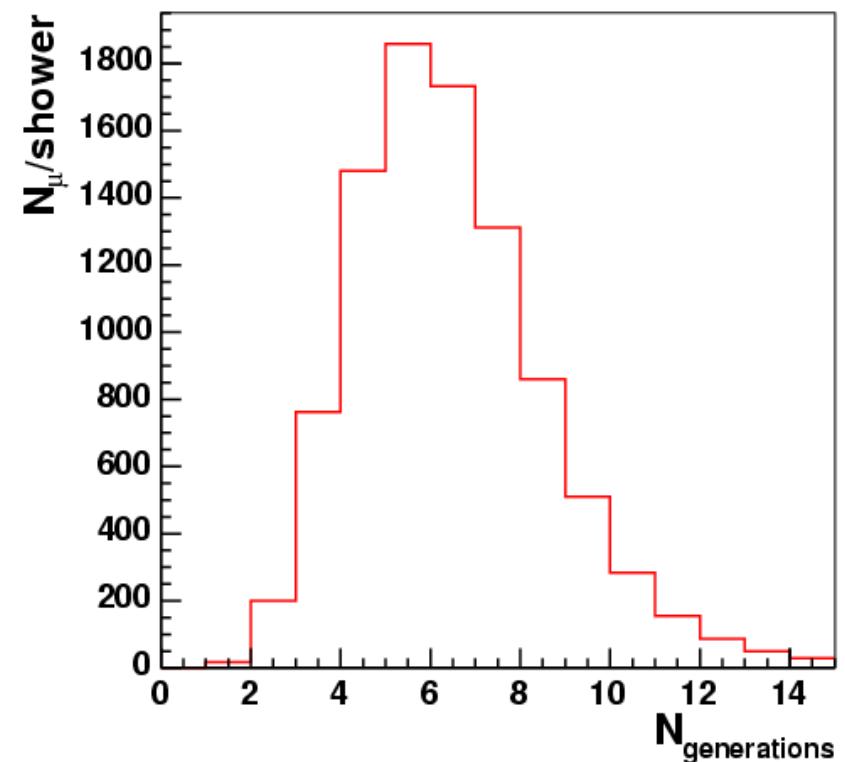
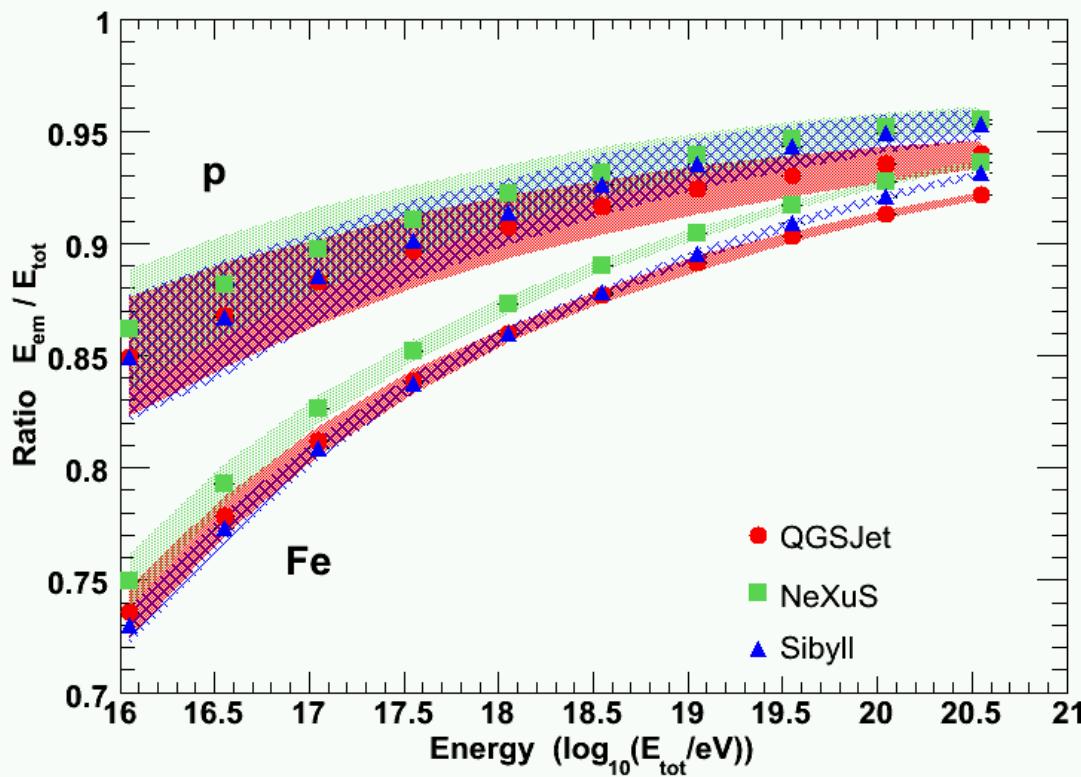
$$E_{\text{em}} = \left[1 - \left(\frac{2}{3} \right)^n \right] E_0$$

Energy Transfer

From Heitler's model we know :

- ⊕ ~ 90 % energy in the em. particles
- ⊕ ~ 5 - 6 hadronic generations

Confirmed by MC simulation
and model independent !



Summary on Heitler's Model

Air-shower characteristics controlled by 2 atmospheric parameters and 3 hadronic model dependent parameters :

- ✚ hadronic cross section : X_{\max}
- ✚ mean total multiplicity : X_{\max} and number of generations
- ✚ ratio charged/total number of particles : muon number

Hadronic interaction MC models disagree on this parameters !

Hadronic Interaction Models

Realistic Monte-Carlo simulation needs hadronic interaction models

High energy models :

DPMJET II.5 and III (Ranft / Roesler, Engel & Ranft)

NEXUS 2 and 3 (Drescher, Hladik, Ostapchenko, TP & Werner)

QGSJET 01 and II (Kalmykov & Ostapchenko / Ostapchenko)

SIBYLL 2.1 (Engel / Engel, Fletcher, Gaisser, Lipari & Stanev)

Low/intermediate energy models:

GHEISHA (Fesefeldt)

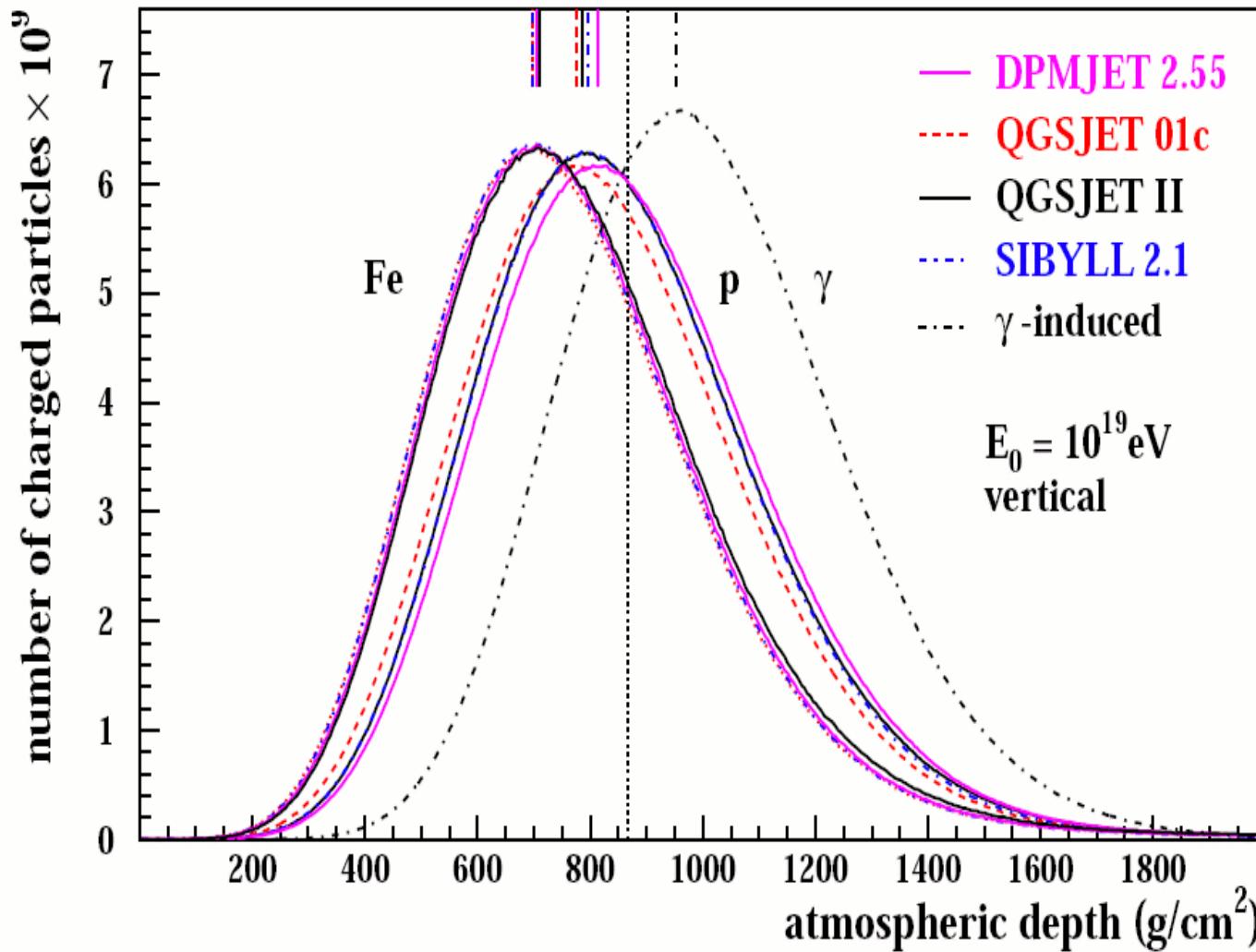
Hillas' splitting algorithm (Hillas)

FLUKA (Fasso, Ferrari, Ranft, Roesler & Sala)

UrQMD (Bleicher et al.)

- ➡ Gribov-Regge type models
- ➡ QCD-inspired minijet production
- ➡ Parametrization of data
(see S. Ostapchenko's talk)

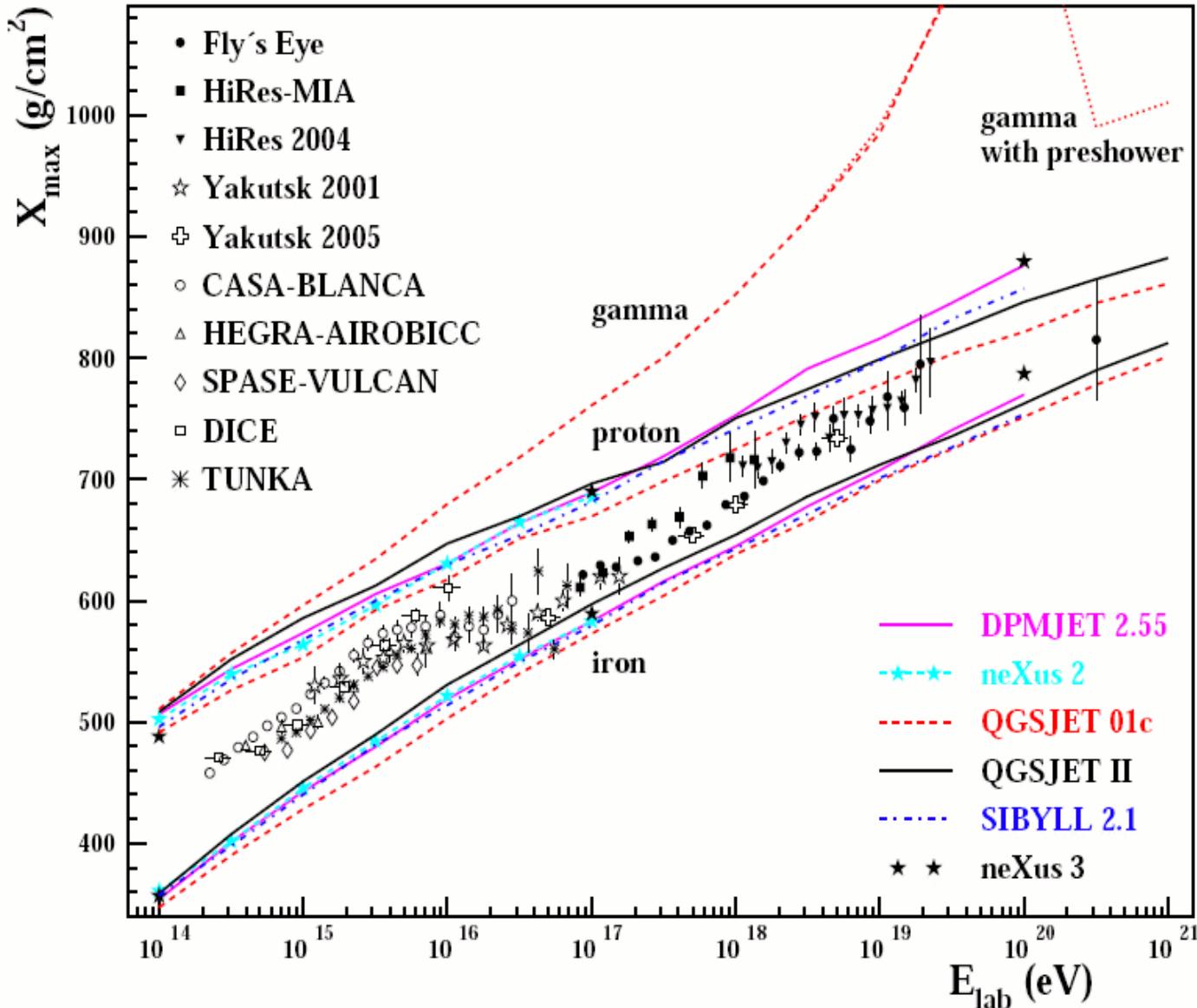
Longitudinal Development



Shape is model-independent (em. processes)

Differences between models appears in X_{\max}

Mean Depth of Shower Maximum



Prediction are strongly model dependent at very high energy !

QGSJET 01 : highest n_{tot} and lowest $\langle X_{\text{max}} \rangle$

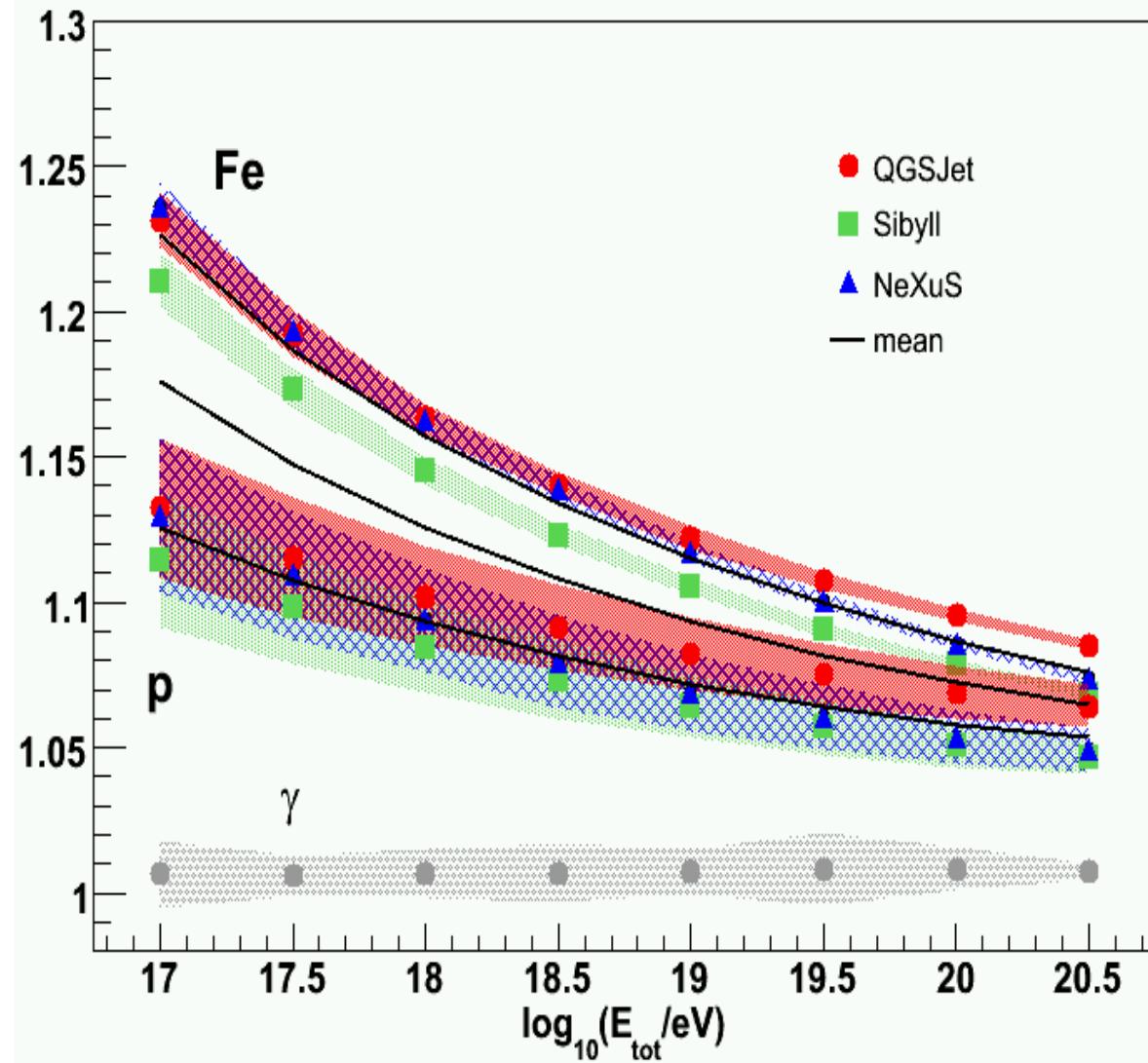
DPMJET : lowest n_{tot} and highest $\langle X_{\text{max}} \rangle$

$n_{\text{tot}} = \text{multiplicity} \otimes \text{elasticity}$

Mass interpretation
very difficult ...

Energy Deposit

$E_{\text{tot}}/E_{\text{depo}}$ Correction factor

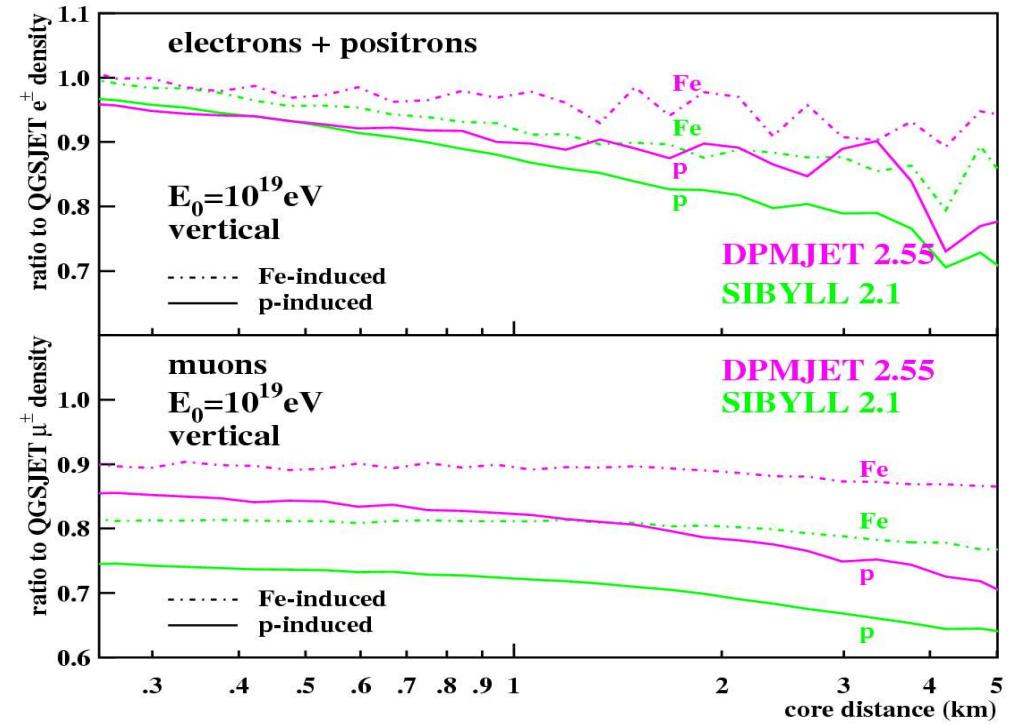
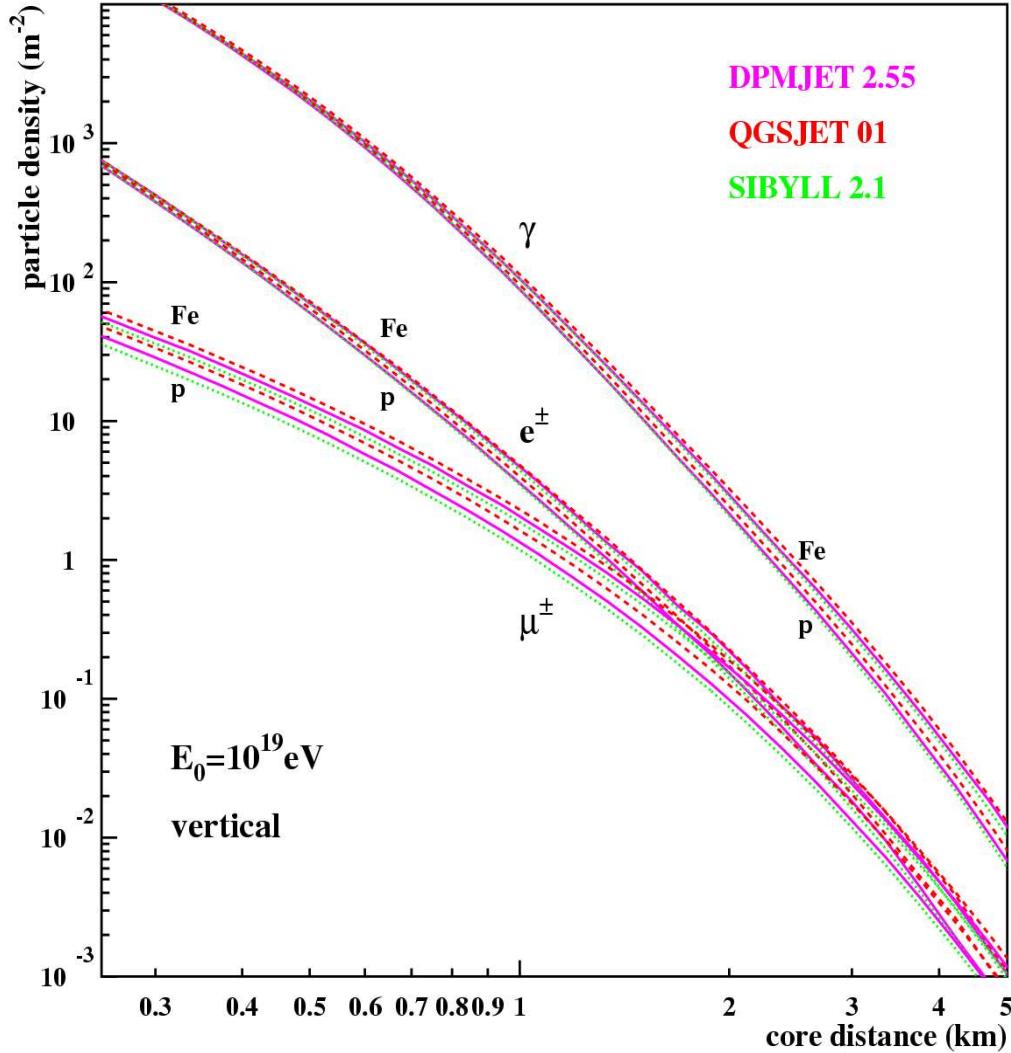


Energy deposit related to fluorescence light emitted :
energy estimator

model uncertainties ($\sim 2\%$)
lower than uncertainty due
to mass ($\sim 6\%$)

Fluorescence light based
energy measurement are
almost hadronic model
independent !

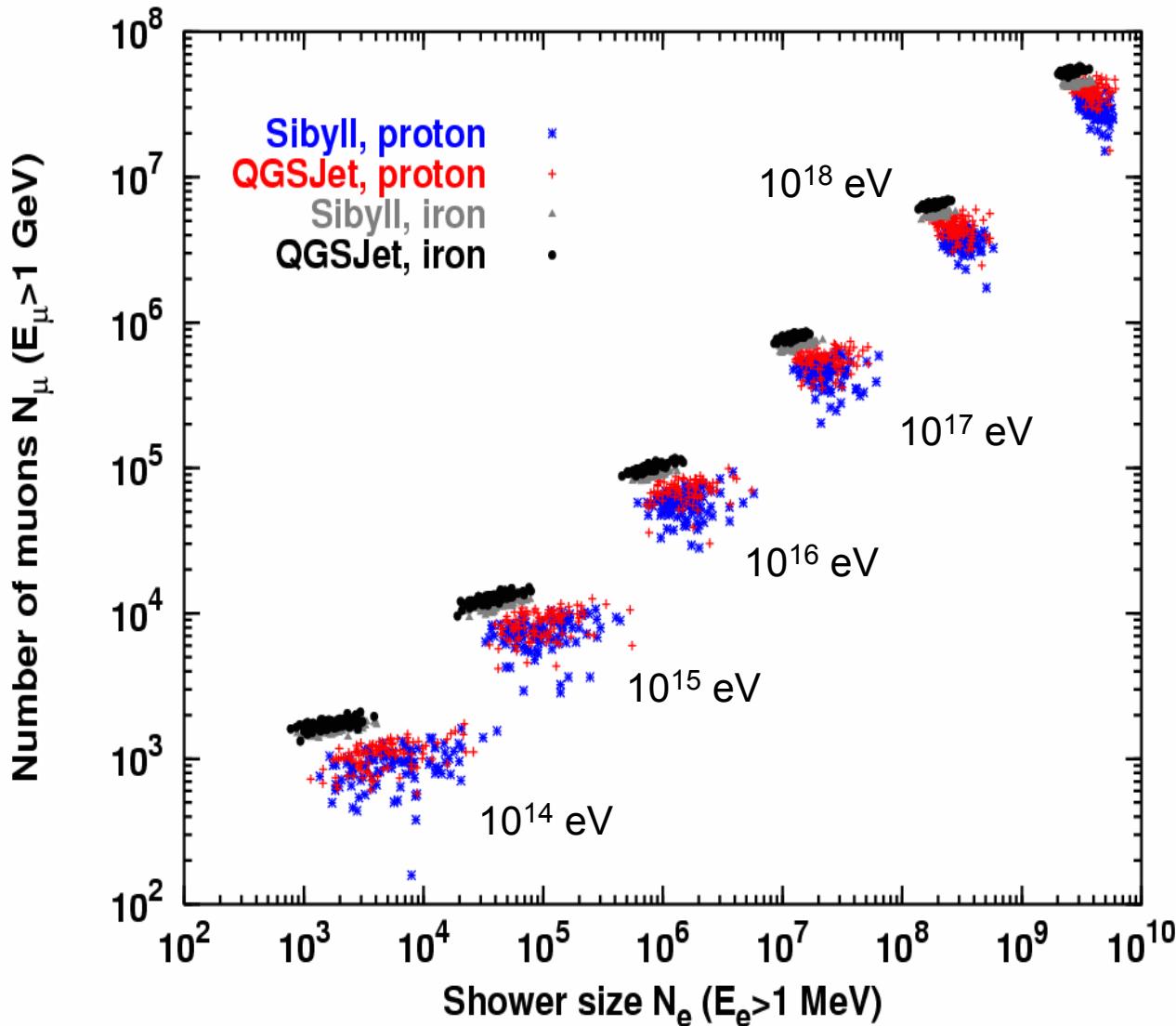
Lateral Distribution Function



LDF used for energy determination
but up to 35 % difference at large
distances !

See also Ch. Meurer's talk (dependence on low-energy models)

N_μ vs N_e



From theory, N_e and N_μ do not behave the same : good mass estimator used in many experiment (KASCADE, EAS-TOP, AKENO ...)

but strong model dependence !

Mass analysis based on N_e and N_μ are highly model dependent

(see H. Ulrich (KASCADE) talk)

Conclusion

Air-shower characteristics controlled by features of hadronic multiparticle production :

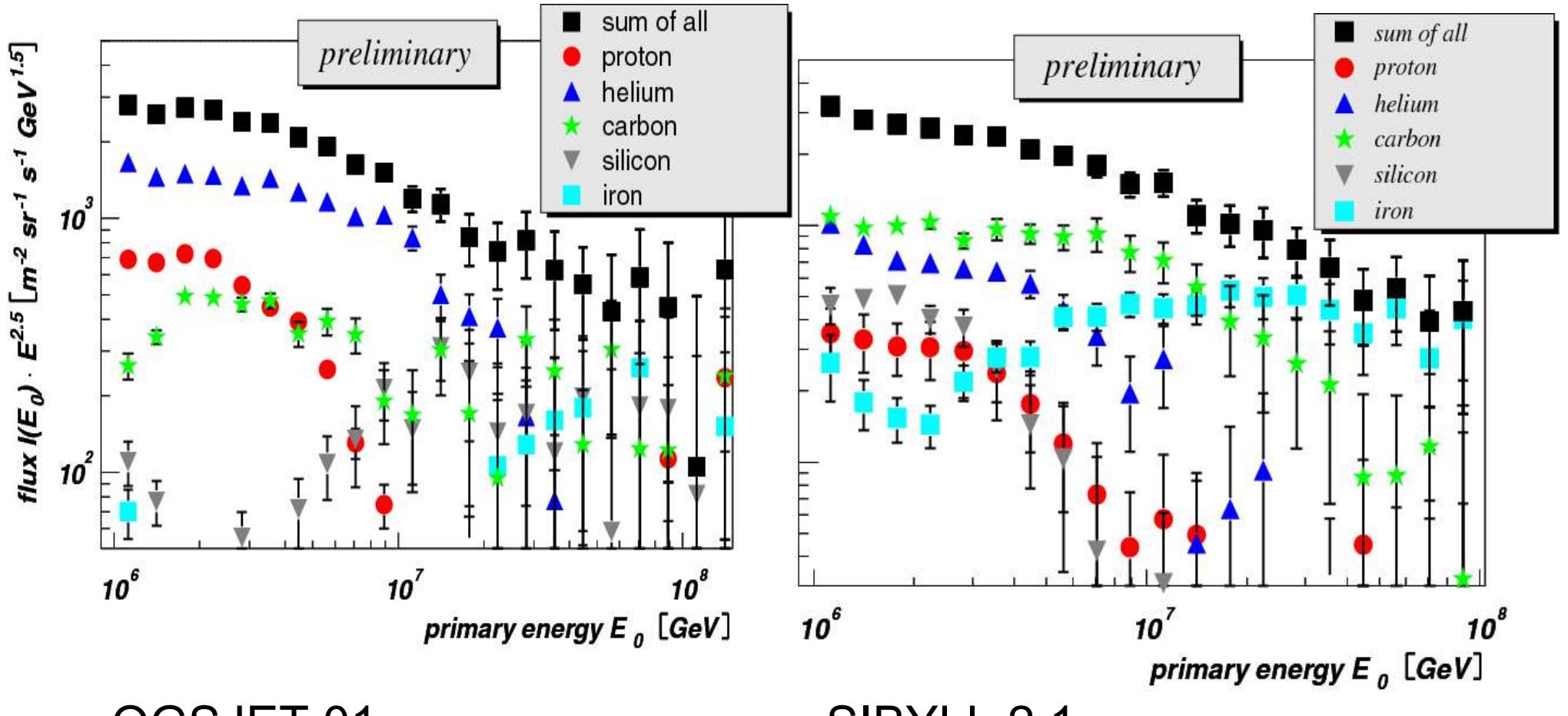
- ➡ Hadronic cross section
- ➡ Inelasticity
- ➡ Diffraction dissociation
- ➡ Total multiplicity
- ➡ Ratio charged/total number of particles
- ➡

Implies large (>10 %) uncertainties in :

- ➡ Energy estimation from surface detectors (LDF)
- ➡ Mass estimation from X_{\max} or from N_e or N_μ

Only energy estimation from calorimetric measurement are almost MC independent.

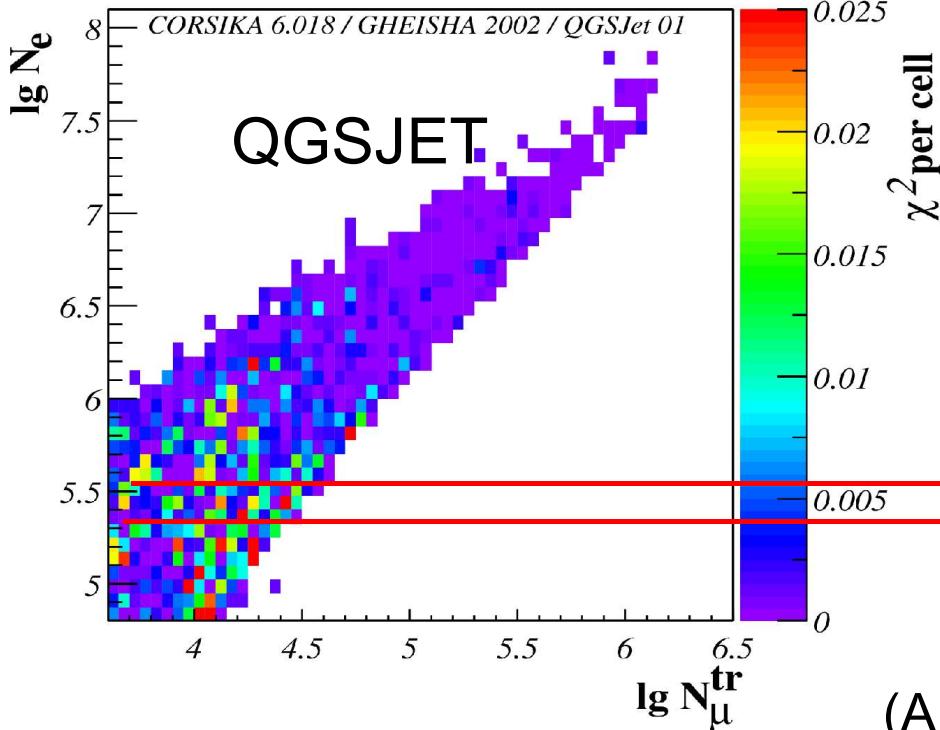
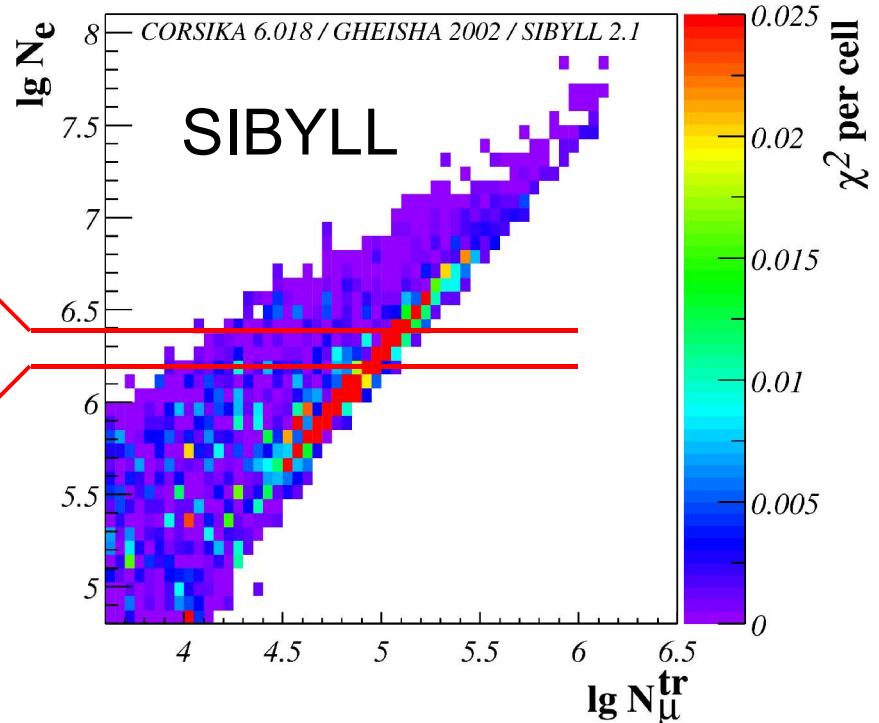
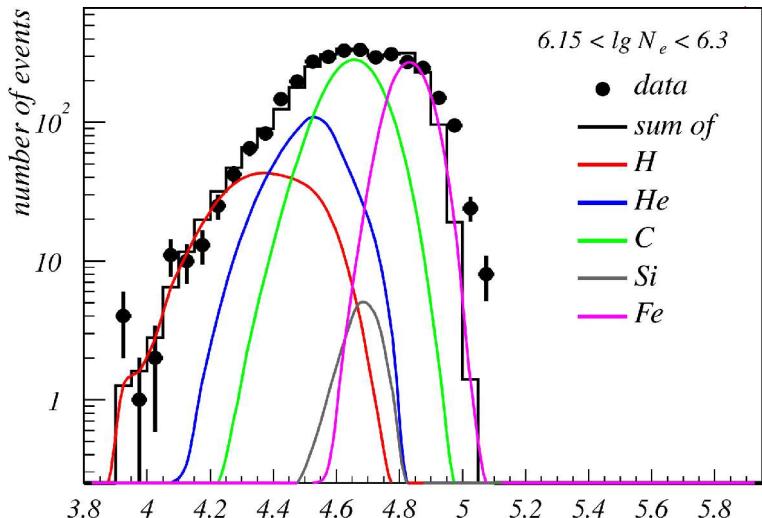
Example: KASCADE Composition



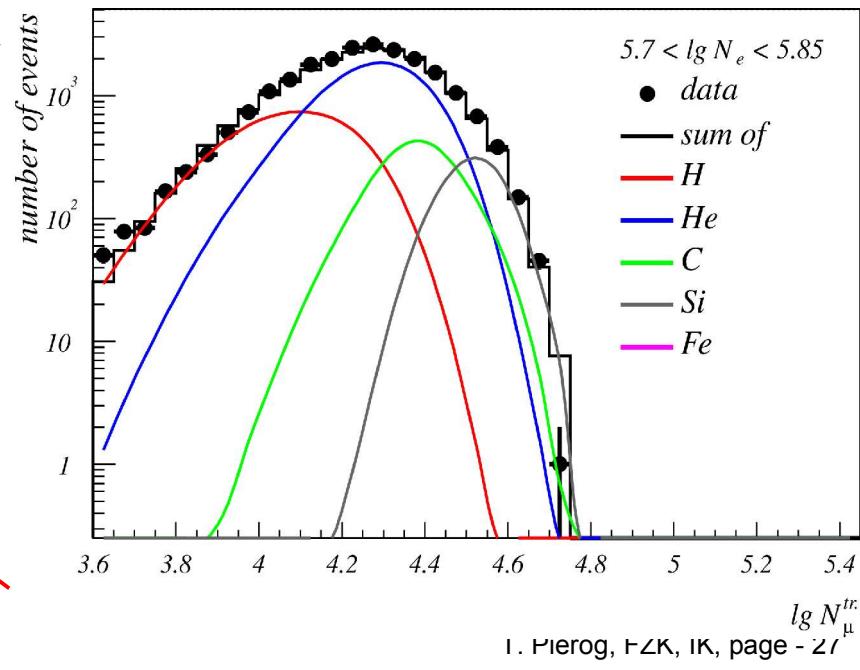
KASCADE: high resolution/precision data,
analysis limited by model uncertainties

Energy spectra KASCADE:

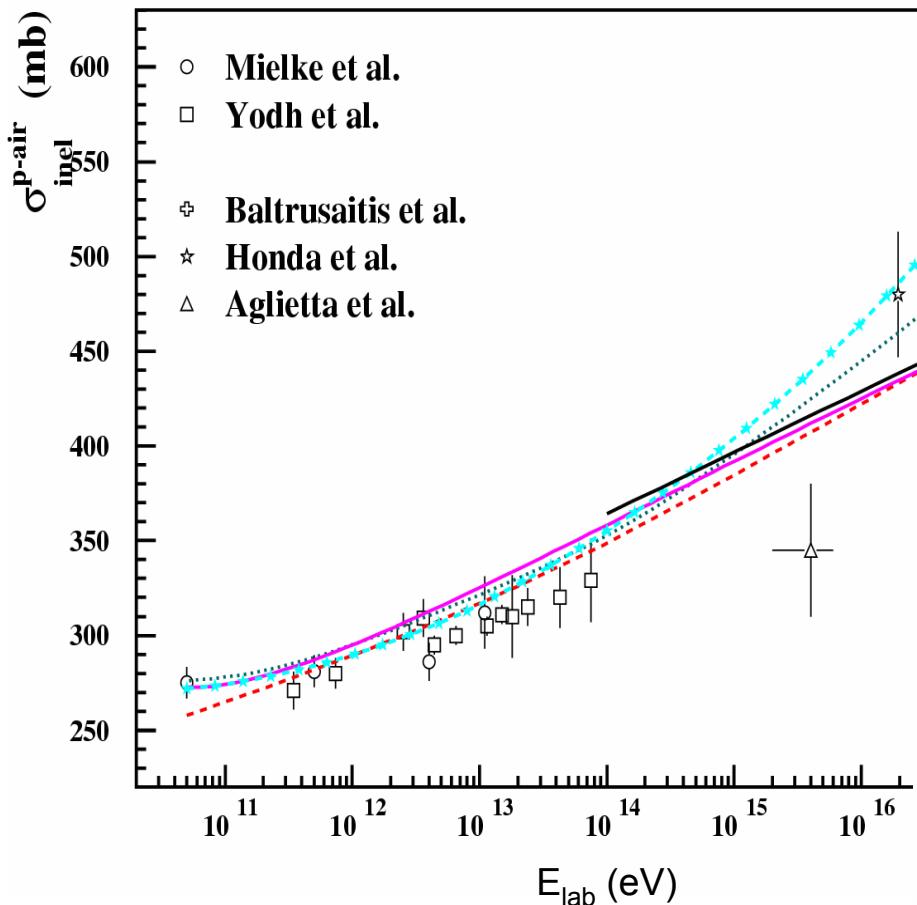
quality of description of data



(A. Haungs)

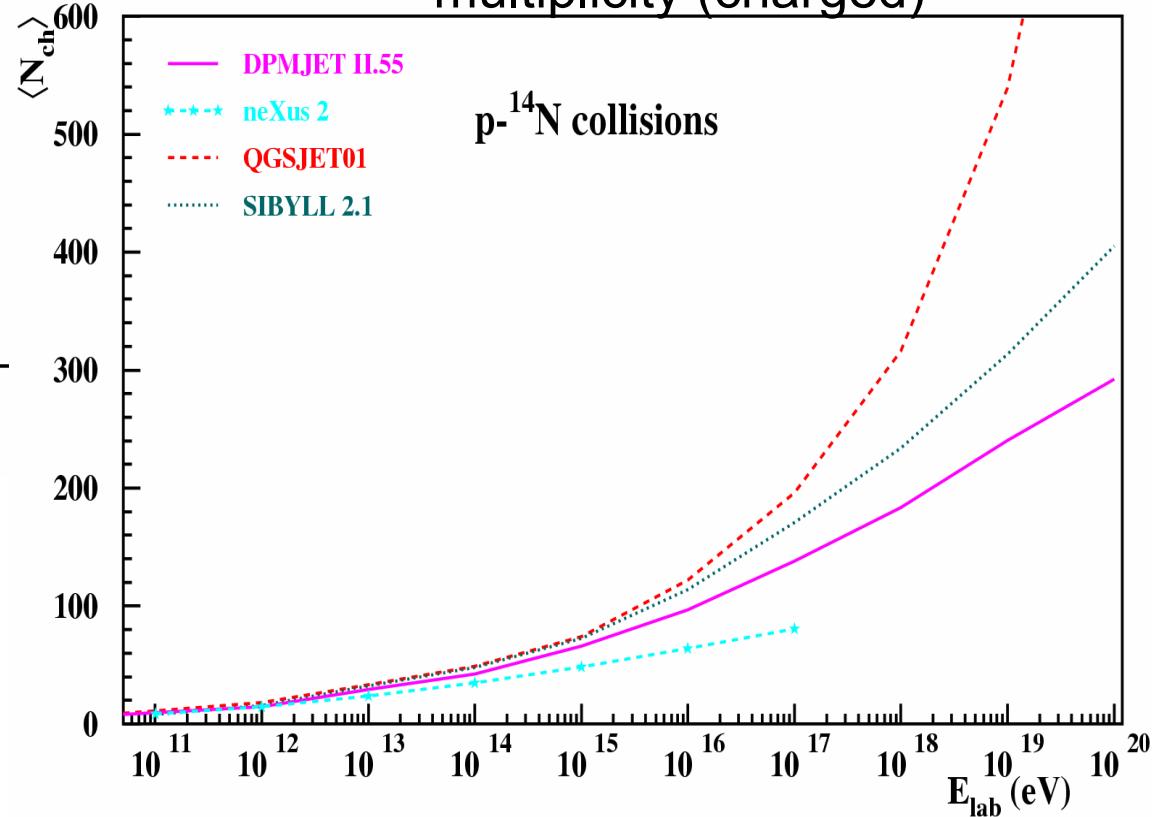


Cross Section and Multiplicity Predictions



Proton-air production cross section

Average secondary particle multiplicity (charged)



QGSJET:

- rapid increase of multiplicity with energy
- smallest elongation rate of all models