

The COMPASS Experiment and the SPIN structure of the NUCLEON

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Trieste University and INFN

Prague Scientific Symposium 26/10/2012





*CO*mmon
Muon and
Proton
Apparatus for
Structure and
Spectroscopy

fixed target experiment at the CERN SPS



SPS

LHC



*CO*mmon
Muon and
Proton
Apparatus for
Structure and
Spectroscopy

fixed target experiment at the CERN SPS



physics programme:

hadron spectroscopy (p , π , K)★

- light mesons, glue-balls, exotic mesons
- polarisability of pion and kaon

nucleon structure (μ)

- longitudinal spin structure
- transverse momentum and transverse spin structure

The Spectrometer

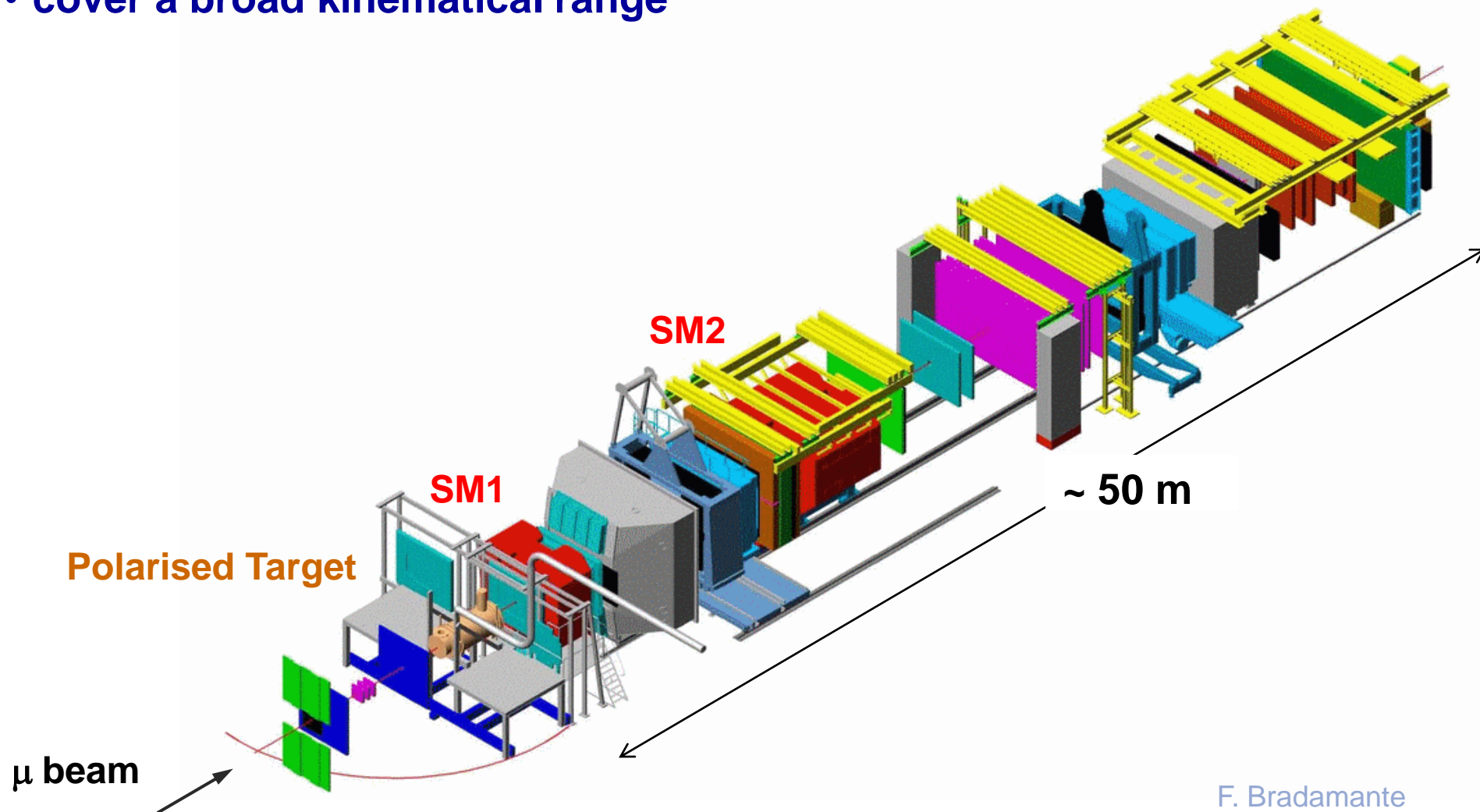


designed to

- use high energy beams
- have large angular acceptance
- cover a broad kinematical range

two stages spectrometer

- Large Angle Spectrometer (**SM1**)
- Small Angle Spectrometer (**SM2**)





The Spectrometer

designed to

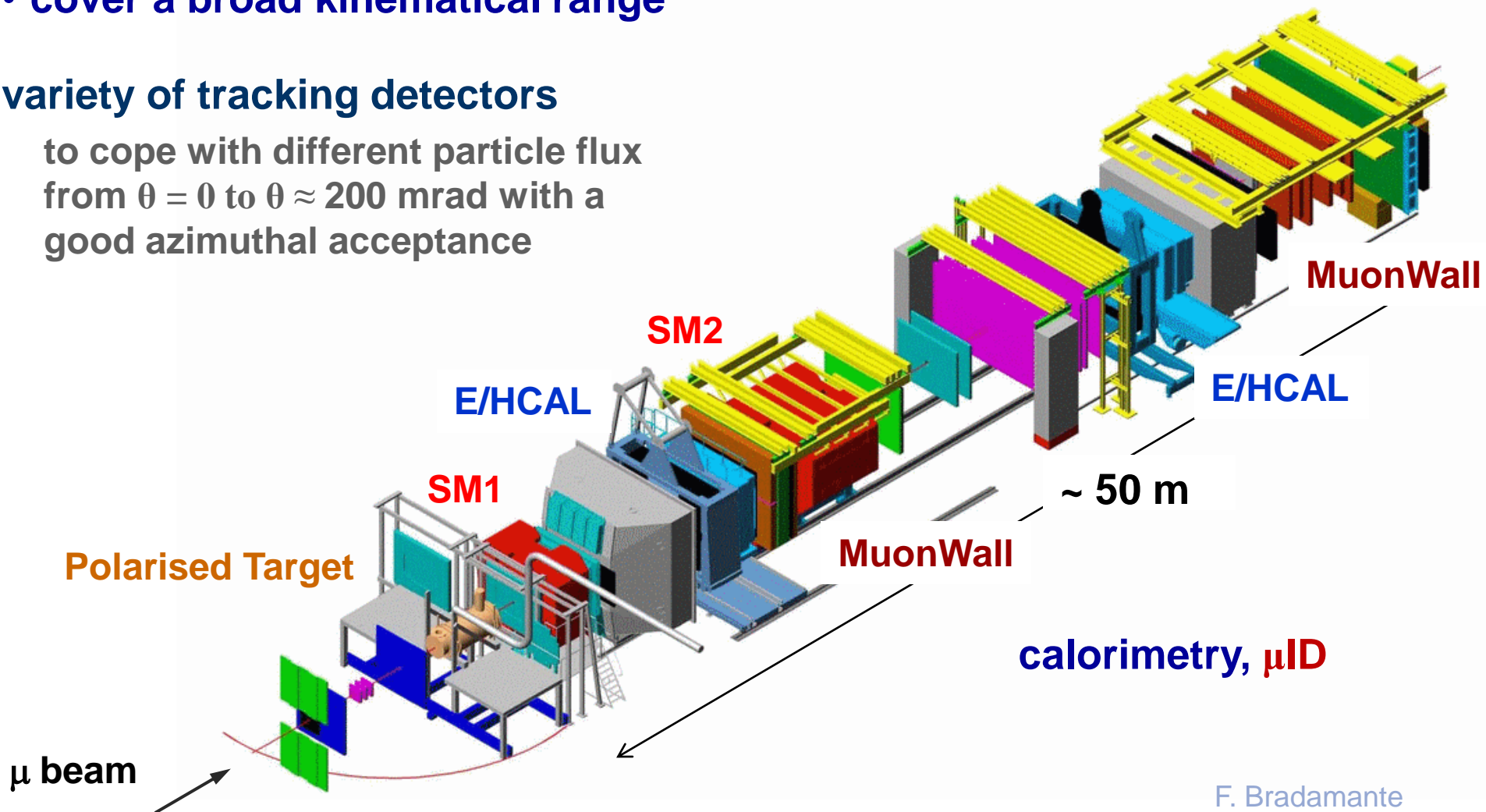
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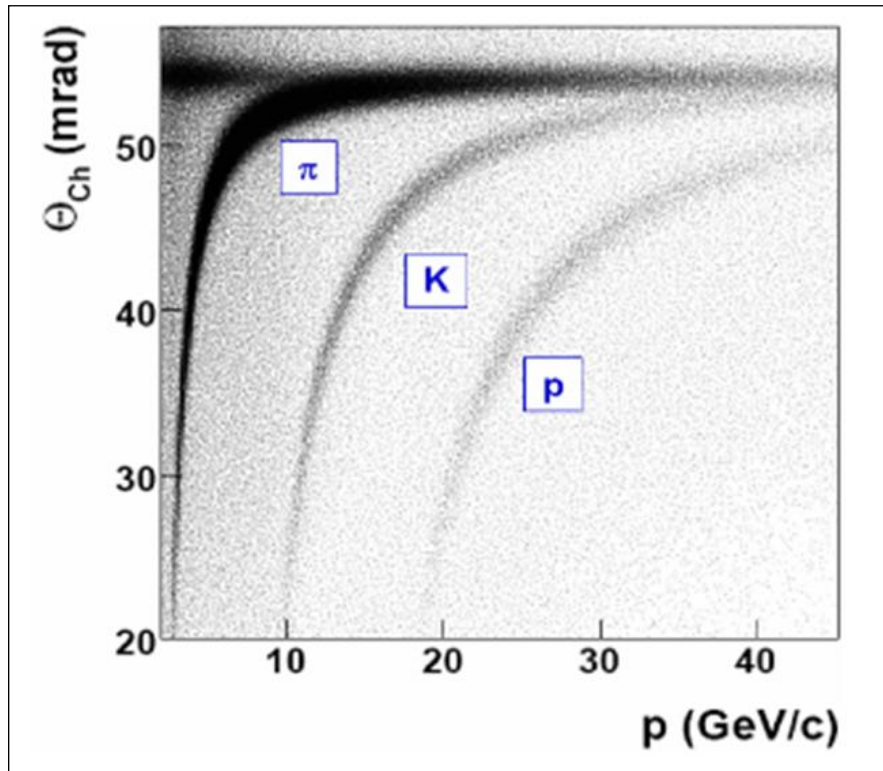
variety of tracking detectors

to cope with different particle flux from $\theta = 0$ to $\theta \approx 200$ mrad with a good azimuthal acceptance



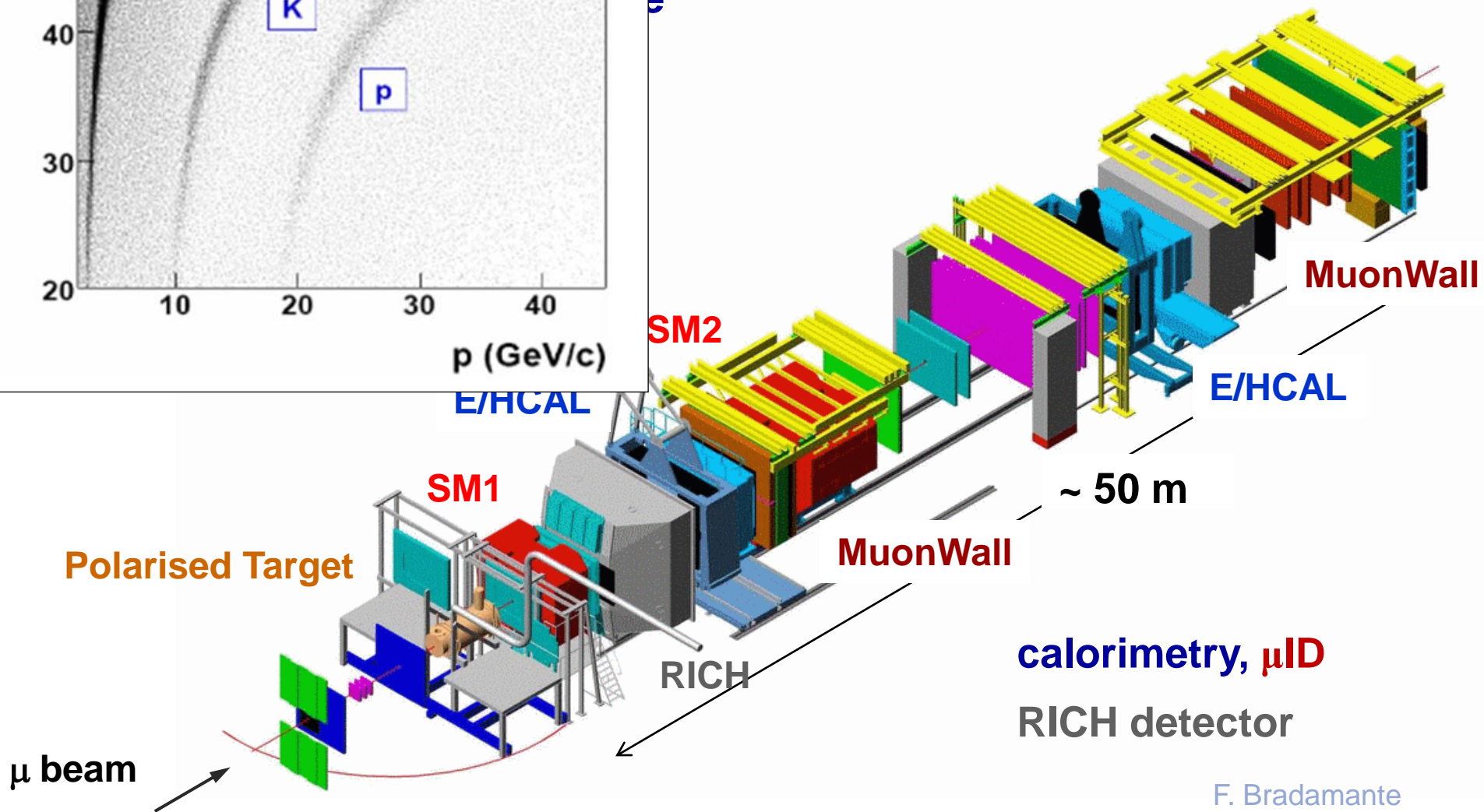


The Spectrometer

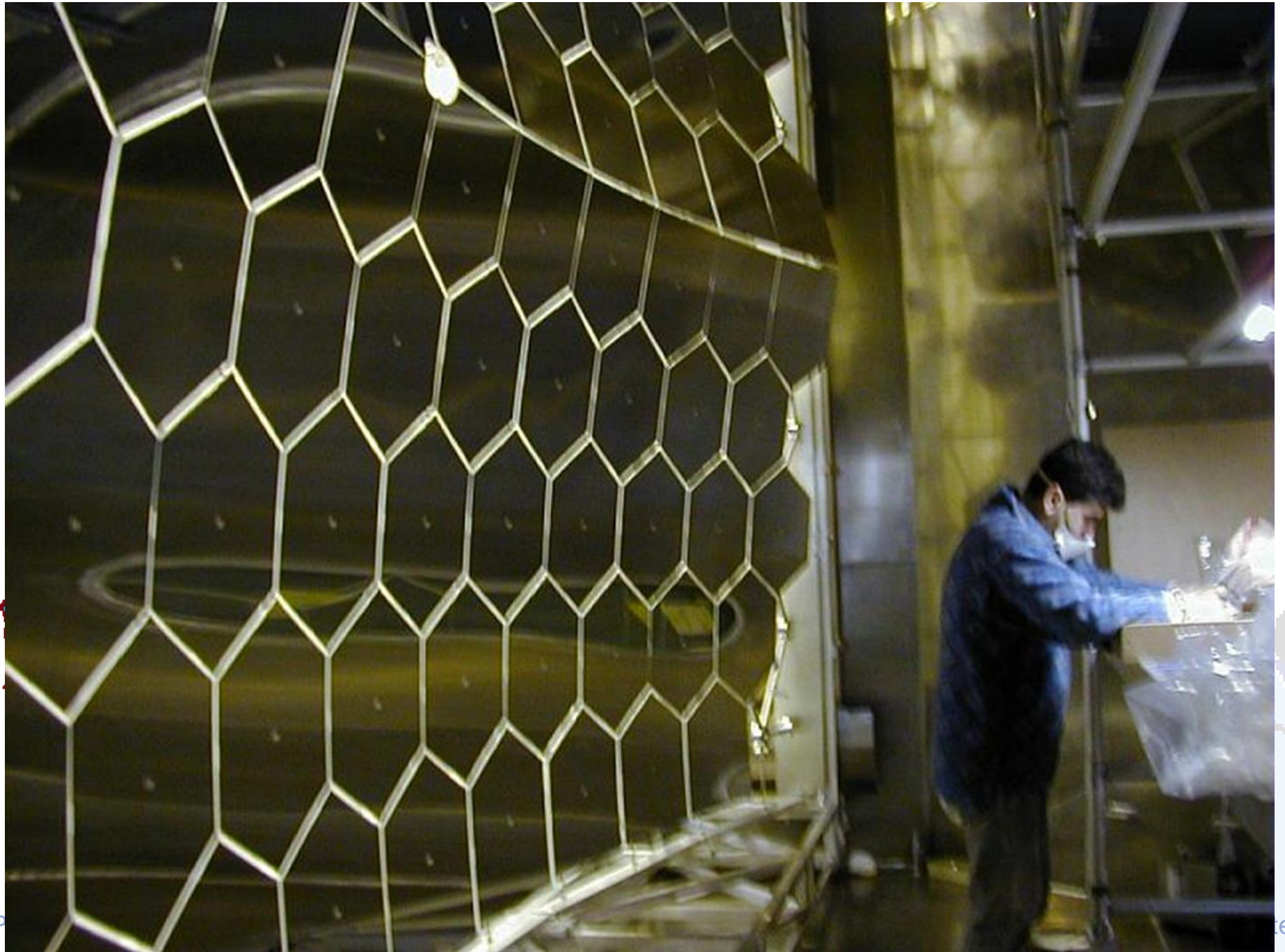


two stages spectrometer

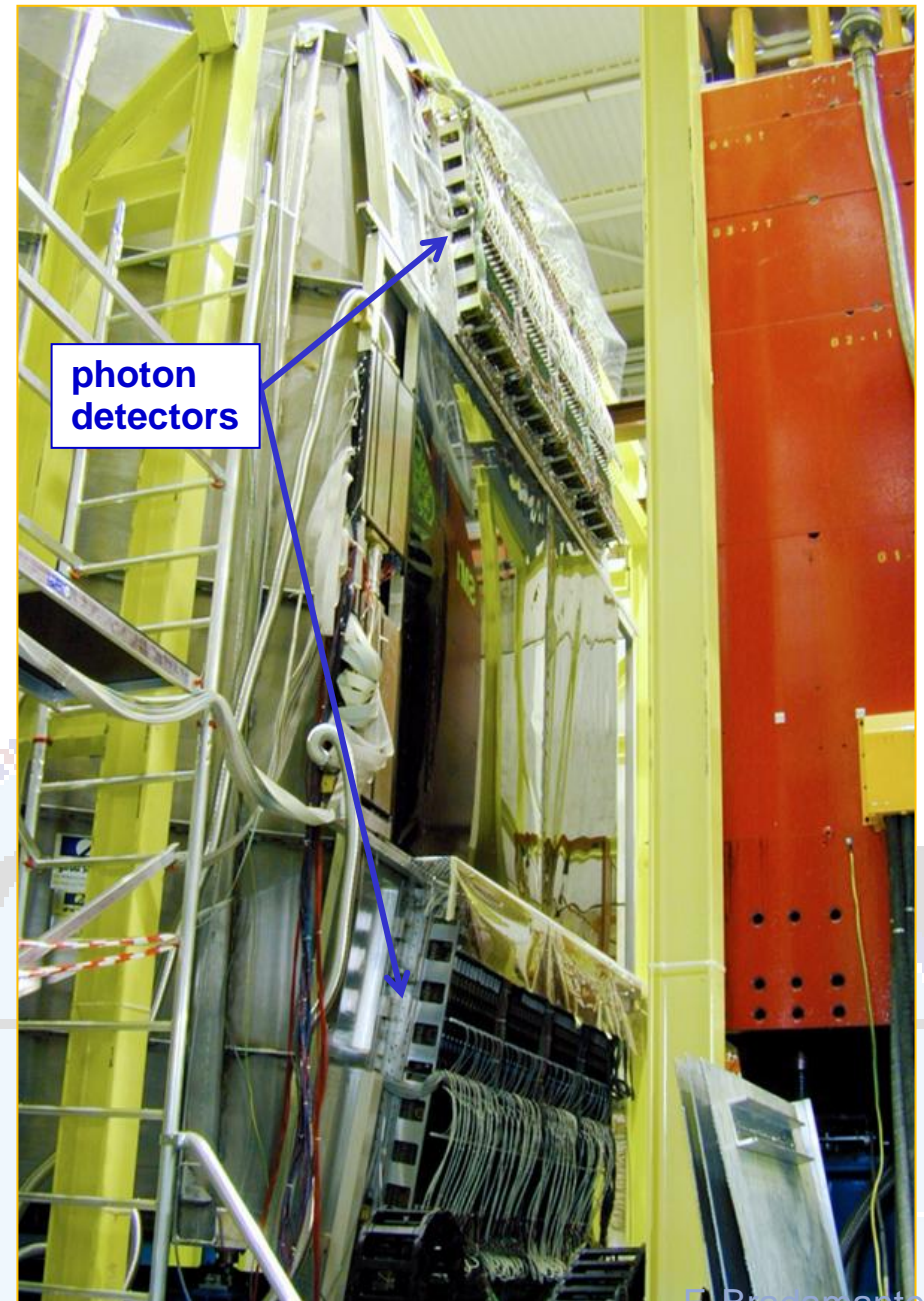
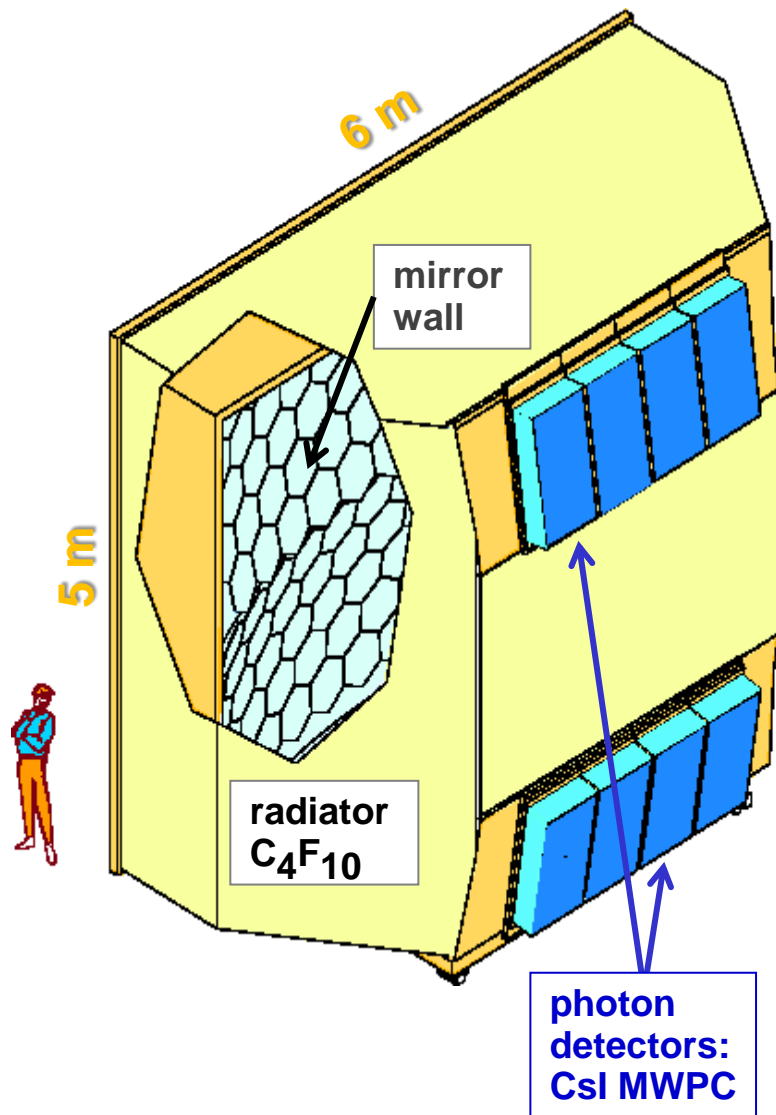
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- Small Angle Spectrometer (**SM2**)



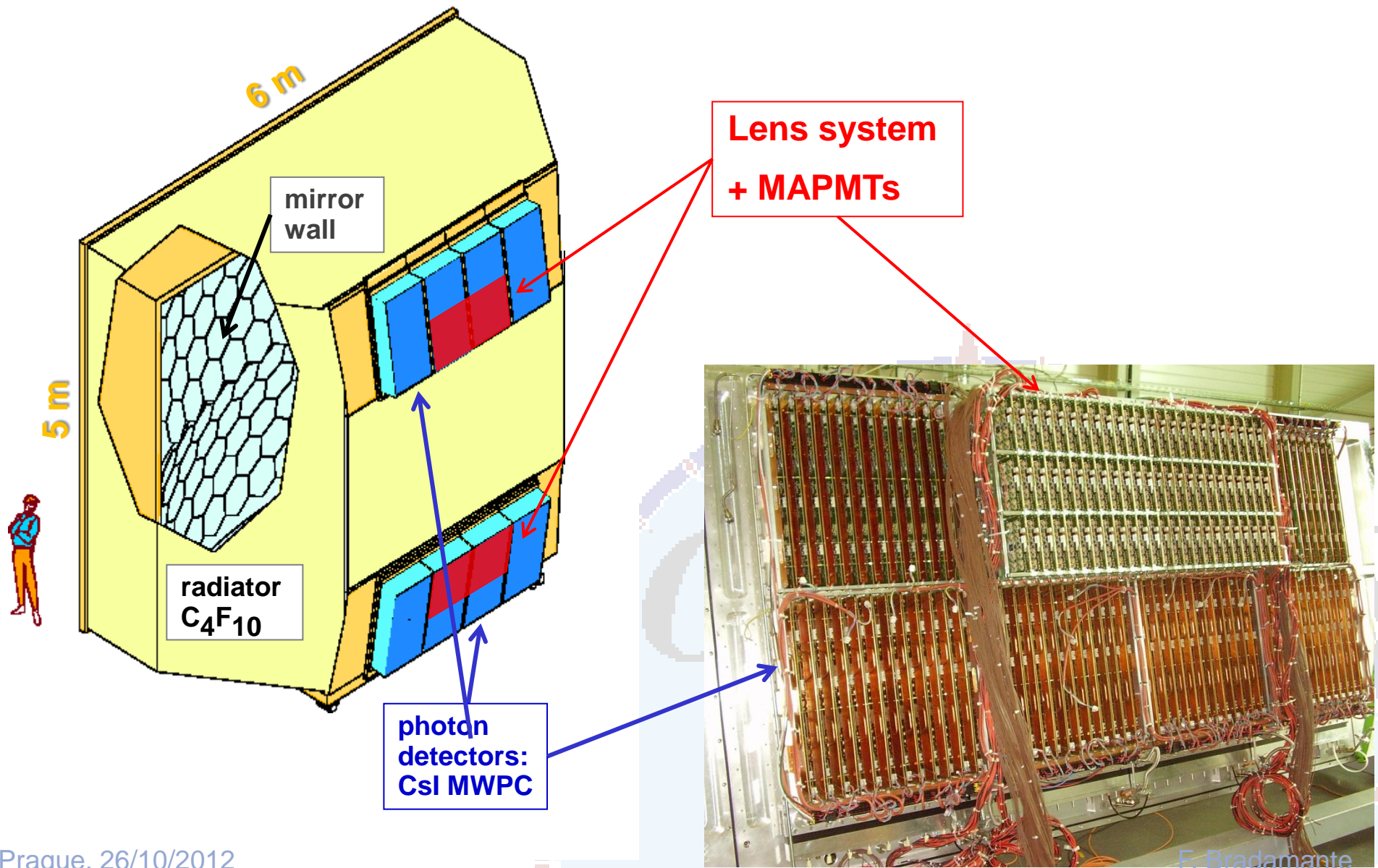
The RICH 1 (<2005)



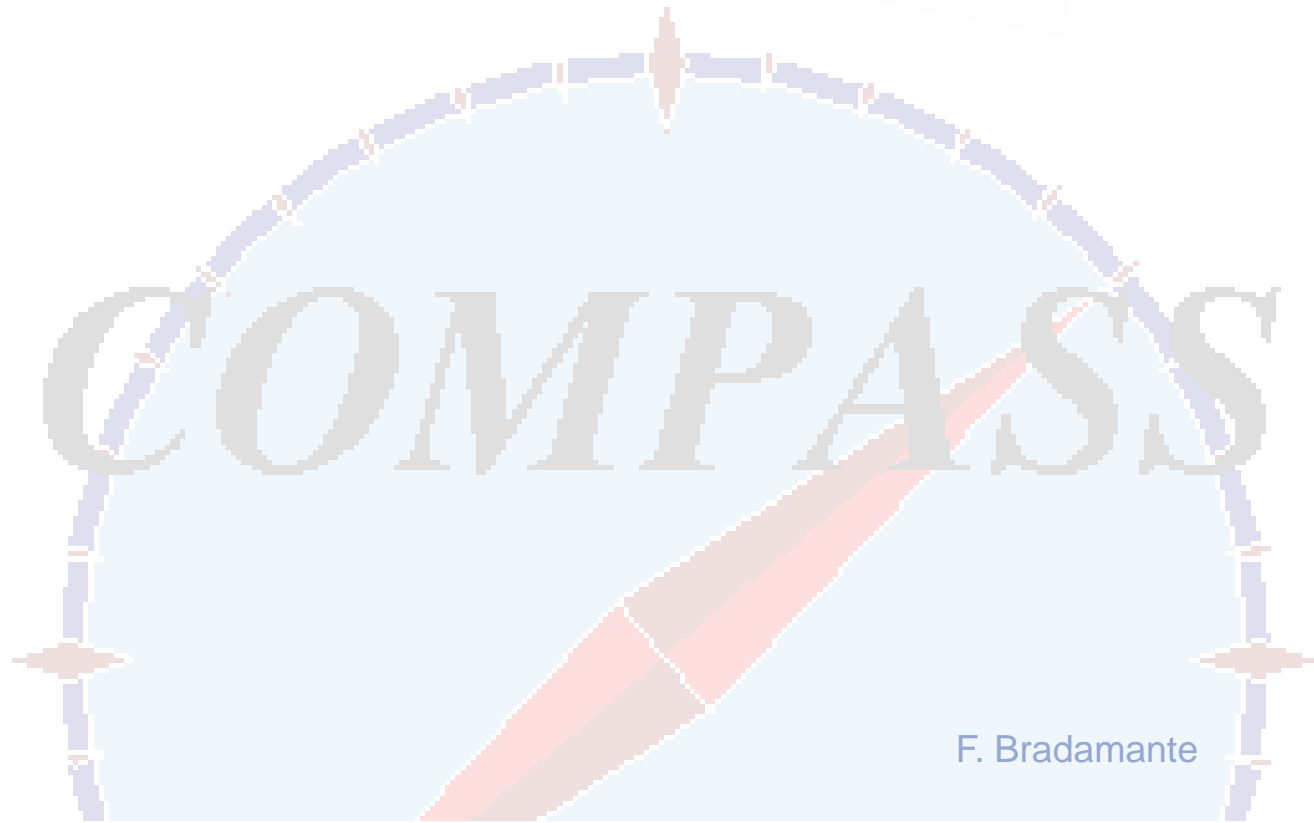
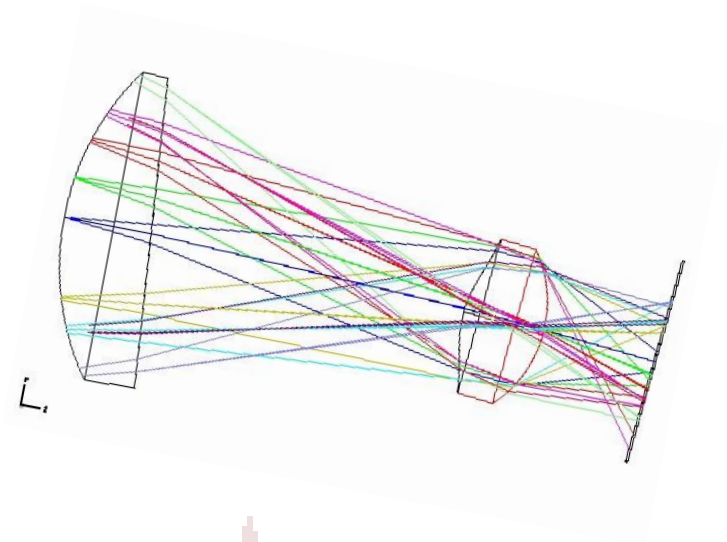
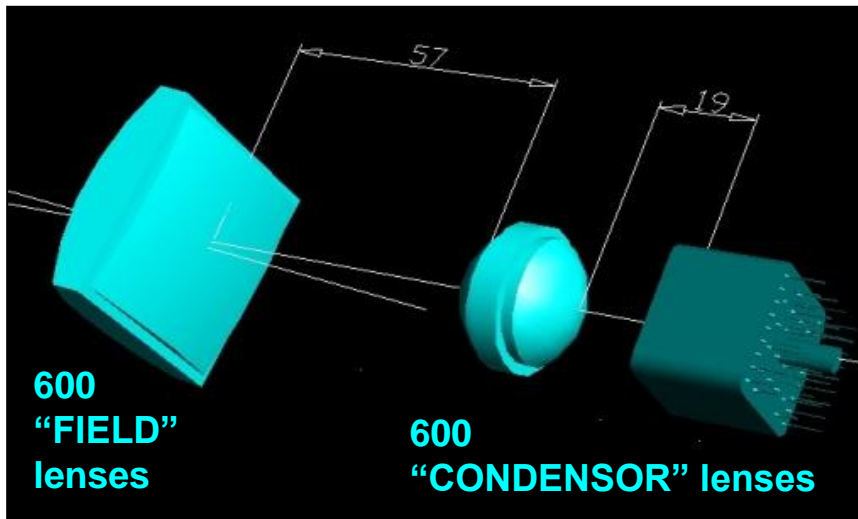
The RICH 1 (<2005)



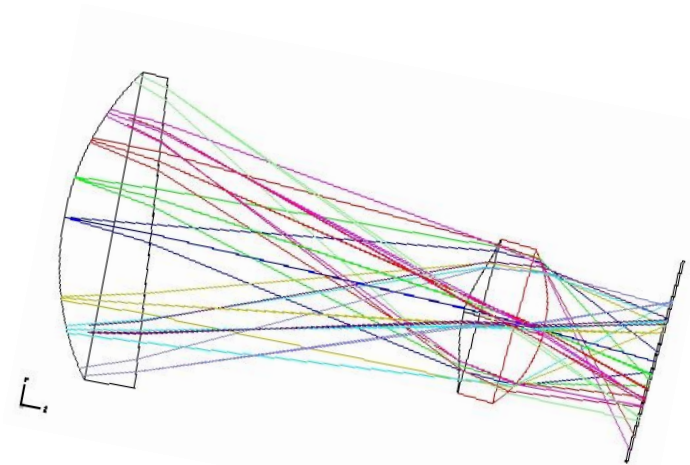
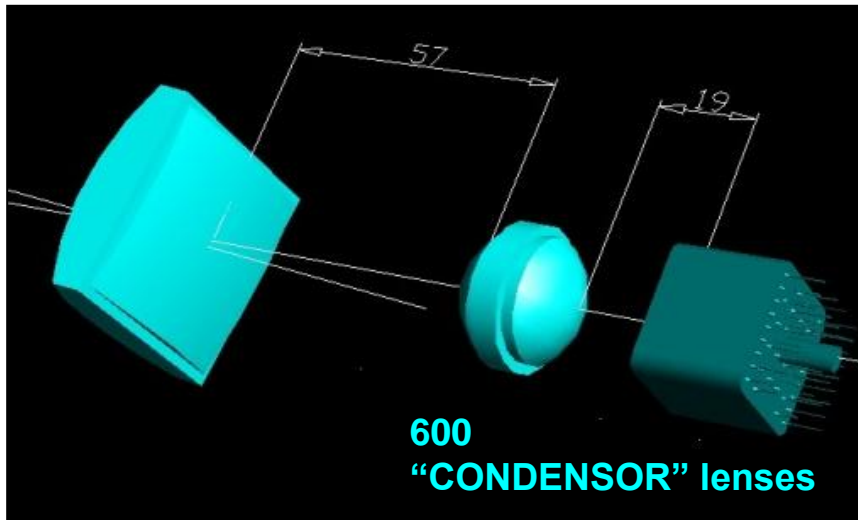
The RICH 1 Upgrade (2005)



The RICH 1 Upgrade (2005)



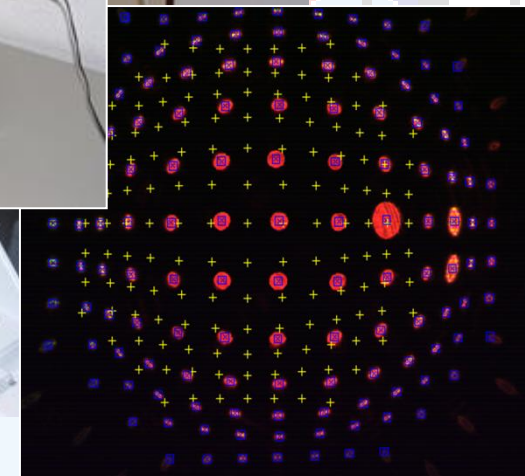
The RICH 1 Upgrade (2005)



optical parameters,
interferometric measurements
**Optical Development
Workshop of CAS
(Turnov, CZ)**

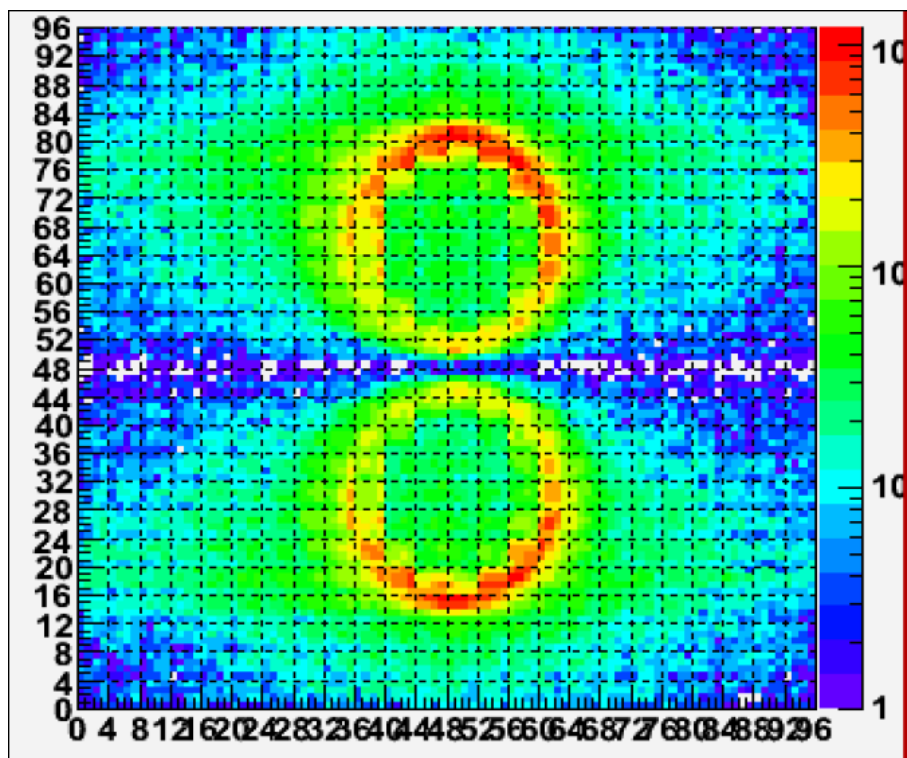


optical parameters,
Hartmann test,
**University of
Liberec (CZ)**

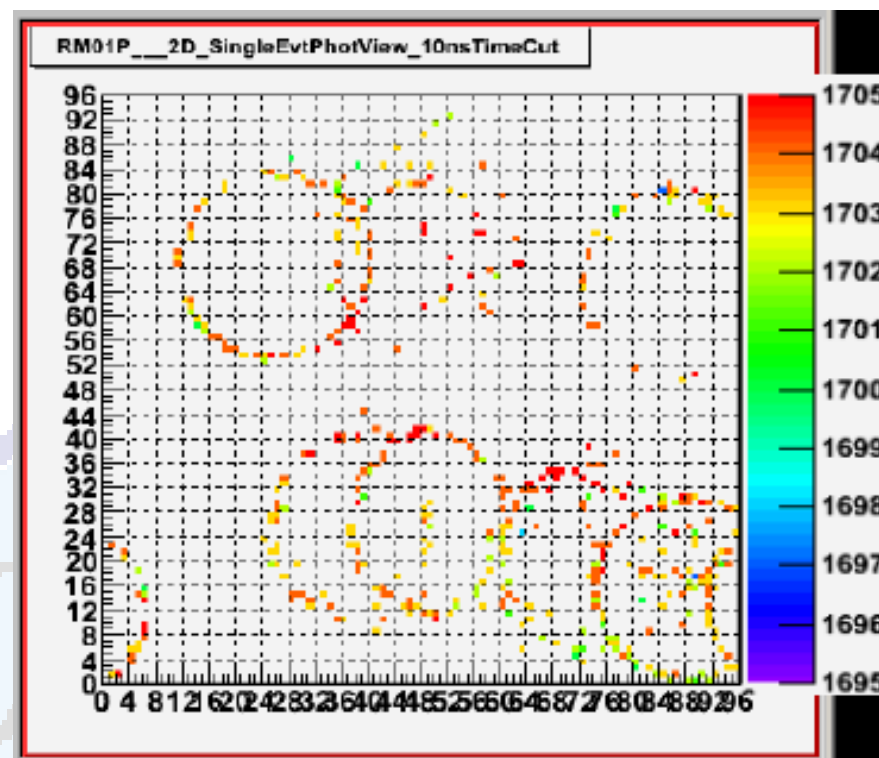


The RICH 1 (>2005)

sum over several events



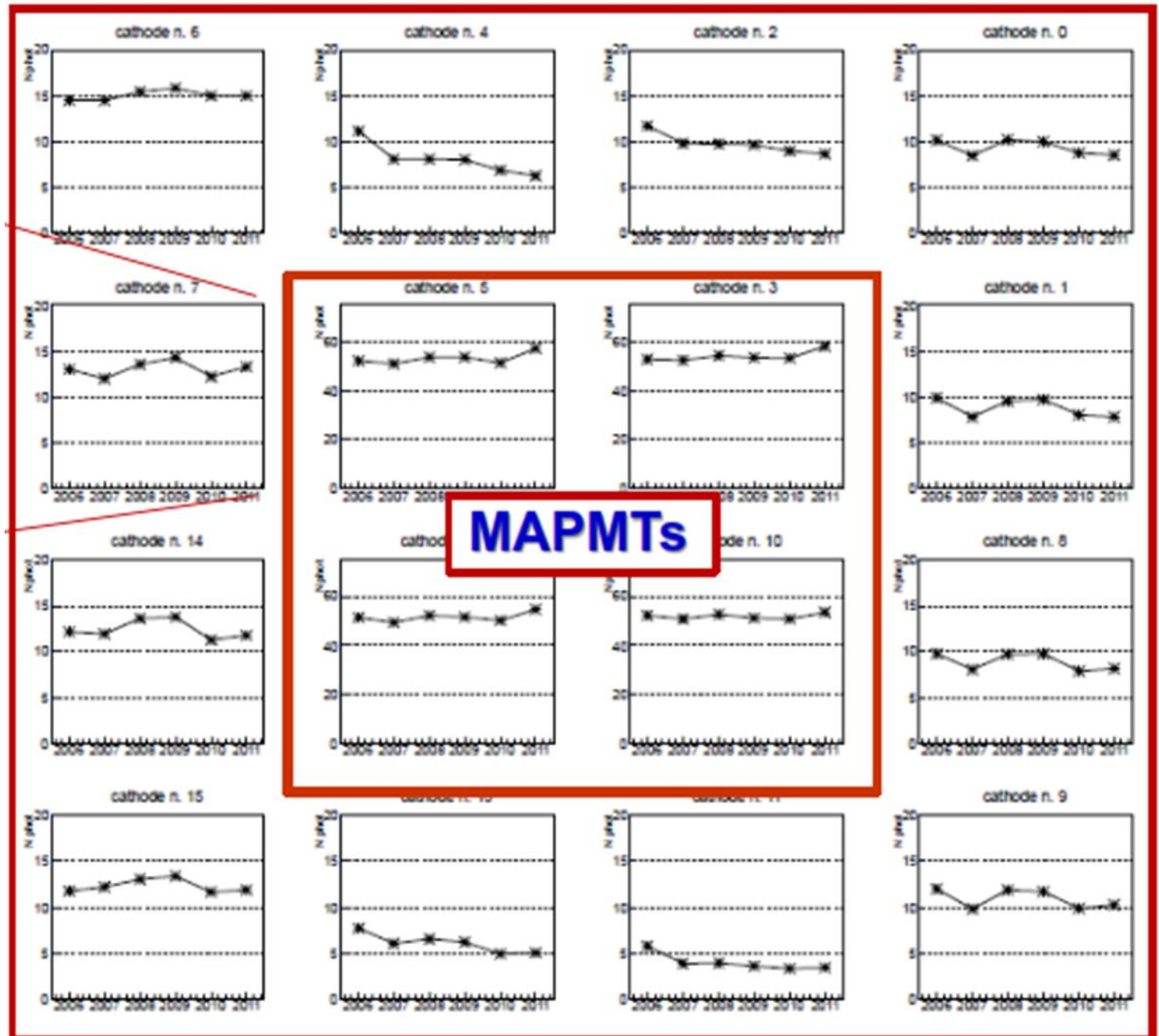
a single event



The RICH 1 (2006-2011)

mean number of photons / ring in the different photon detector sectors

~ 10 on MWPC
~ 60 on MAPMTs



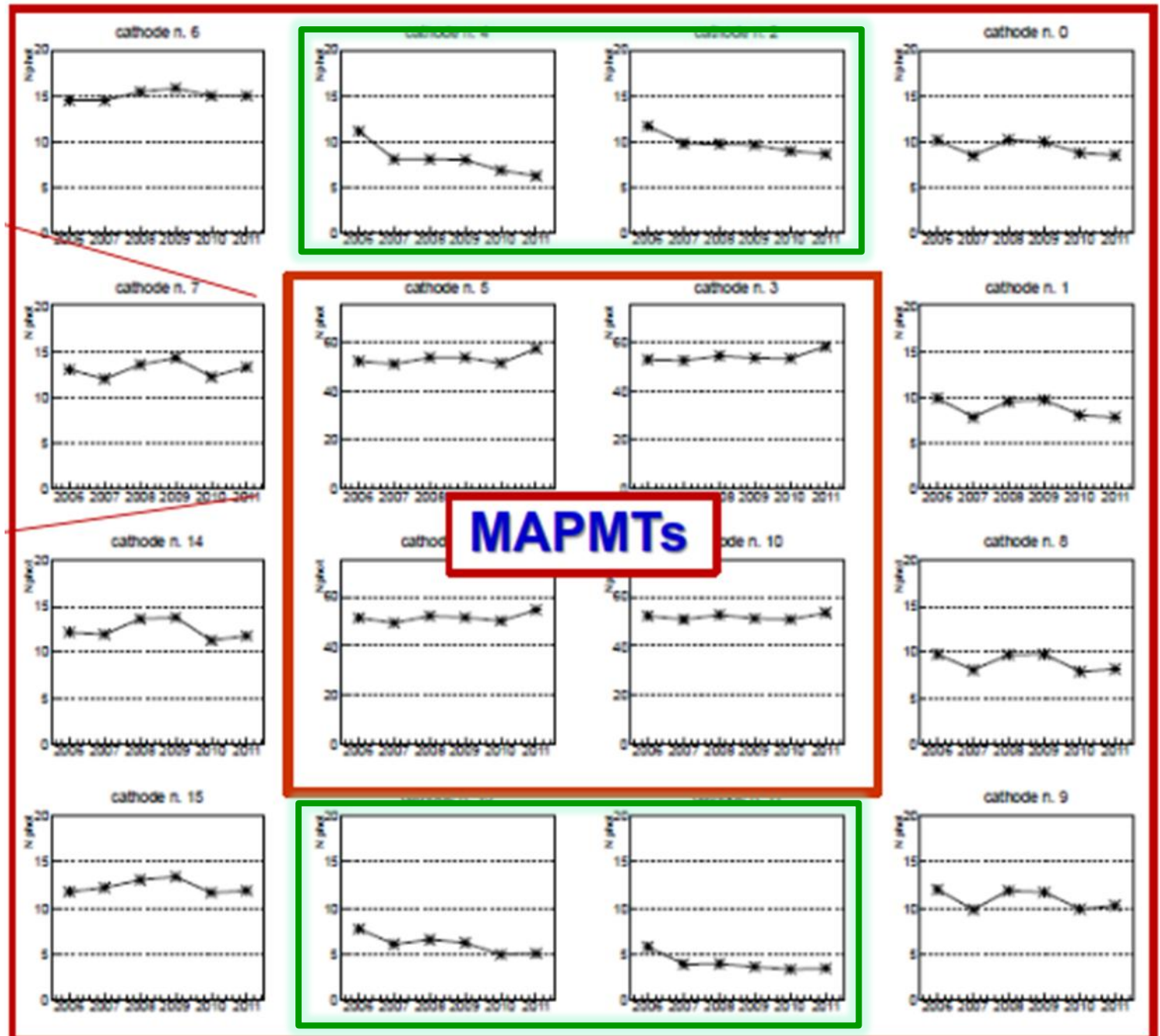
The RICH 1 (2006-2011)

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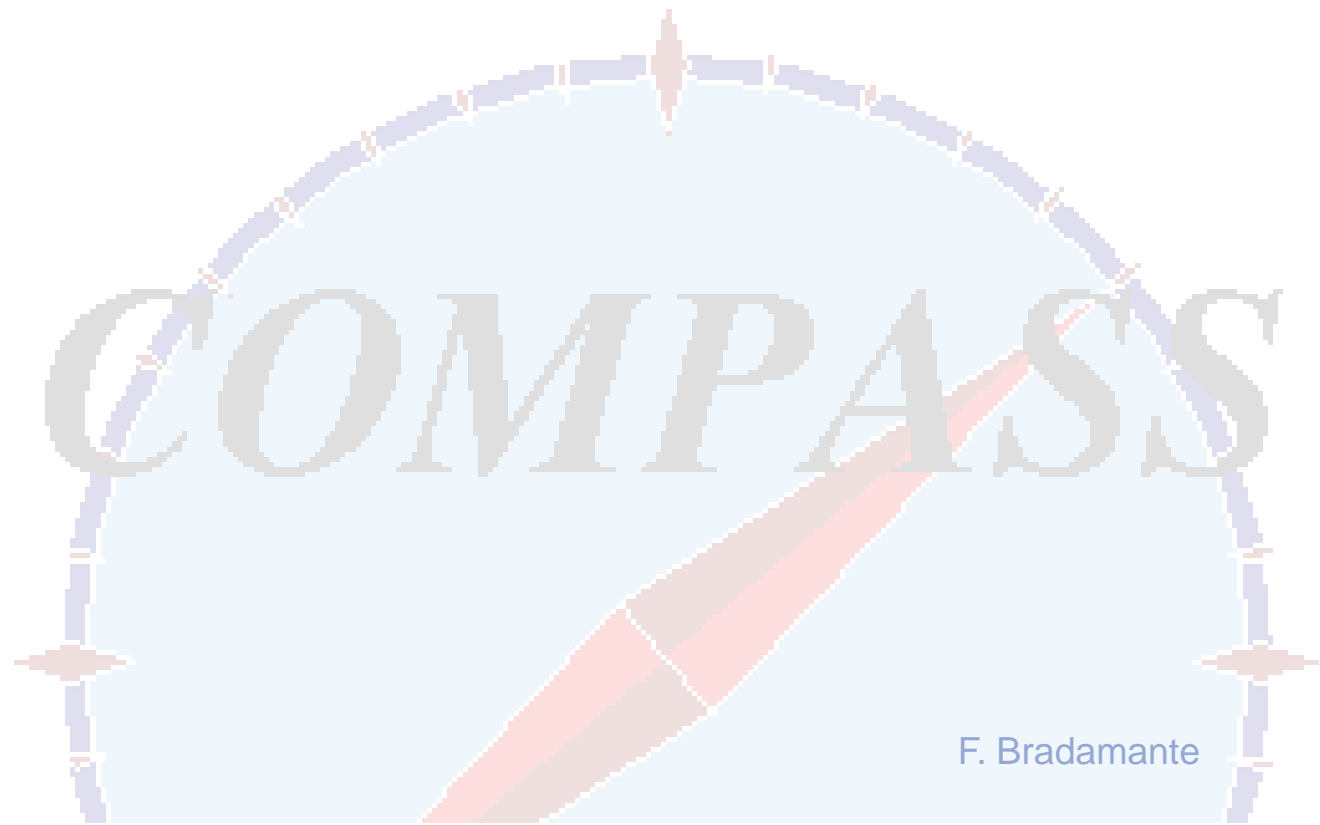
~ 10 on MWPC
~ 60 on MAPMTs

upgrade for DVMP/SIDIS 2015-...

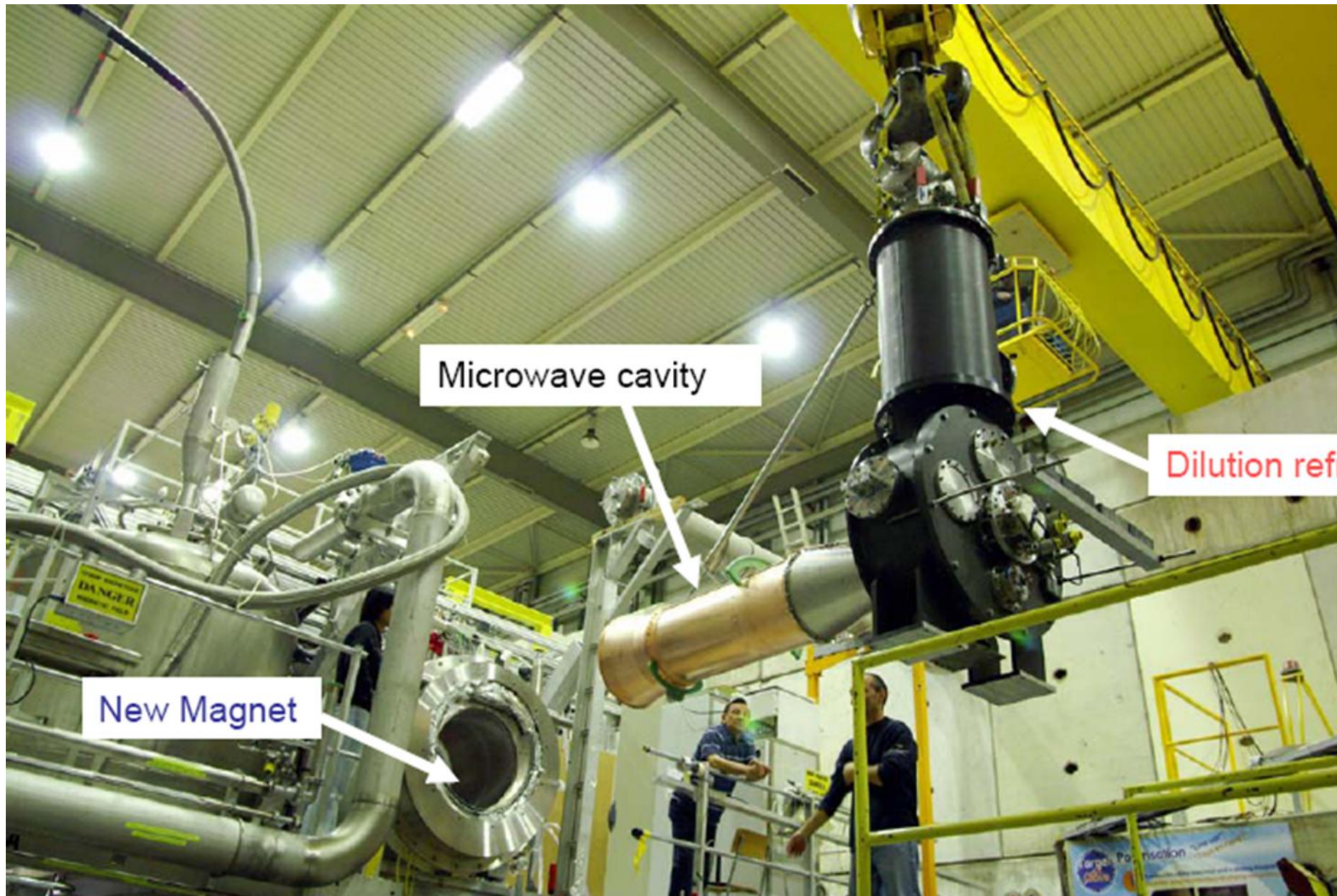
Trieste
Turin
Freiburg
Prague



The Polarized Target System (>2005)



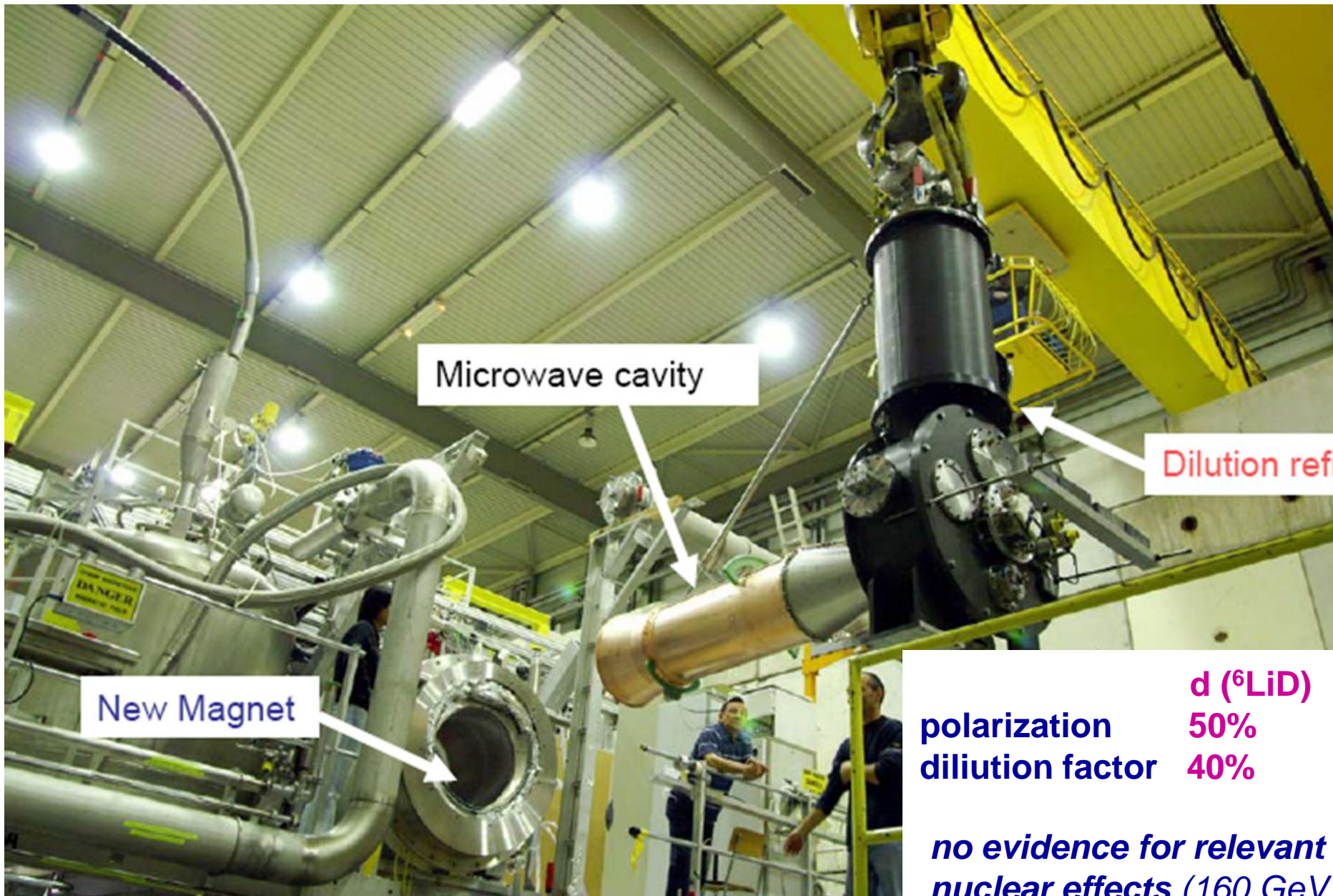
The Polarized Target System (>2005)



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The Polarized Target System (>2005)



Microwave cavity

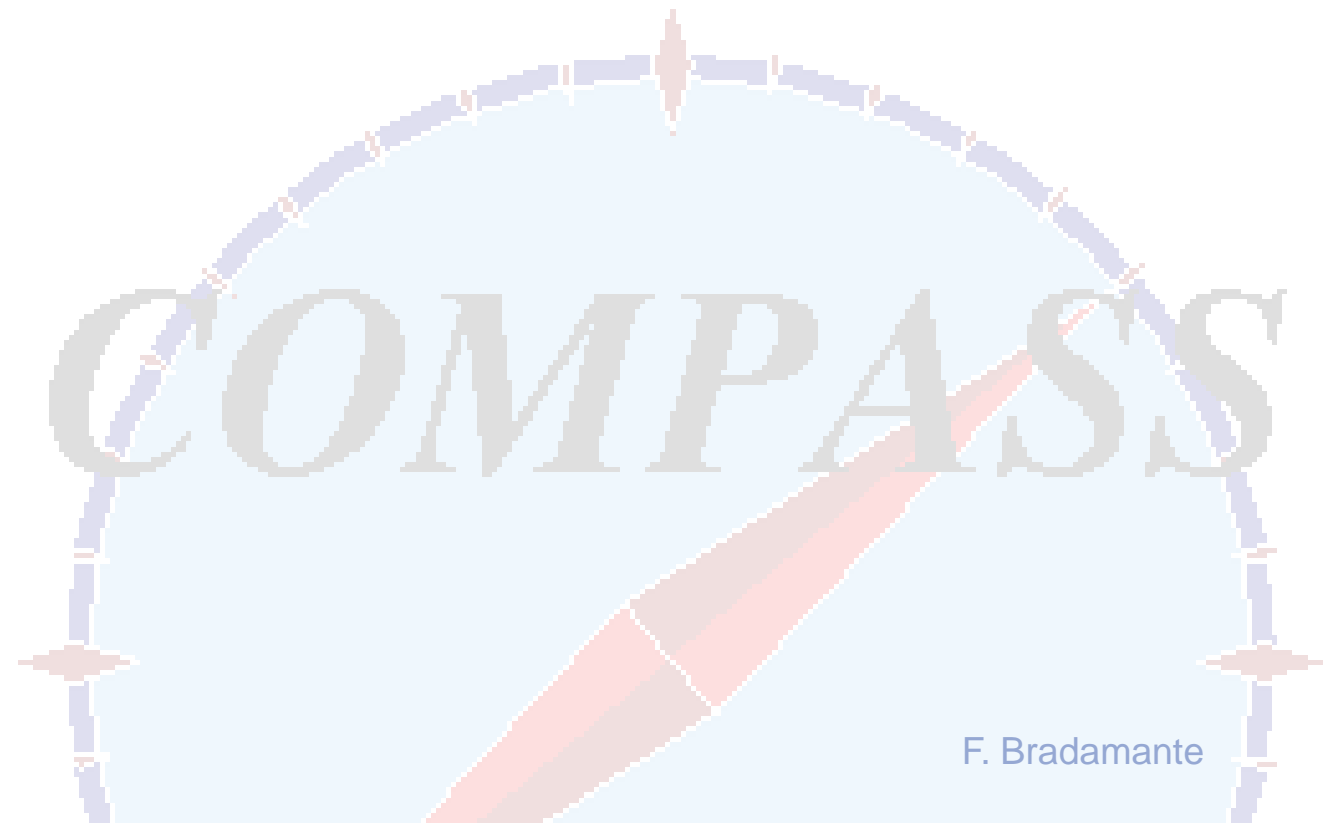
Dilution refrigerator

New Magnet

	d (⁶LiD)	p (NH₃)
polarization	50%	90%
dilution factor	40%	16%

no evidence for relevant nuclear effects (160 GeV)

Data Taking



Data Taking

2002	nucleon structure with	160 GeV μ	L&T	polarised deuteron target
2003	nucleon structure with	160 GeV μ	L&T	polarised deuteron target
2004	nucleon structure with	160 GeV μ	L&T	polarised deuteron target
2005	<i>CERN accelerators shut down</i>			
2006	nucleon structure with	160 GeV μ	L	polarised deuteron target
2007	nucleon structure with	160 GeV μ	L&T	polarised proton target
2008	hadron spectroscopy			
2009	hadron spectroscopy			
2010	nucleon structure with	160 GeV μ	T	polarised proton target
2011	nucleon structure with	190 GeV μ	L	polarised proton target
2012	Primakoff & DVCS / SIDIS test			

COMPASS Physics case

... from a talk at SPIN 2000

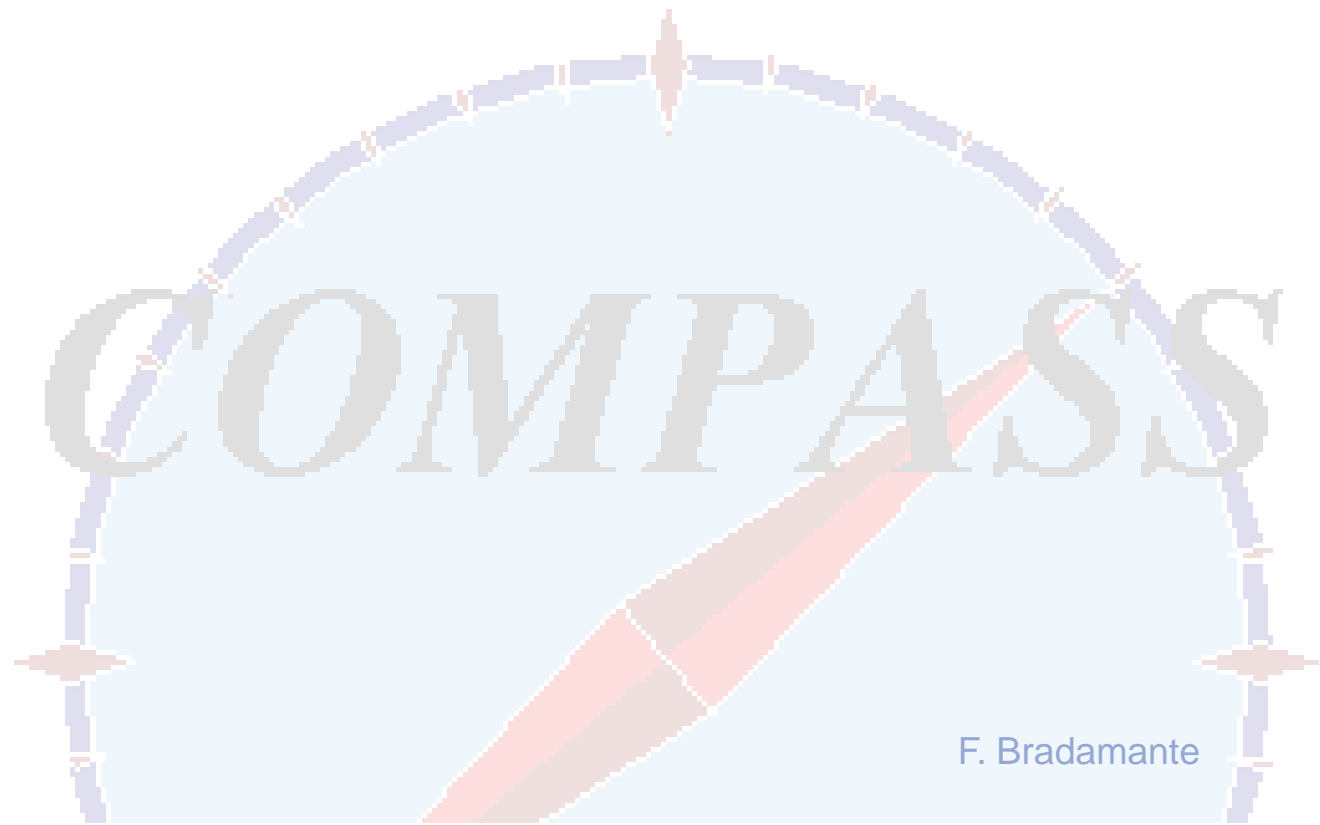
acknowledgement that it is still worthwhile to invest in trying to answer some questions which have been with us since 30 years

- **how is a nucleon made up ?**
! spin structure !
- **are there non qqq or $q\bar{q}$ hadrons ?**
! exotics !

fundamental problems of QCD

- **role of the axial anomaly to the spin of the proton**
- **glueballs \longleftrightarrow non abelian nature of QCD**

the COMPASS experiment and the HADRON PROGRAM



Hadron Spectroscopy

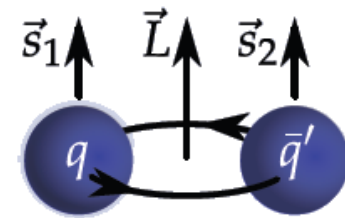
mesons

quantum numbers in CQM

$$S = 0, 1; \quad \vec{J} = \vec{L} + \vec{S}; \quad P = (-1)^{L+1}; \quad C = (-1)^{L+S}$$

forbidden (exotic QN's)

$$J^{PC} = 0^{--}, 0^{+-}, 1^{-+}, 2^{+-}, 3^{-+}, \dots$$



Hadron Spectroscopy

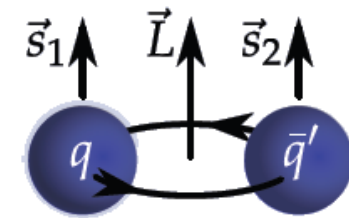
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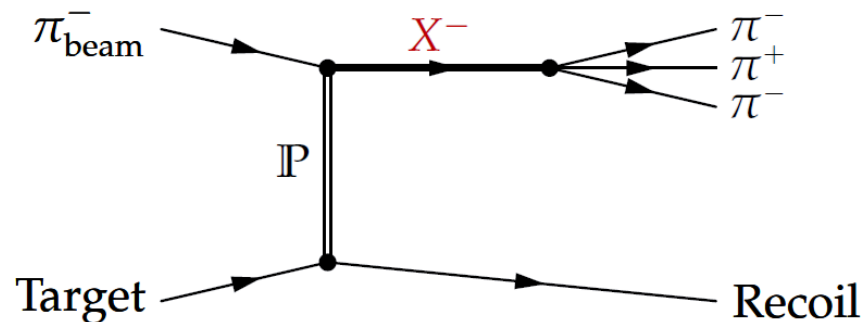
$$J^{PC} = 0^{--}, 0^{+-}, 1^{-+}, 2^{+-}, 3^{-+}, \dots$$



more states in QCD:

hybrids $|q\bar{q}g\rangle$, glueballs $|gg\rangle$, multiquark states $|qq\bar{q}\bar{q}\rangle$

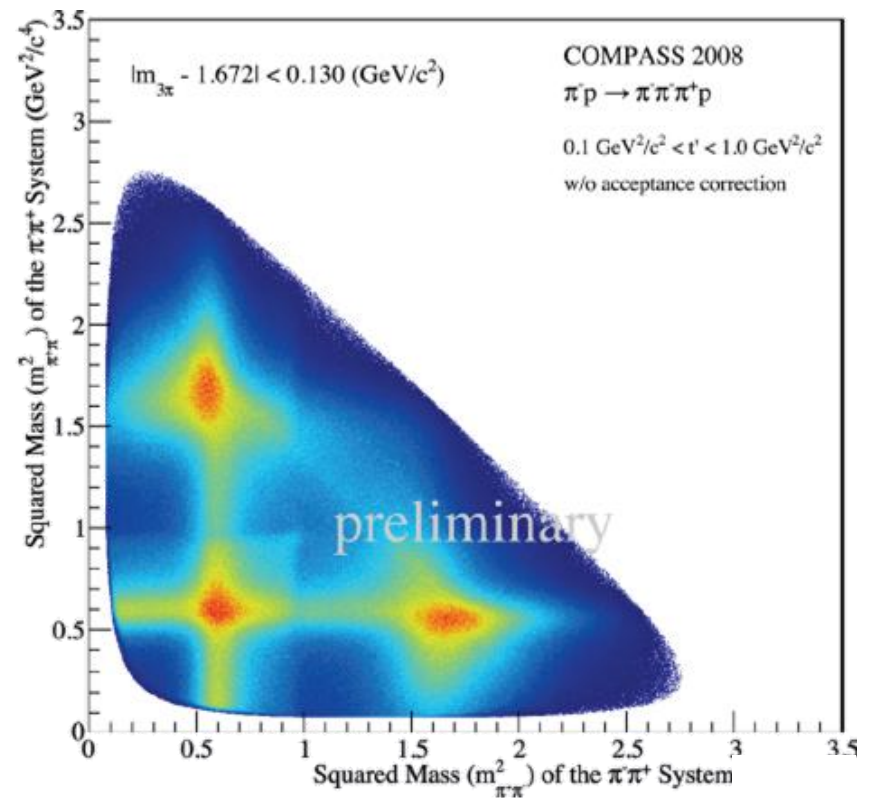
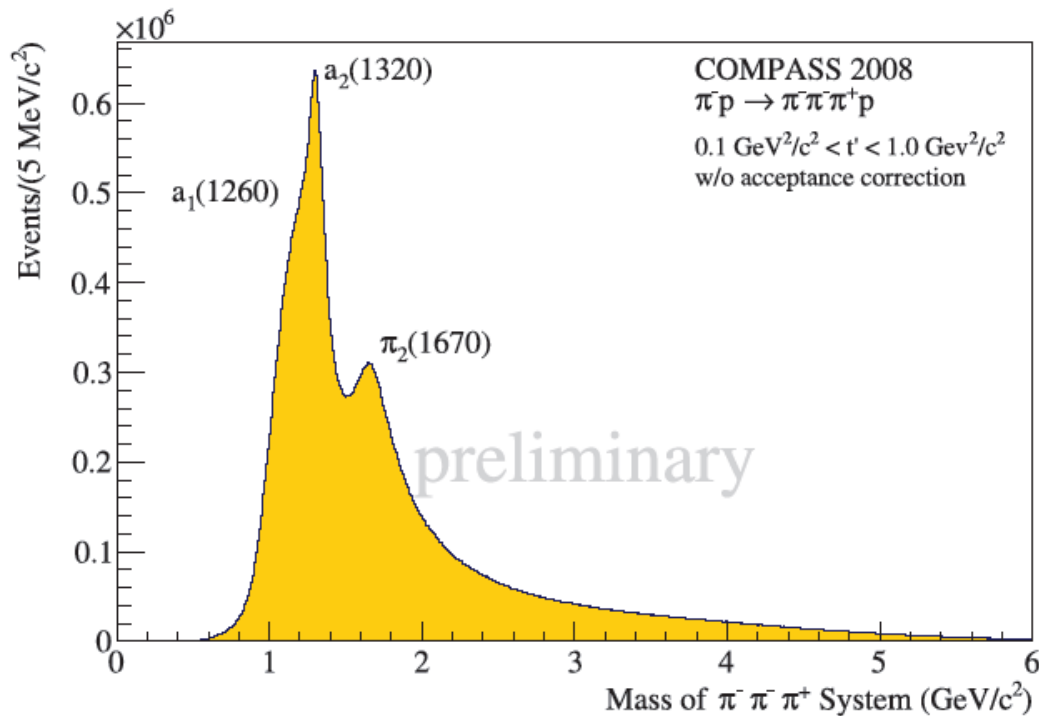
diffractive dissociation:



Hadron Spectroscopy

example: $\pi^- p \rightarrow \pi^- \pi^+ \pi^- p$

sample with 100M events

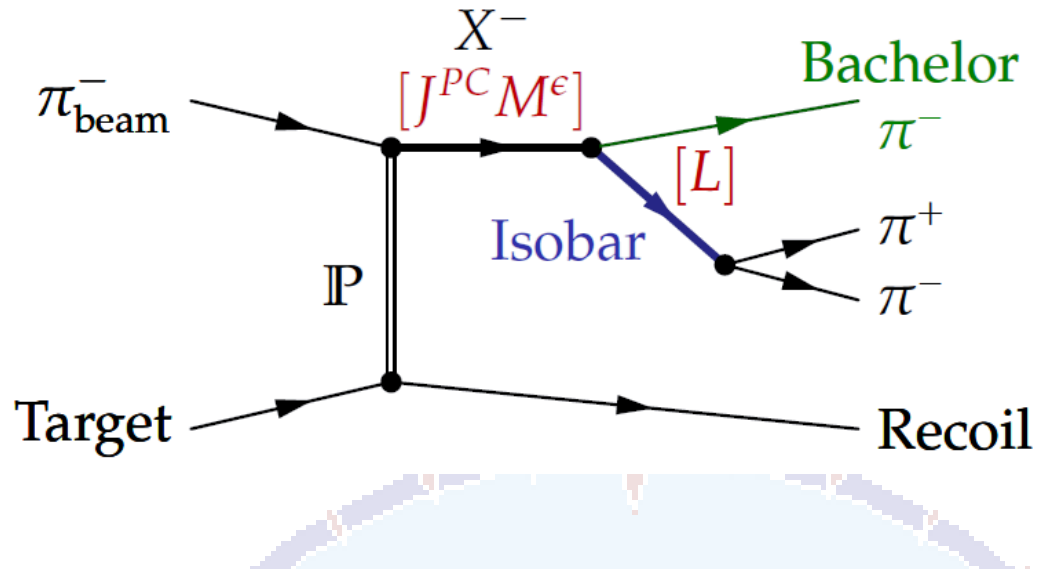


Hadron Spectroscopy

example: $\pi^- p \rightarrow \pi^- \pi^+ \pi^- p$

- **Isobar model:**

X decay is chain of successive two-body decays



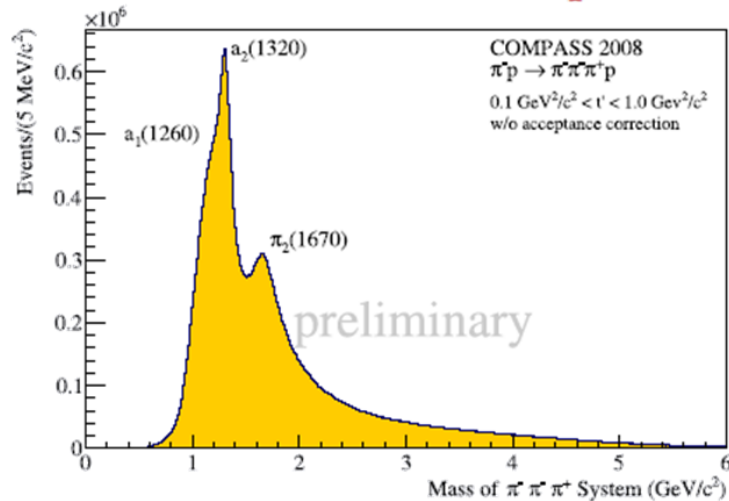
- **Analysis:**

- **Partial Wave Analysis (PWA) in mass bins with up to 53 waves**
- **fit of spin-density matrix for major waves with Breit-Wigner**

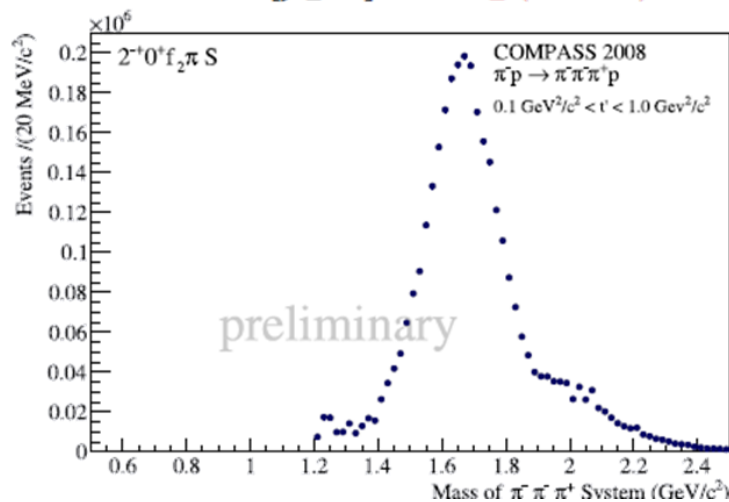
Hadron Spectroscopy

example: $\pi^- p \rightarrow \pi^- \pi^+ \pi^- p$ major waves

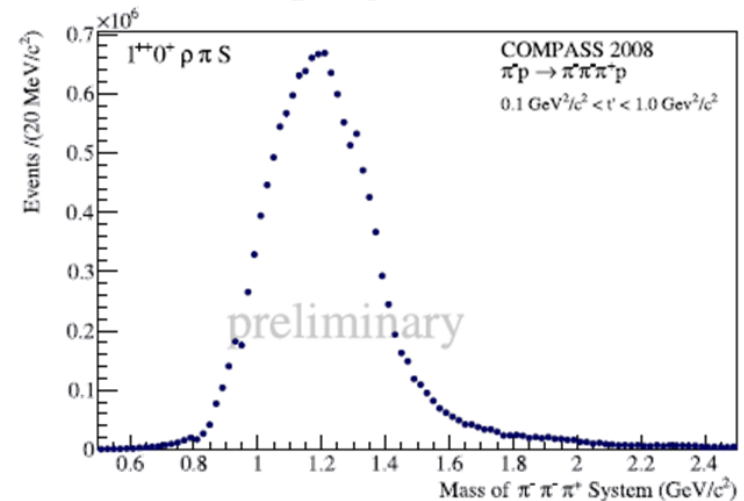
$\pi^- \pi^+ \pi^-$ invariant mass spectrum



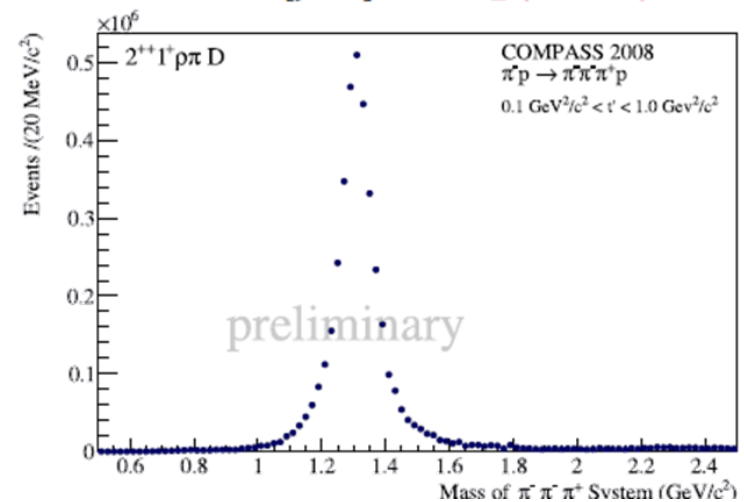
$2^-+ 0^+ [f_2 \pi] S : \pi_2(1670)$



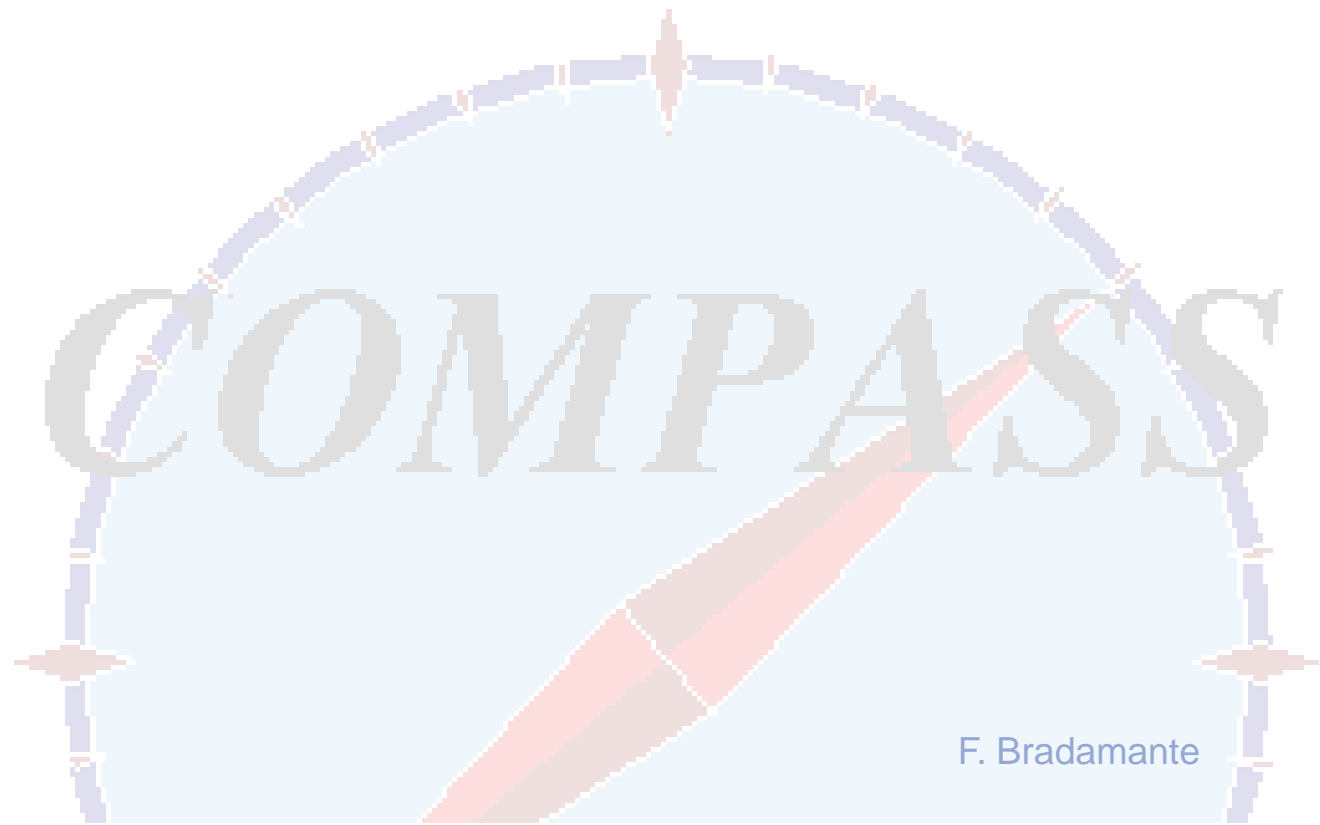
$1^{++} 0^+ [\rho \pi] S : a_1(1260)$



$2^{++} 1^+ [\rho \pi] D : a_2(1320)$



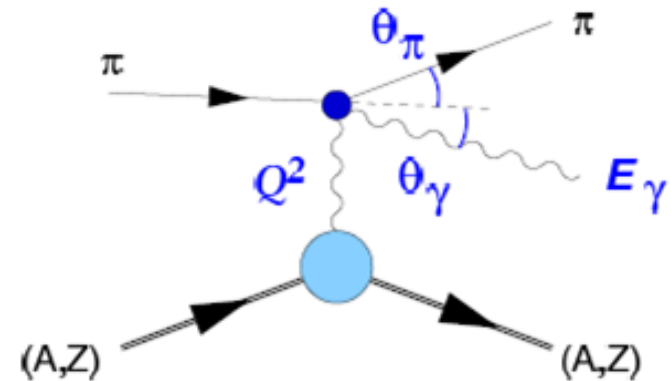
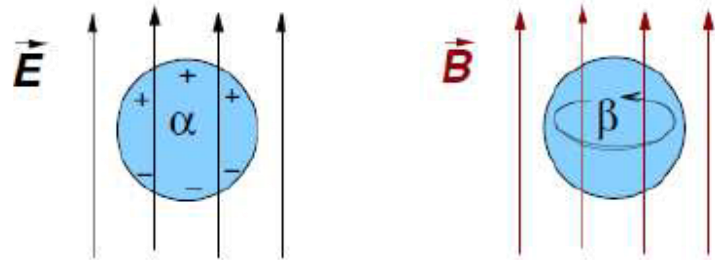
COMPASS hadron programme: Tests of χ PT



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COMPASS hadron programme: Tests of χ PT



pion (and kaon) polarisability via Primakoff scattering

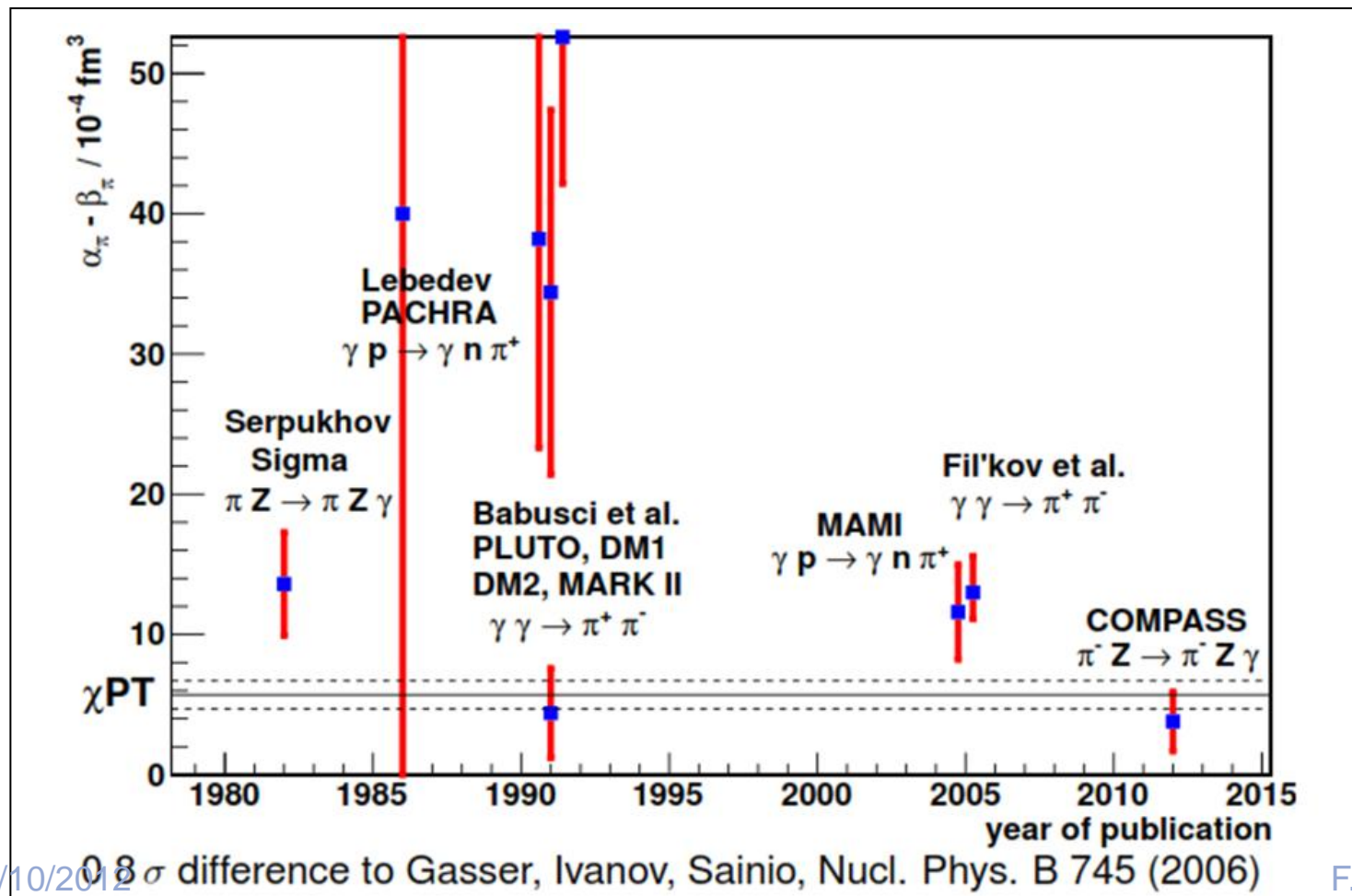
present exp. situation confused

	$\alpha_\pi - \beta_\pi$ (10^{-4} fm^3)	$\alpha_\pi + \beta_\pi$ (10^{-4} fm^3)	$\alpha_2 - \beta_2$ (10^{-4} fm^3)
2-loop ChPT prediction	5.7 ± 1.0	0.16 ± 0.10	16
COMPASS sensitivity	± 0.66	± 0.025	± 1.94

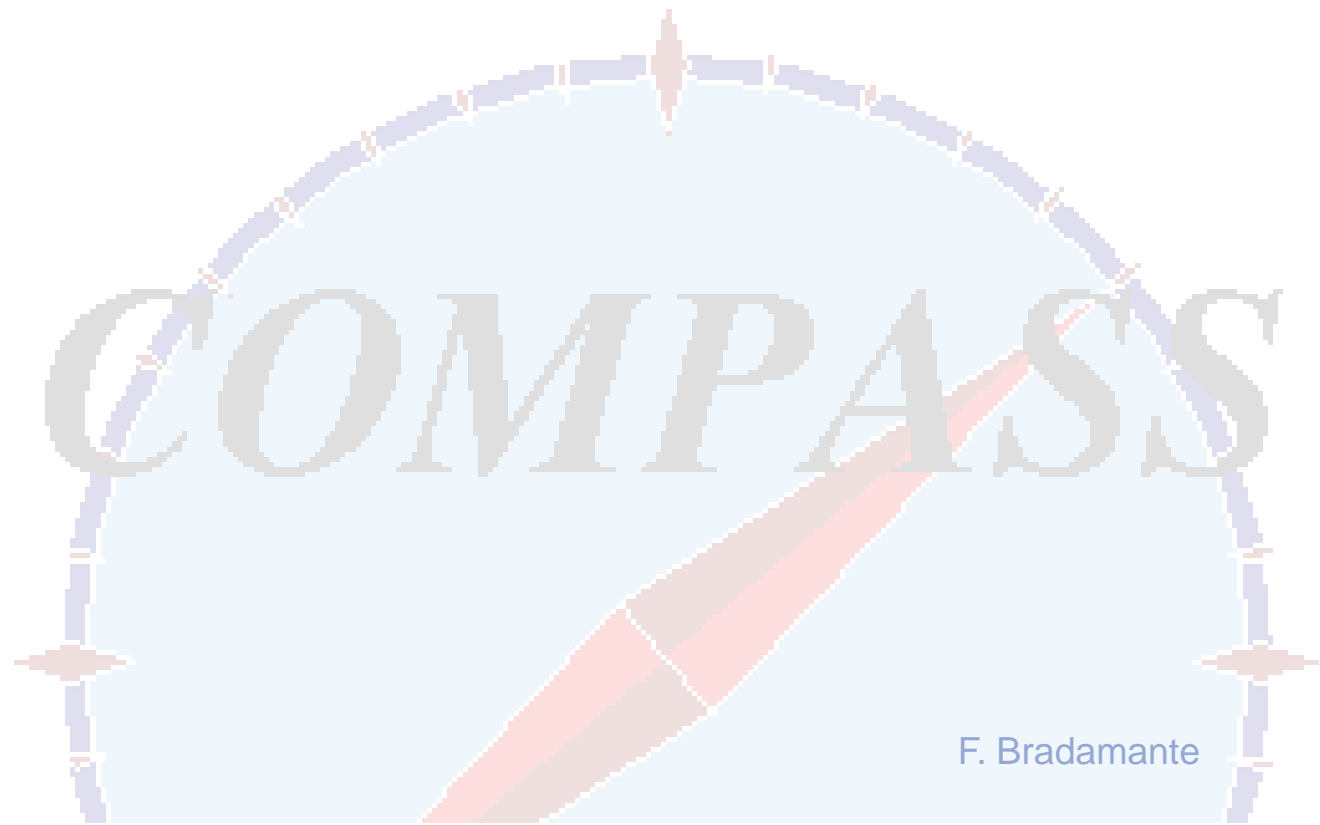
COMPASS hadron programme: Tests of χ PT

pion polarisability via Primakoff scattering
preliminary result, 2009 data

$$\alpha_\pi = 1.9 \pm 0.7_{stat} \pm 0.8_{syst} \cdot 10^{-4} fm^3$$

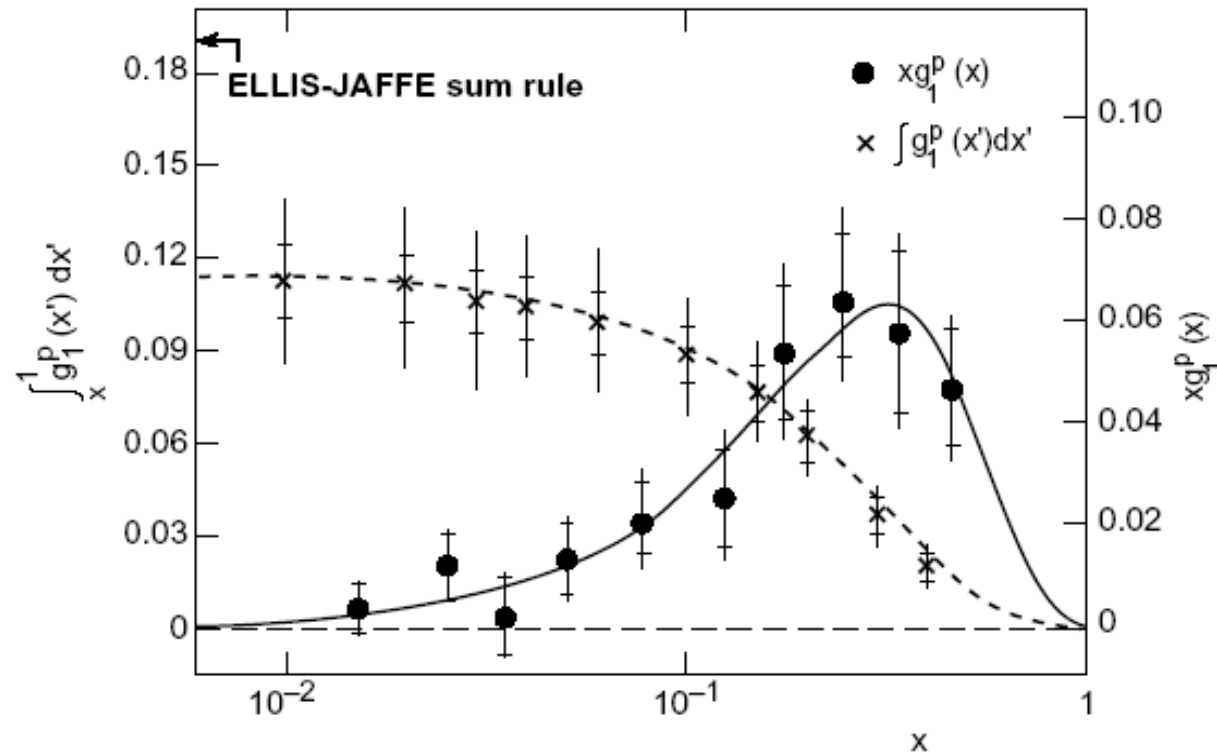


the **COMPASS** experiment and the **SPIN** Structure of the **NUCLEON**



The Quark Contribution to the Nucleon Spin

EMC
1988



$$\Gamma_1^p = 0.123 \pm 0.013 \pm 0.019$$

$$\Delta\Sigma = 0.12 \pm 0.17$$

→ SPIN CRISIS

The Quark Contribution to the Nucleon Spin

QUARK MODEL

$$|p\rangle = |uud\rangle$$

$$|n\rangle = |udd\rangle$$

SU(6)

magnetic moments:

$$\mu_p = \frac{4}{3}\mu_u - \frac{1}{3}\mu_d$$

$$\mu_n = \frac{4}{3}\mu_d - \frac{1}{3}\mu_u$$

assuming u and d Dirac particles with

$$m \cong \frac{1}{3}M_N$$

$$\mu_u = \frac{q\hbar}{2m_u c} = 2\mu_N \quad \mu_d = -\mu_N$$

$$\mu_p = 3\mu_N$$

$$\mu_n = -2\mu_N$$

$$\frac{\mu_p}{\mu_n} = -\frac{3}{2}$$

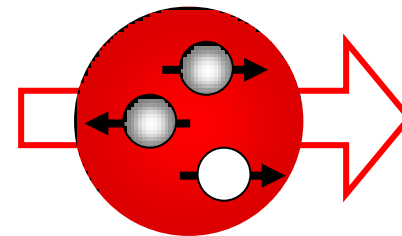
The Quark Contribution to the Nucleon Spin

in this model the **spin of the nucleon**
is given by the spin of the quarks

probability of finding a quark in a given state of polarization

$$\begin{array}{lll} \vec{\uparrow} & \vec{\uparrow} & \vec{\uparrow} \quad \vec{\downarrow} \\ \text{u} = \frac{5}{3} & \text{u} = \frac{1}{3} & \Delta\text{u} = \text{u} - \text{u} = \frac{4}{3} \\ \vec{\uparrow} & \vec{\downarrow} & \vec{\uparrow} \quad \vec{\downarrow} \\ \text{d} = \frac{1}{3} & \text{d} = \frac{2}{3} & \Delta\text{d} = \text{d} - \text{d} = -\frac{1}{3} \end{array}$$

$$\Delta\Sigma = \Delta\text{u} + \Delta\text{d} = 1$$



The Quark Contribution to the Nucleon Spin

CRITICAL REVIEW OF THE CONCEPTS

$$\frac{1}{2} = \frac{1}{2}\Delta\Sigma + \Delta G + L_z$$

NECESSITY

- to remeasure the proton
- to measure the neutron
- to measure ΔG

The Quark Contribution to the Nucleon Spin

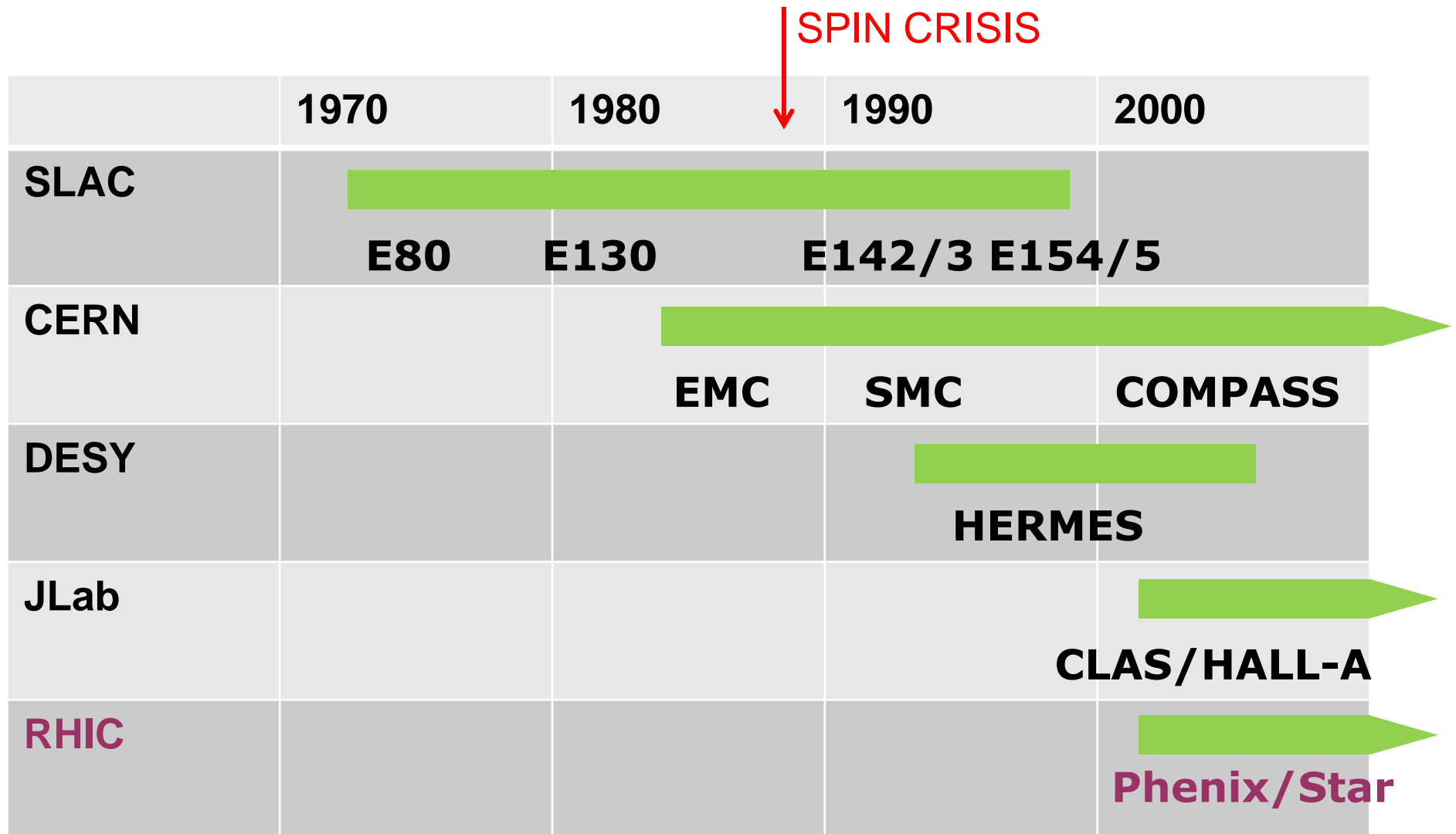
experiments

a worldwide effort since decades

The Quark Contribution to the Nucleon Spin

experiments

a worldwide effort since decades



The Quark Contribution to the Nucleon Spin

$$\Delta\Sigma = \Delta u + \Delta d + \Delta s \quad \text{Helicity PDFs}$$

$$\Delta q = \int_0^1 [\Delta q(x) + \Delta \bar{q}(x)] dx = \int_0^1 [q^+(x) - q^-(x) + \bar{q}^+(x) - \bar{q}^-(x)] dx$$

Δq 's can be extracted from the DIS cross-section asymmetry for parallel and antiparallel lepton and nucleon spins

Inclusive DIS, beam and target longitudinally polarized

$$\frac{d\Delta\sigma}{dx dy} = \lambda \cdot \frac{e^4}{4\pi^2 Q^2} \cdot \left[\left(1 - \frac{y}{2} - \frac{y^2}{4} \cdot \gamma^2 \right) \cdot \mathbf{g}_1 - \frac{y}{2} \cdot \gamma^2 \cdot \mathbf{g}_2 \right] \quad \begin{array}{l} d\sigma = d\bar{\sigma} \pm d\Delta\sigma \\ \uparrow \\ \text{beam/target helicity} \end{array}$$

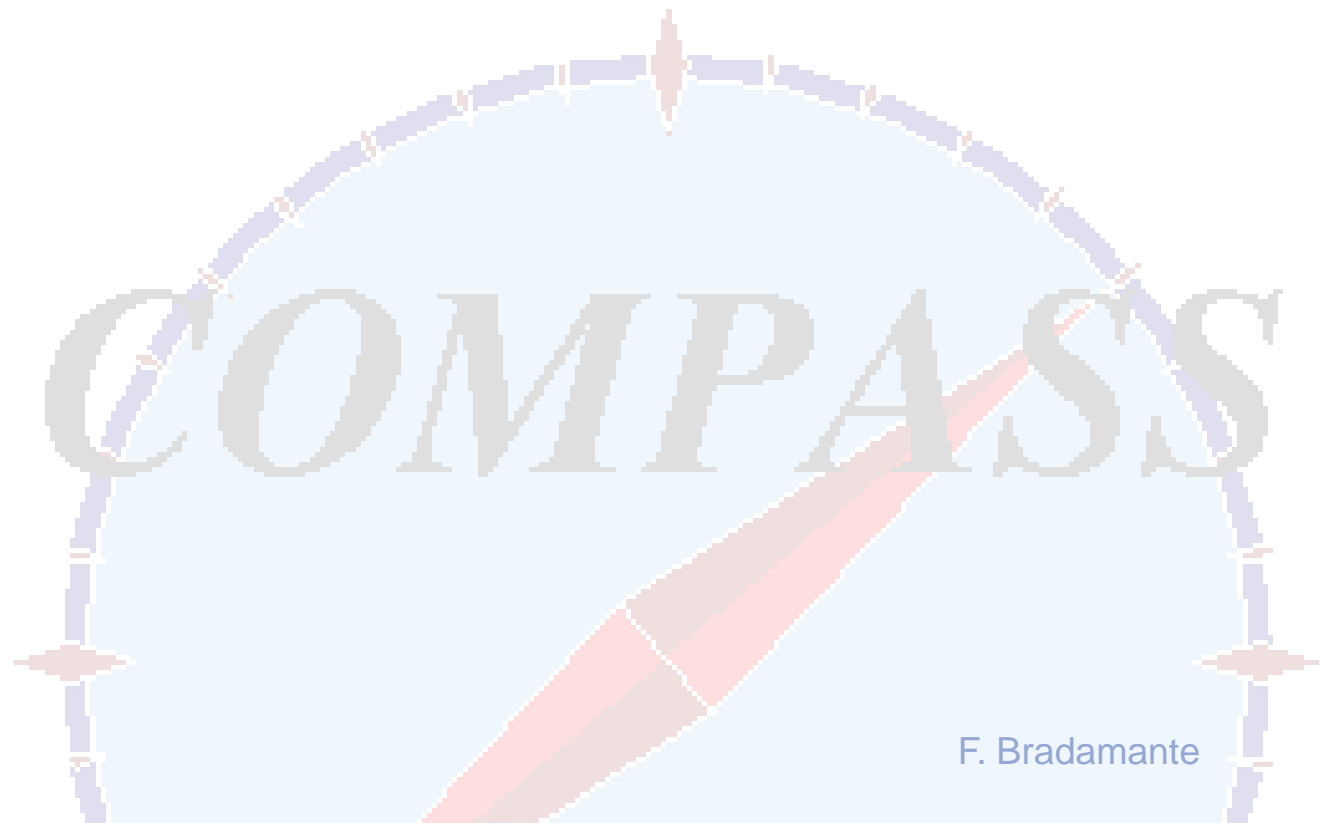
\mathbf{g}_1 measured at SLAC, EMC, SMC, HERMES, COMPASS

\mathbf{g}_2 suppressed by a factor $\gamma^2 \approx 0.01$ at 100 GeV (SMC, SLAC)

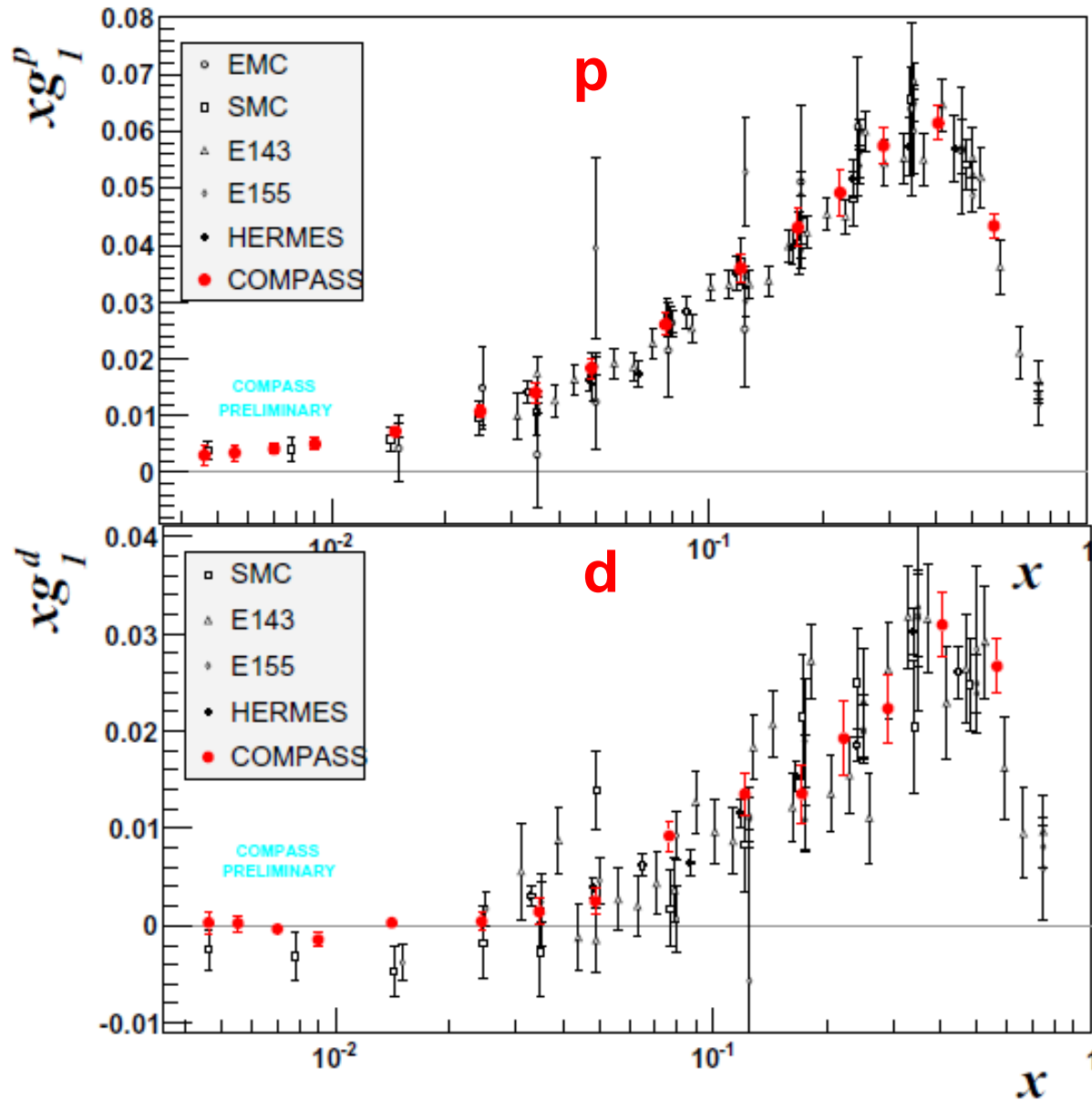
in the parton model

$$\mathbf{g}_1(x) = \frac{1}{2} \sum_q e_q^2 \cdot [\Delta q(x) + \Delta \bar{q}(x)]$$

The Structure Function $g_1(x, Q^2)$

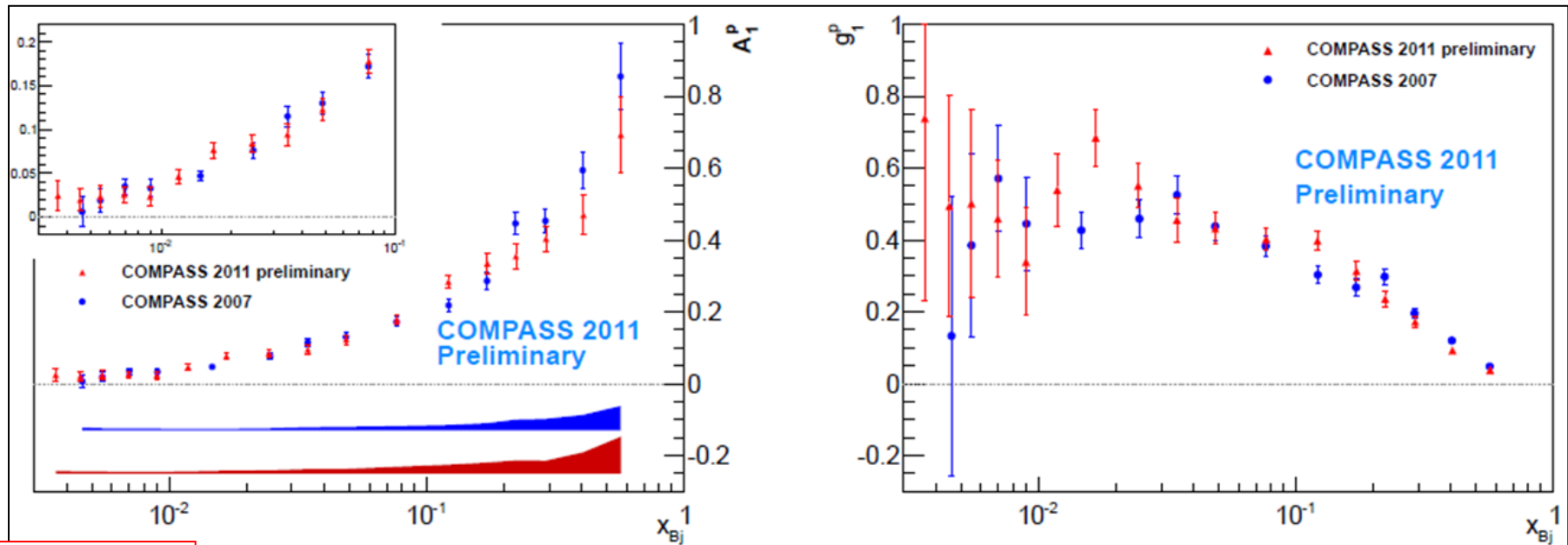


The Structure Function $g_1(x, Q^2)$



- very precise data
- only COMPASS for $x < 0.01$ ($Q^2 > 1$)
- deuteron data:
 $\Delta\Sigma = 0.33 \pm 0.03 \pm 0.05$
 $\Delta s + \Delta \bar{s} = -0.08 \pm 0.01 \pm 0.02$

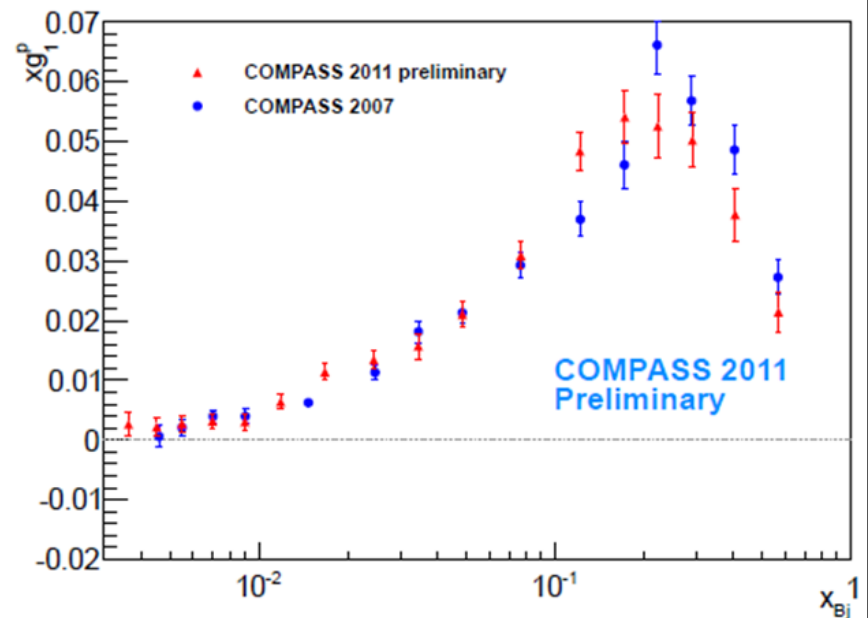
The Structure Function $g_1(x, Q^2)$



SPIN2012

$$g_1^p = \frac{F_2^p}{2x(1+R)} A_1^p$$

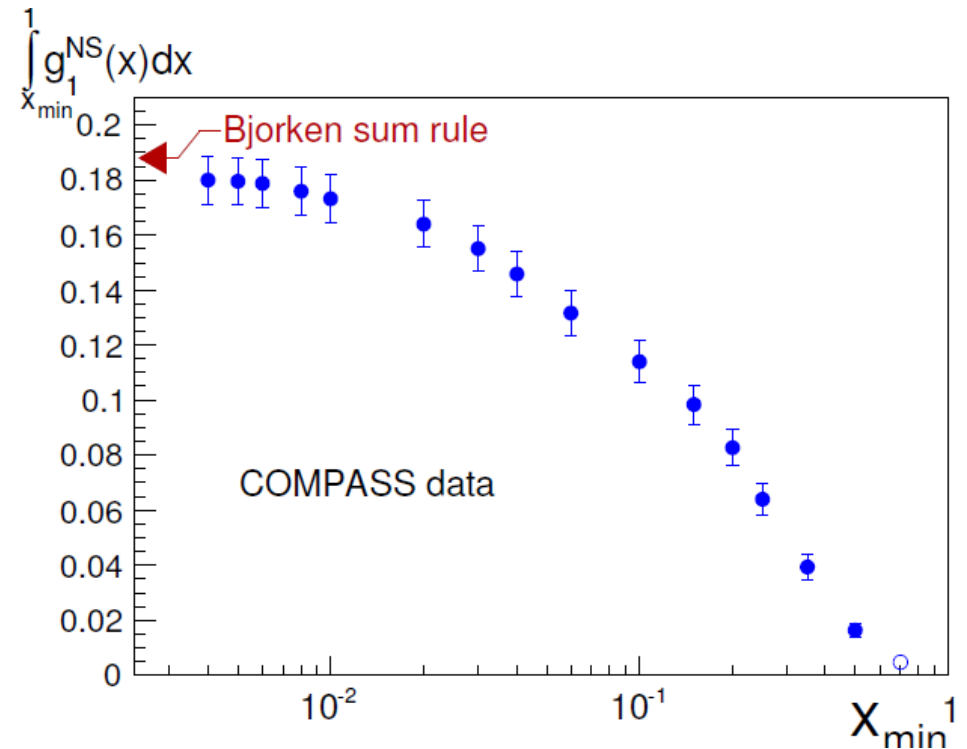
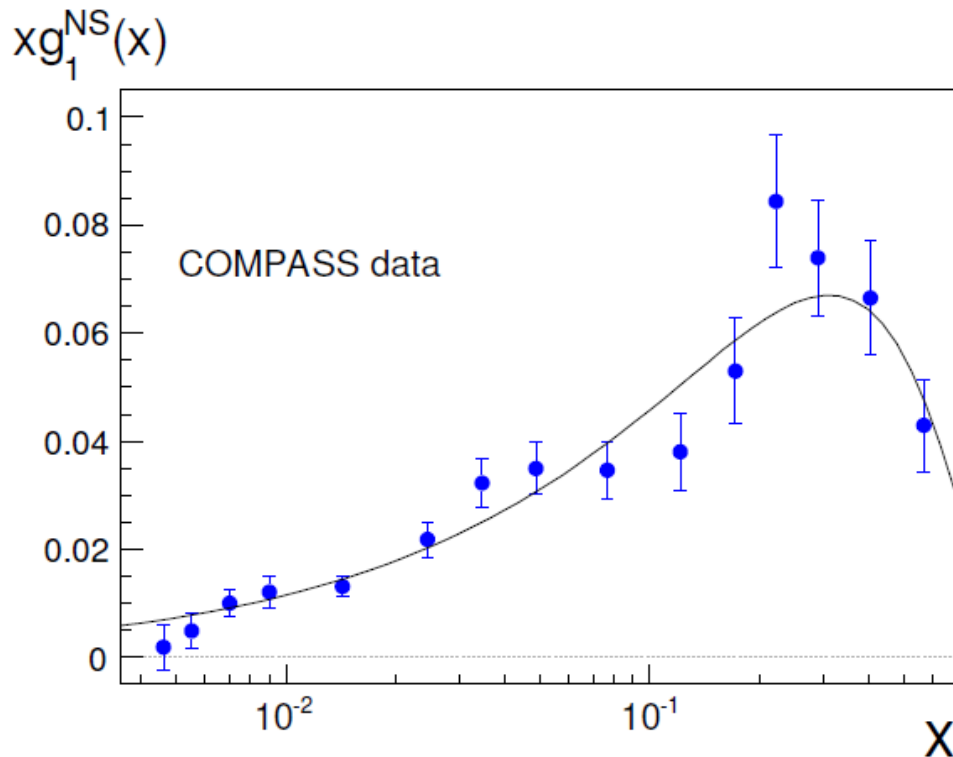
F_2^p from SMC parametrisation;
for R the same parametrisation as
in the evaluation of the depolarisation
factor D



The Structure Function $g_1(x, Q^2)$

Bjorken sum rule

$$g_1^{NS}(x, Q^2) = g_1^p(x, Q^2) - g_1^n(x, Q^2)$$



$$|g_A/g_V| = 1.28 \pm 0.07(\text{stat.}) \pm 0.10(\text{syst.})$$

PLB 690 (2010) 466

$$|g_A/g_V| = 1.269 \quad \text{from neutron } \beta \text{ decay}$$

The Gluon Contribution to the Nucleon Spin

Measurements of $\Delta G/G$

direct

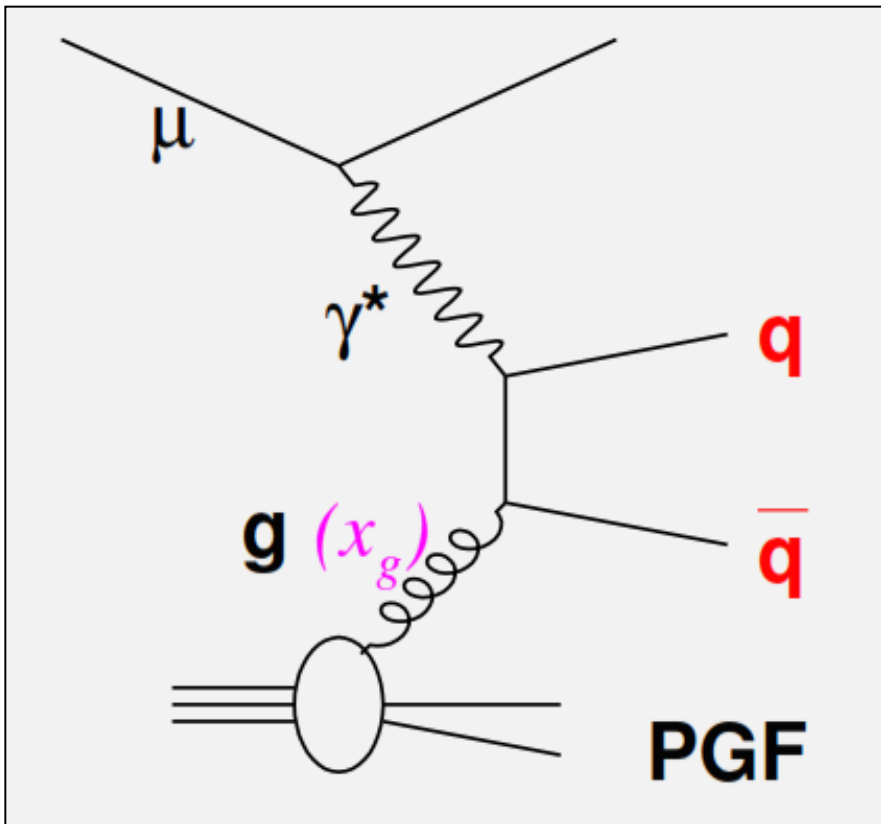
- cross-section asymmetry for charm production
- cross-section asymmetry for high p_T hadron pairs

indirect

- scale violation of $g_1(x, Q^2)$ (QCD fit)
SMC, SLAC Experiments, HERMES, COMPASS

Direct Measurements of $\Delta G/G$

Photon Gluon Fusion

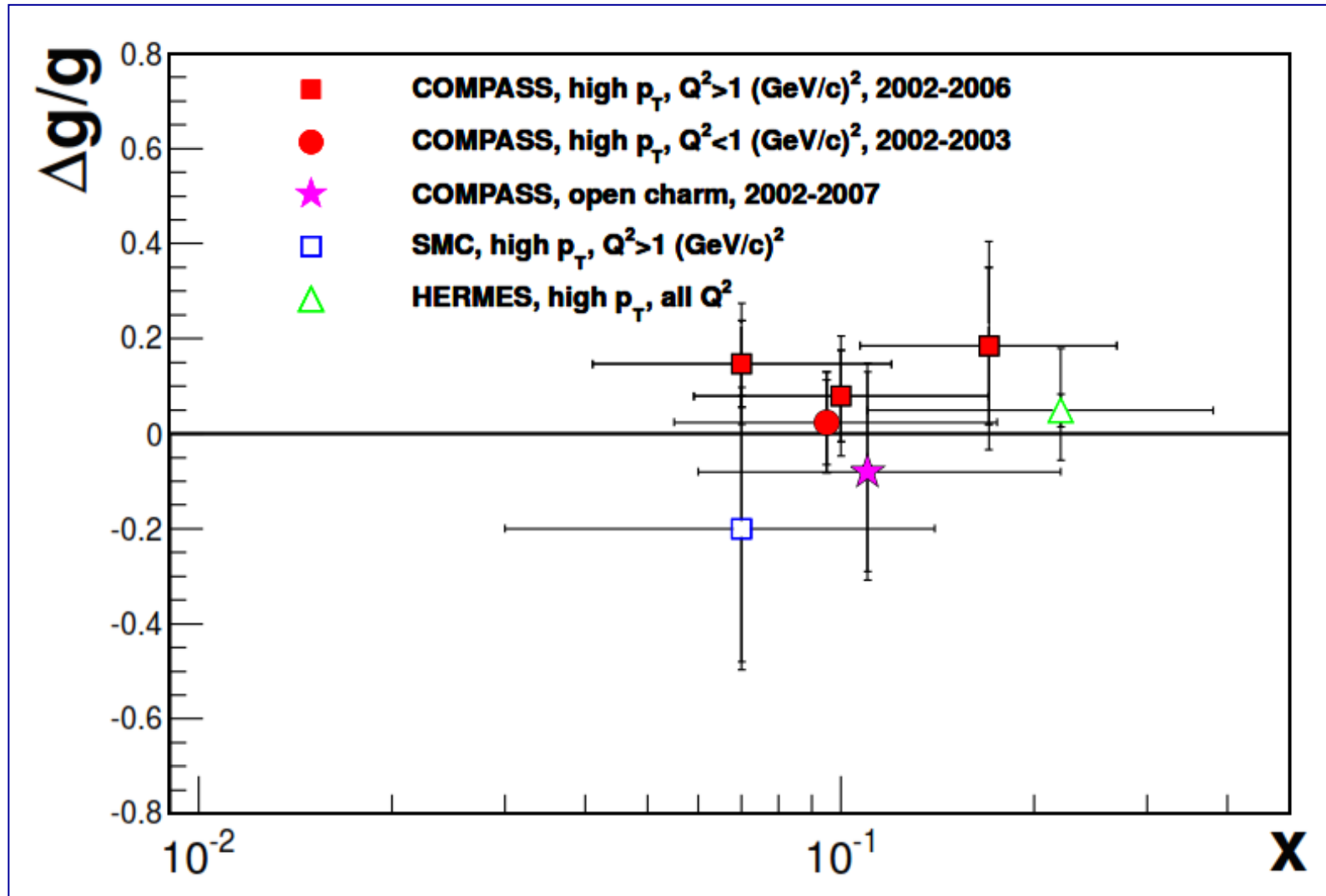


$q = c$ cross section difference
in **charmed meson production**
→ theory well understood
→ experiment challenging
COMPASS

$q = u, d, s$ cross section
difference in 2+1 jet
production: events with
hadrons with high- p_t
→ experiment easy
→ theory more difficult
COMPASS & HERMES

The Gluon Contribution to the Nucleon Spin

all LO results



THE SPIN PUZZLE

$$\frac{1}{2} = \frac{1}{2}\Delta\Sigma + \Delta G + L_z$$




$$\Delta\Sigma \sim 1/3$$

ΔG small

- necessity to address L_q and L_g → Ji sum rule
GPD
- **GROWING INTEREST IN
TRANSVERSE SPIN PHENOMENA**

The Structure of the Nucleon

three distribution functions are necessary to describe the quark structure of the nucleon at LO in the collinear case (Jaffe and Ji, 1991)

		nucleon polarisation		
		U	L	T
quark polarisation	U	f_1  number density q		
	L		g_1  -  helicity Δq	
	T			

The Structure of the Nucleon

three distribution functions are necessary to describe the quark structure of the nucleon at LO in the collinear case (Jaffe and Ji, 1991)






transversity PDF

correlation between the transverse polarisation of the nucleon and the transverse polarisation of the quark

chiral odd

can not be measured in inclusive DIS, still poorly known

quark polarisation

		nucleon polarisation		
		U	L	T
quark polarisation	U	f_1  number density q		
	L		g_1  -  helicity Δq	
	T			h_1  -  transversity $\Delta_T q$

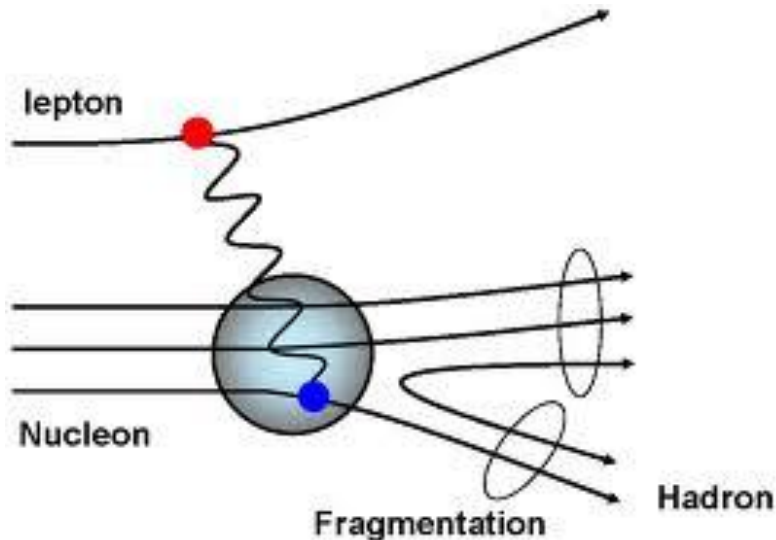
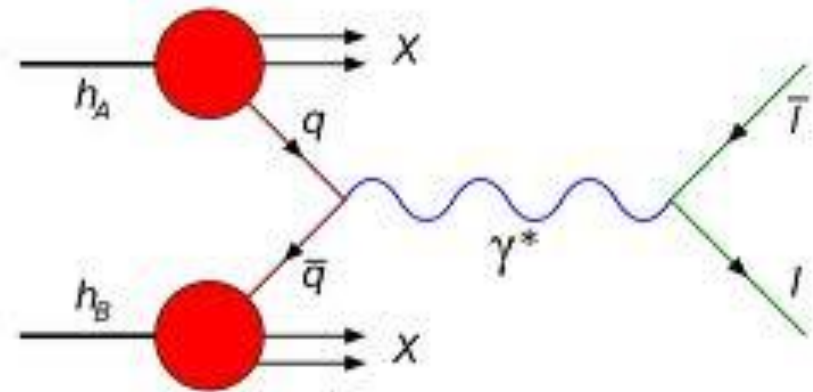
Probing Transversity

transversely polarized DY

(as proposed by Ralston & Soper)

A_{TT} in pp DY very small

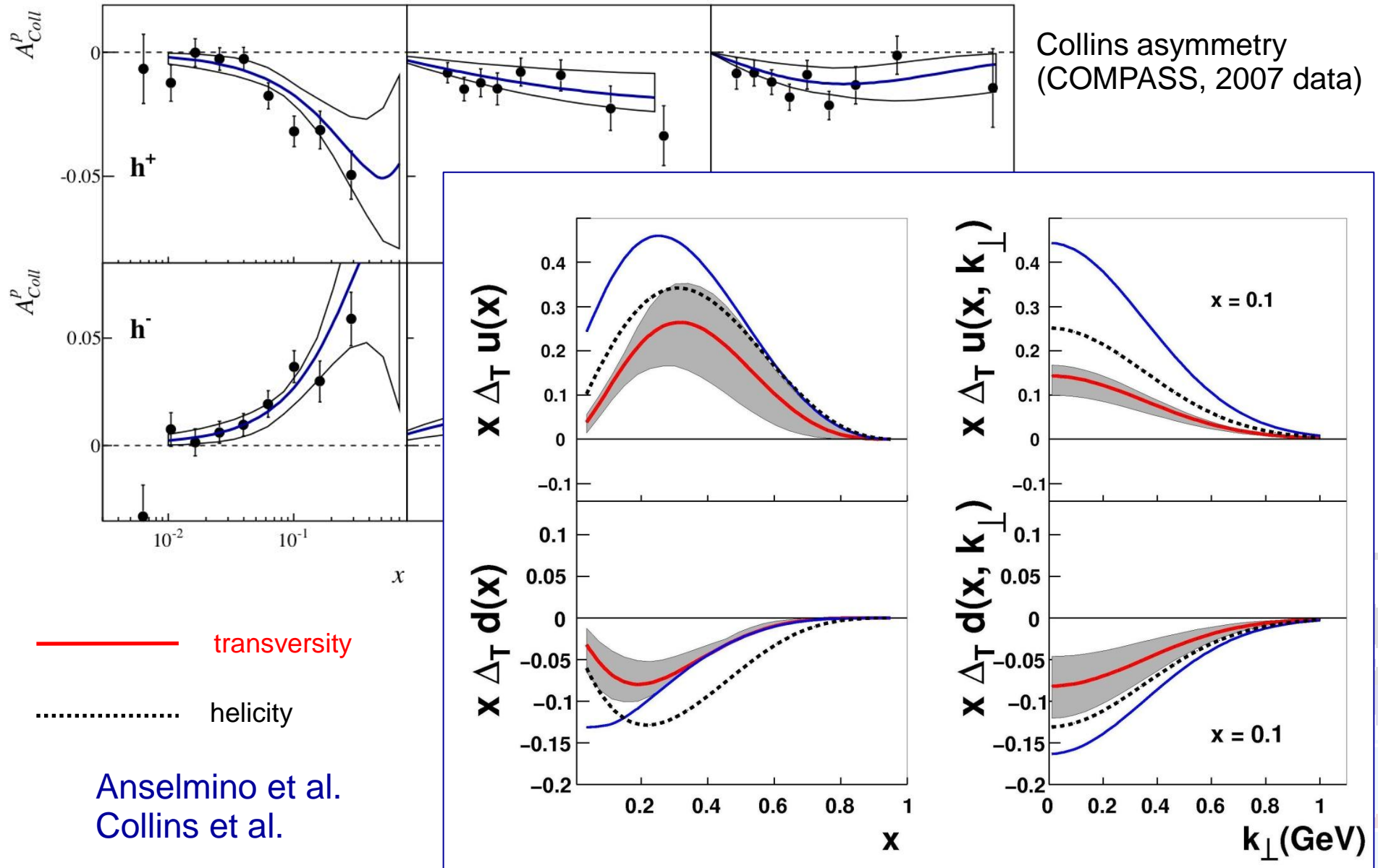
dream option: measure ppbar DY



Semi-Inclusive Deep Inelastic Scattering

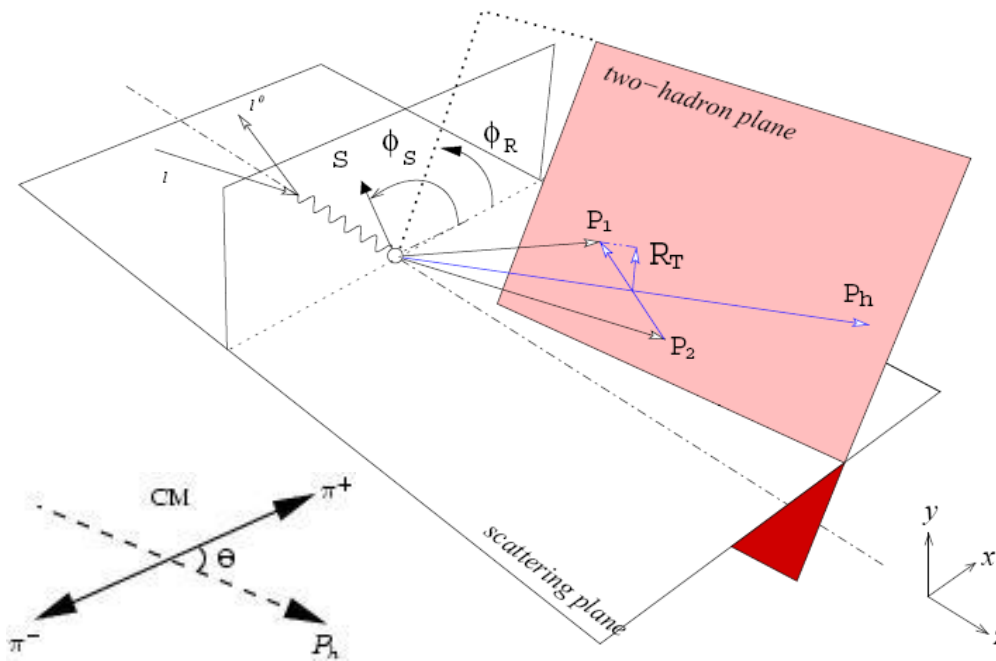
Collins asymmetry: $\sin(\phi_h + \phi_s)$
correlation between spin of quark and
transverse momentum of hadrons
convolution of Transversity and Collins FF
**measurable as an azimuthal
modulation of the hadrons**

Transversity distribution extracted from SIDIS



Alternative way to access Transversity in SIDIS

two-hadron asymmetry



azimuthal asymmetry in

$$\phi_{RS} = \phi_{R\perp} + \phi_S$$

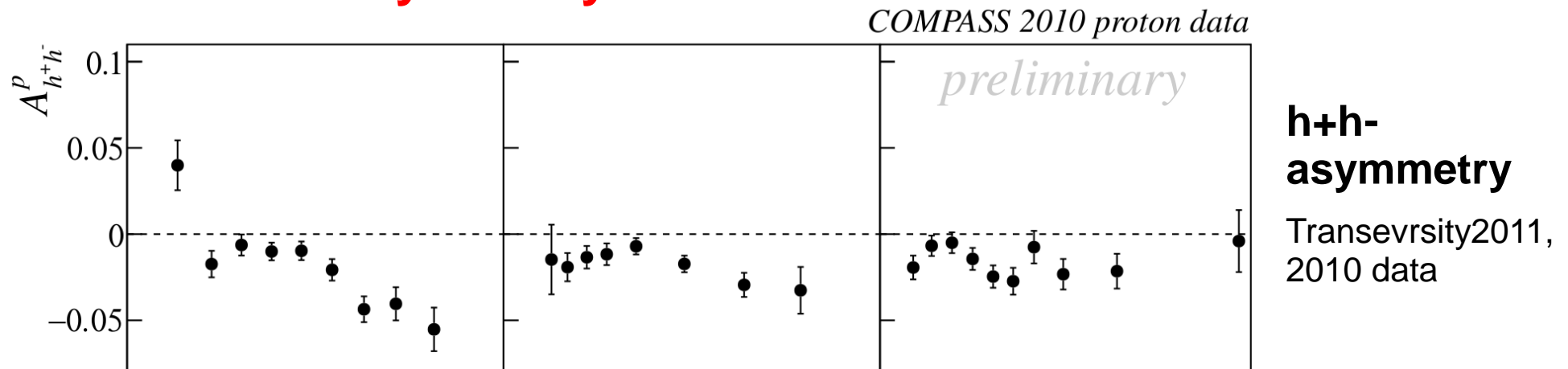
$\phi_{R\perp}$ is the azimuthal angle of the plane defined by the two hadrons

product of Transversity and Interference FF

A. Bacchetta, M. Radici, hep-ph/0407345
X. Artru, hep-ph/0207309

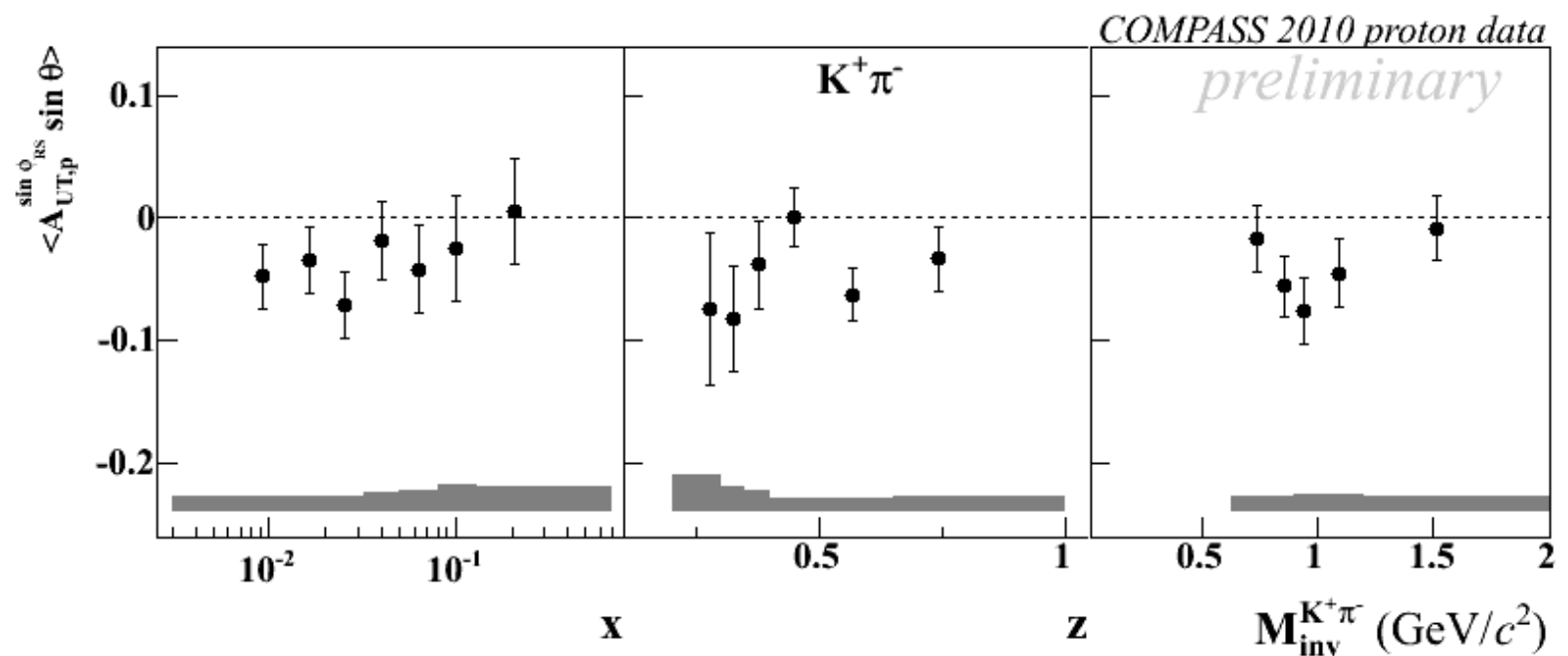
Alternative way to access Transversity in SIDIS

two-hadron asymmetry



K+ π - asymmetry

SPIN2012,
2010 data

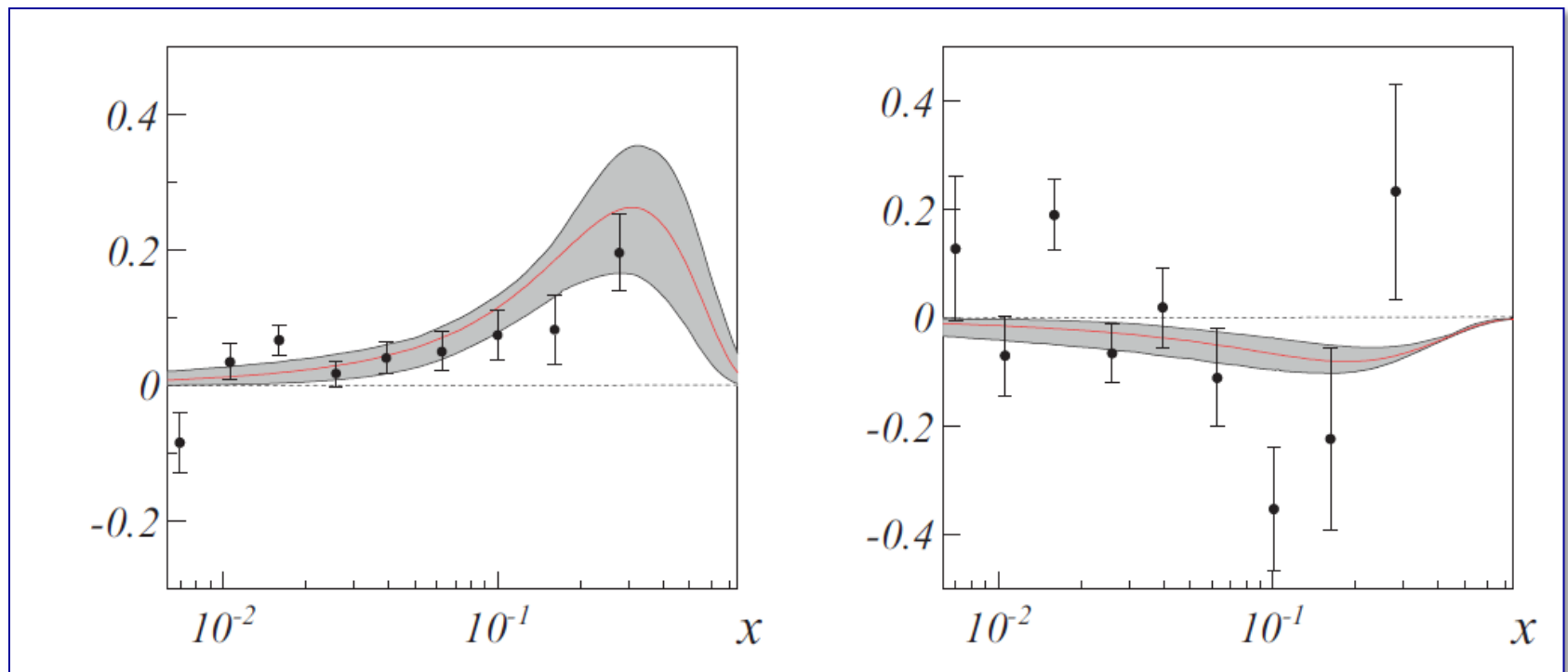


Transversity distribution extracted from two-hadron asymmetries in SIDIS

following the work of A Bacchetta, A. Courtoy, M Radici
Phys.Rev.Lett.107:012001,2011

first direct extraction of u and d transversity
using deuteron 2002-2004 and proton 2010 COMPASS results

C.Elia Ph.D Thesys, Trieste 2012



curves: global fits of COMPASS d, HERMES p, BELLE e+e- data
(M. Anselmino et al. PRD 75 (2007) 054032)

... BUT ...

... going beyond collinear pQCD...

- **There are some phenomena involving transverse momenta (and transverse spin) which are not accounted for by a collinear pQCD description**
- **When the observed transverse momentum P_T is much smaller than the hard scale Q (two-scale process), one has to introduce the **transverse-momentum dependent** (or “**unintegrated**”) **distributions (TMDs)****
- **The TMD physics was prompted by the study of transverse spin phenomena but has also led to an improved QCD knowledge of ordinary, unpolarized, TMDs**

... going beyond collinear pQCD...

transverse spin couples to transverse momentum
giving rise to a number of possible correlations

Sivers function:

unpolarized quarks in a transversely polarized nucleon

$$f_{1T}^\perp = \begin{array}{c} \uparrow \\ \circ \\ \bullet \end{array} - \begin{array}{c} \uparrow \\ \bullet \\ \circ \end{array} \quad f_{q/p^\dagger}(x, \mathbf{p}_T) = f_1^q(x, p_T^2) - f_{1T}^{\perp q}(x, p_T^2) \frac{(\hat{\mathbf{P}} \times \mathbf{p}_T) \cdot \mathbf{S}}{M}$$

→ azimuthal single-spin $\sin(\phi_h - \phi_S)$ asymmetry

Boer-Mulders function:

transversely polarized quarks in an unpolarized nucleon

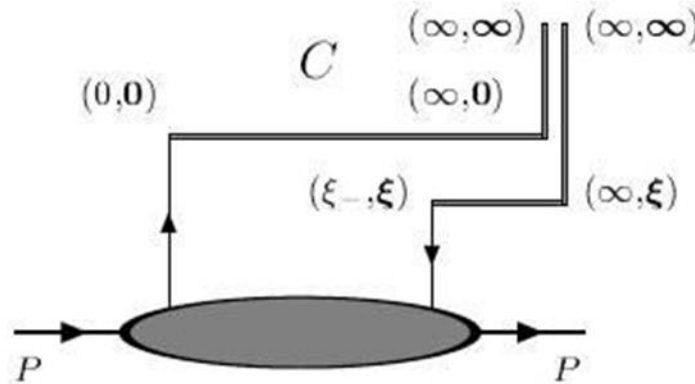
$$h_1^\perp = \begin{array}{c} \uparrow \\ \bullet \\ \circ \end{array} - \begin{array}{c} \bullet \\ \uparrow \\ \circ \end{array} \quad f_{q^\dagger/p}(x, \mathbf{p}_T) = \frac{1}{2} \left[f_1^q(x, p_T^2) - h_1^{\perp q}(x, p_T^2) \frac{(\hat{\mathbf{P}} \times \mathbf{p}_T) \cdot \mathbf{S}_q}{M} \right]$$

→ azimuthal $\cos \phi_h$ and $\cos 2\phi_h$ asymmetries

... going beyond collinear pQCD...

the gauge structure of TMDs

$$(x, \mathbf{k}_T) = \int \frac{d\xi^-}{2\pi} \int \frac{d^2\xi_T}{(2\pi)^2} e^{ixP^+ \xi^-} e^{-i\mathbf{k}_T \cdot \xi_T} \langle P, S | \bar{\psi}(0) \mathcal{W}[0, \xi] \psi(\xi) | P, S \rangle |_{\xi^+ = 0}$$



$$\text{SIDIS : } \mathcal{W}[0, \xi] = \mathcal{W}^- [0, \infty] \mathcal{W}^T [0_T, \infty_T] \mathcal{W}^T [\infty_T, \xi_T] \mathcal{W}^- [\infty, \xi]$$

the existence of Sivers and Boer-Mulders functions is a consequence of the gauge link structure, which also implies

$$\text{TMD (SIDIS)} = - \text{TMD (DY)}$$

this is a fundamental test of gauge invariance

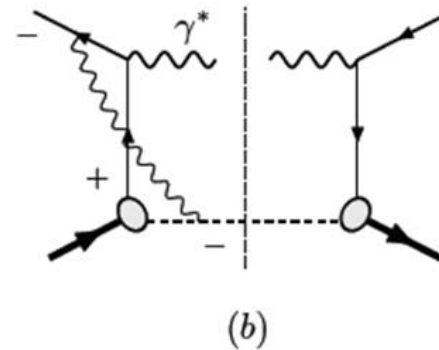
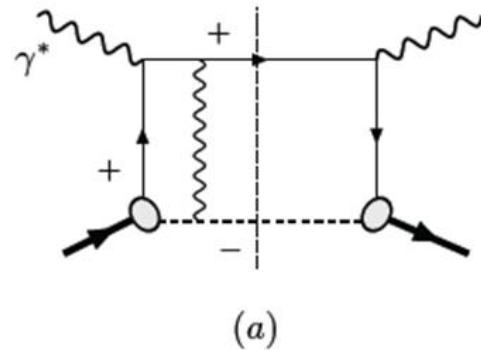
... going beyond collinear pQCD...

Crucial role of gauge-links in TMDs

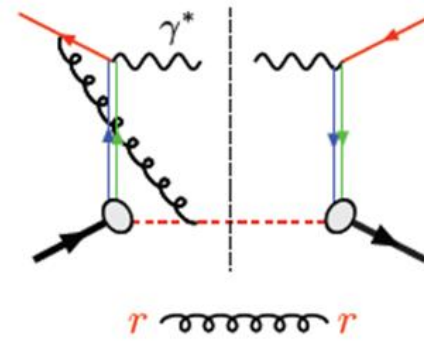
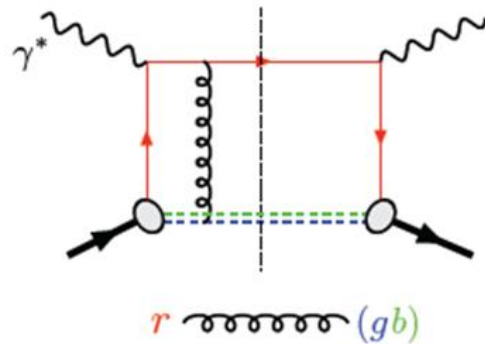
Brodsky, Hwang, Schmidt;
Collins; Belitsky, Ji, Yuan;
Boer, Mulders, Pijlman

process-dependence of Sivers functions

DIS:
"attractive"


















D-Y:
"repulsive"



$$[f_{1T}^{q\perp}]_{\text{SIDIS}} = -[f_{1T}^{q\perp}]_{\text{DY}}$$

The Structure of the Nucleon

taking into account the quark intrinsic transverse momentum k_T ,
 at leading order other 6 TMD PDFs are needed for a full description
 of the nucleon quark structure

		nucleon polarisation			
		U	L	T	
quark polarisation	U	f_1  number density q		f_{1T}^\perp  -  Sivers	$\Delta_0^T q$
	L		g_1  -  helicity Δq	g_{1T}  - 	
	T	h_1^\perp  -  Boer Mulders	h_{1L}^\perp  - 	h_1  -  transversity h_{1T}^\perp  - 	$\Delta_T q$

The Structure of the Nucleon

taking into account the quark intrinsic transverse momentum k_T ,
 at leading order other 6 TMD PDFs are needed for a full description
 of the nucleon quark structure

Sivers function

correlation between the
 transverse spin of the nucleon
 and the transverse momentum
 of the quark








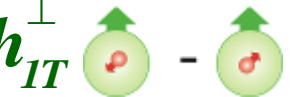
*sensitive to orbital angular
 momentum*

Boer-Mulders function

correlation between the
 transverse spin and the
 transverse momentum
 of the quark in unpol nucleons

T-odd

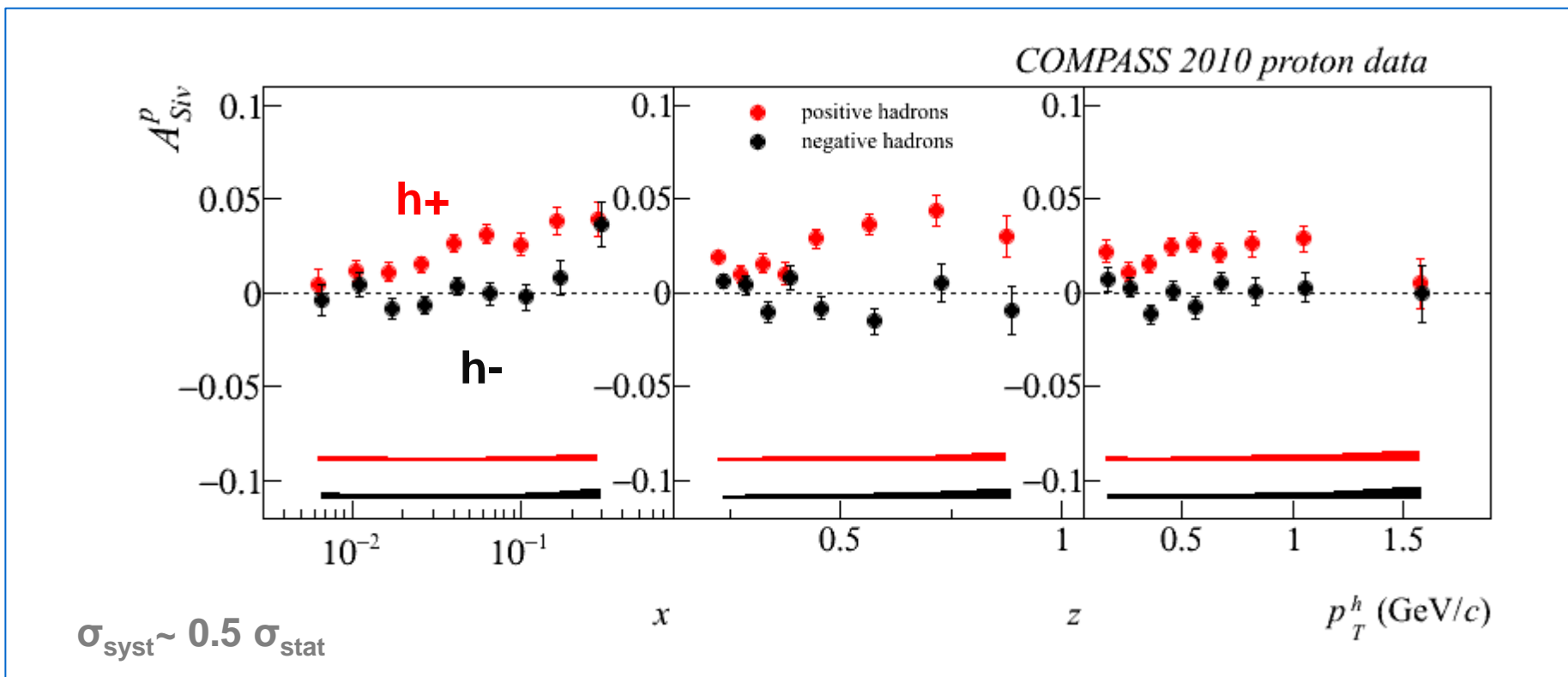
quark
 polarisation

		nucleon polarisation			
		U	L	T	
U	f_1 <i>number density</i> q 	f_{1T}^\perp <i>Sivers</i> 	$\Delta_0^T q$		
L		g_{1T} 			
T	h_1^\perp <i>Boer Mulders</i> 	h_{1L}^\perp 	h_1  <i>transversity</i> h_{1T}^\perp 	$\Delta_T q$	

SIDIS gives access to all of them

Sivers asymmetry on **proton**

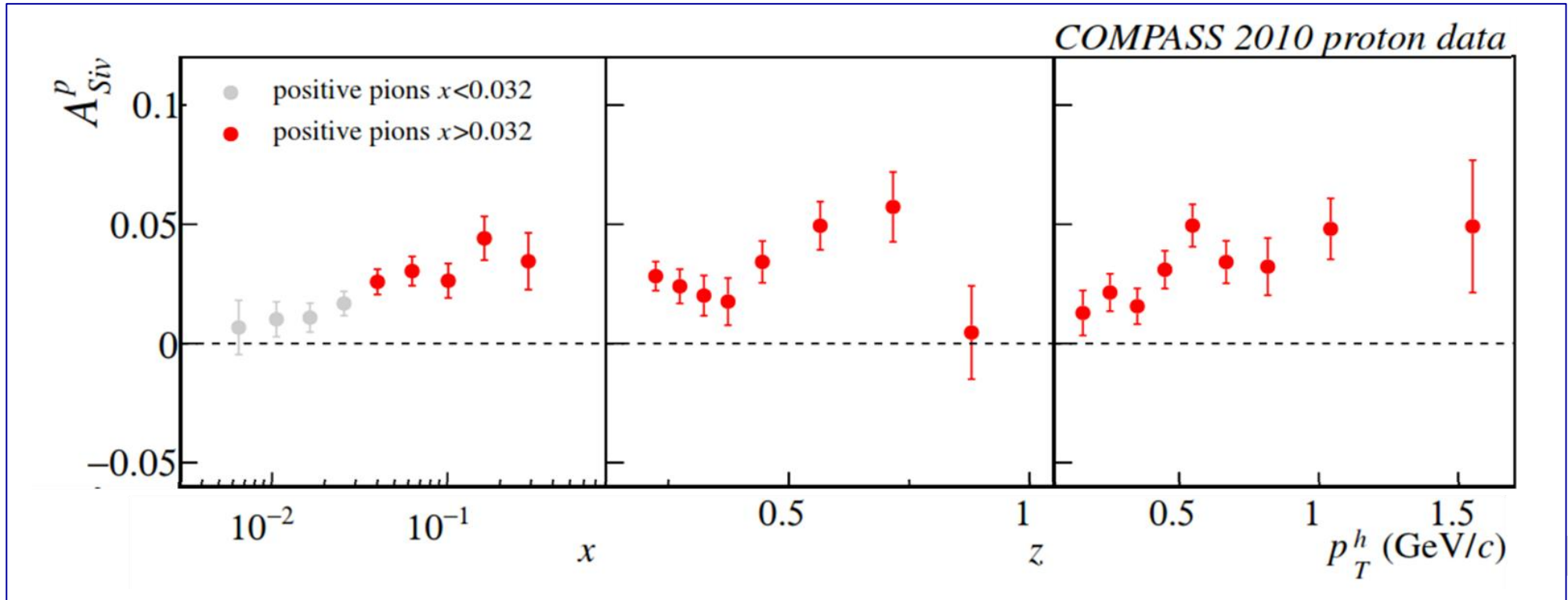
charged hadrons, 2010 data PLB 717 (2012) 383



clear positive signal for h^+ , which extends to small x
EVIDENCE FOR ORBITAL ANGULAR MOMENTUM OF QUARKS

Sivers asymmetry on **proton**

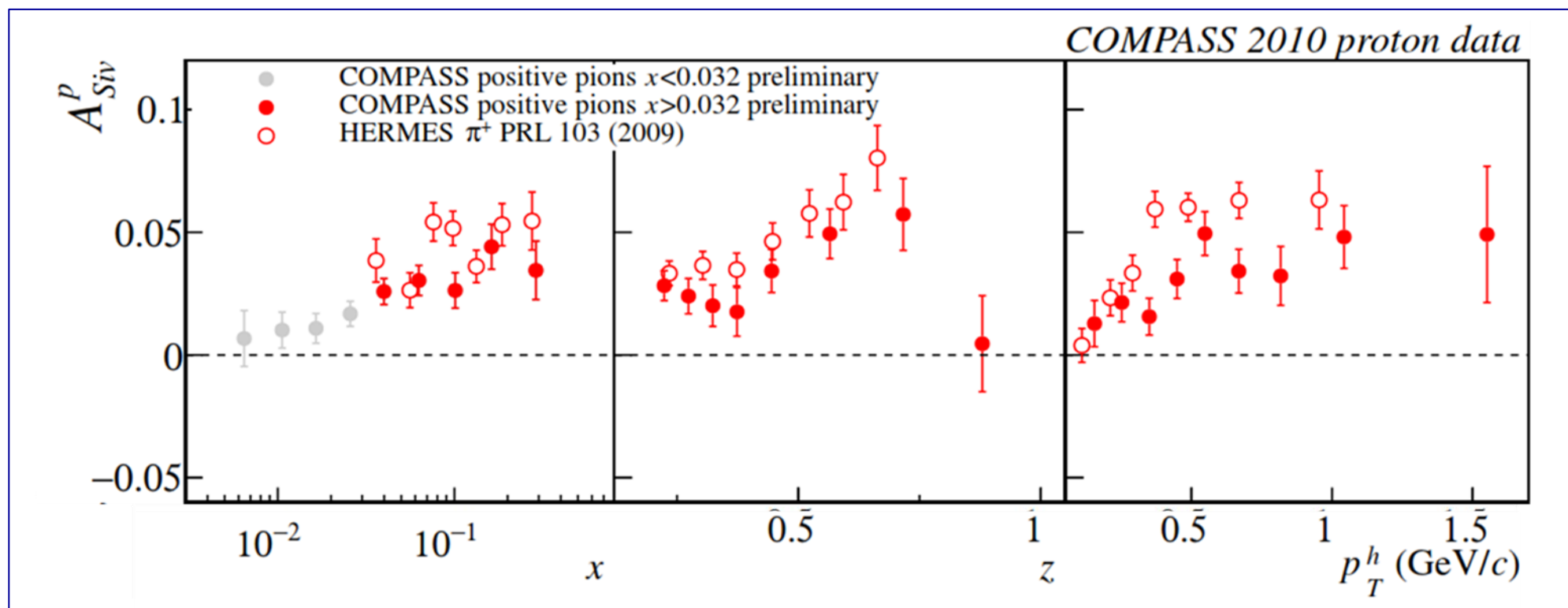
charged pions, 2010 data SPIN2012



comparison with HERMES

Sivers asymmetry on **proton**

charged pions, 2010 data SPIN2012



comparison with HERMES

... TMD Q2 evolution: the very new developments !

CONCLUSIONS

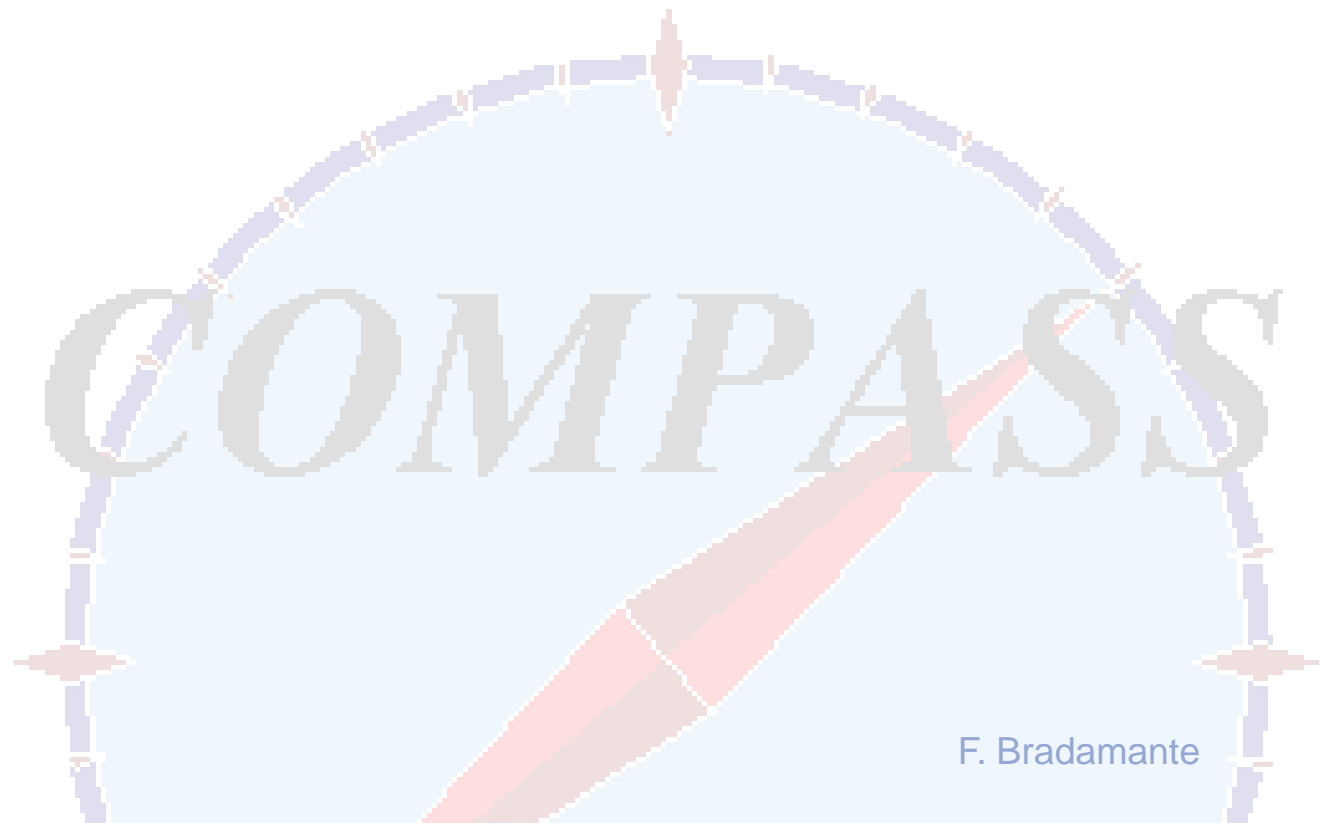
- the quarks contribute 1/3 of the nucleon spin
- the relative size of DG and L not clearly assessed
- COMPASS II will address L via DVCS

- **TRANSVERSE SPIN AND MOMENTUM PHENOMENA:**
NEW Properties of matter have been unveiled
 Transversity and Collins effect, Sivers effect
OTHER correlations are still possible (Boer-Mulders)

- COMPASS II should be able to
 - test the predicted sign change of the Sivers function from SIDIS to D-Y
 - perform precise measurements of unpolarised SIDIS azimuthal asymmetries on p

- COMPASS has a great discovery potential for glueballs and hybrids

THANK YOU !



future measurements

Table 2: Summary of the different physics items for the far and near future. Already approved measurements are in bold.

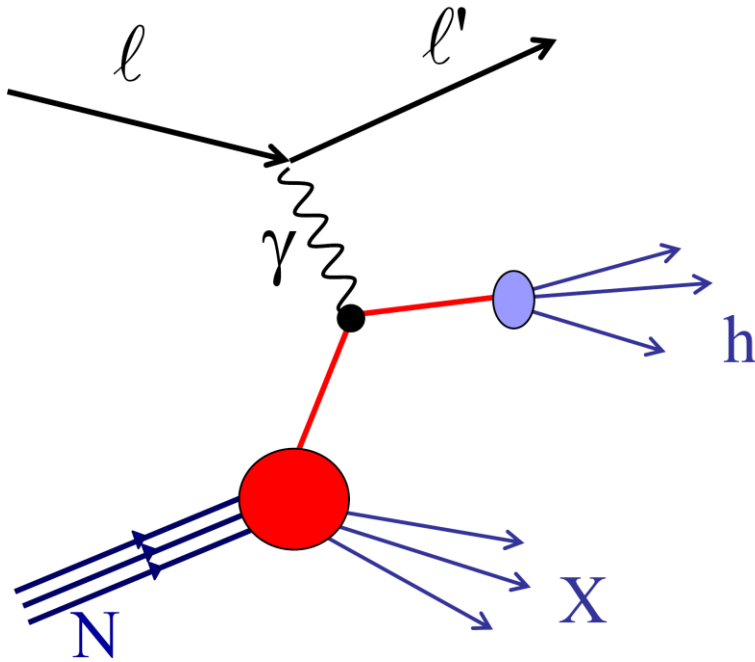
	physics item	key aspects of the measurement
GPD	H <i>t</i>-slope parameter B E	RPD, Beam Charge and Spin Asymmetries $d\sigma/dt$ transversely polarized proton target
SIDIS	hadron multiplicities for π and K $h_{1,u}^\perp, h_{1,d}^\perp$ h_1^d with same accuracy as h_1^u f_1^\perp evolution	PID and absolute acceptance azimuthal modulations and PID transversely polarized deuteron target 100 GeV and transversely polarized proton target
DY	sign change for f_1^\perp and h_1^\perp universality of TMD PDFs flavor separation test of the Lam-Tung relation EMC effect in DY	transversely polarized proton target higher statistics with transversely polarized proton target transversely polarized deuteron target hydrogen target different nuclear targets

P. Newman in Cracow

COMPASS planned measurements in the next five years and longer term perspectives on the study of the nucleon structure

Semi-Inclusive Deep Inelastic Scattering

hard interaction of a lepton with a nucleon via virtual photon exchange

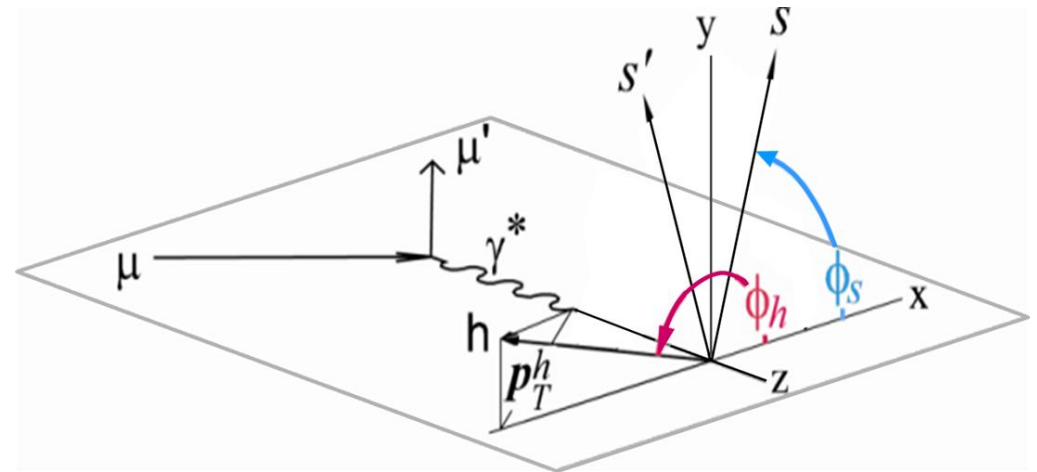


$$x = \frac{Q^2}{2P \cdot q} \quad y = \frac{P \cdot q}{P \cdot \ell} =_{LAB} \frac{E - E'}{E}$$

$$Q^2 = -q^2 \quad W^2 = (P + q)^2$$

$$z = \frac{P \cdot P_h}{P \cdot q} =_{LAB} \frac{E_h}{E - E'}$$

$$\sigma^{\ell N \rightarrow \ell h X} \propto \sum_q q(x) \otimes \sigma^{\ell q \rightarrow \ell q} \otimes D_q^h(z)$$



Semi-Inclusive Deep Inelastic Scattering

$$\begin{aligned}
 \frac{d\sigma}{dx dy d\psi dz d\phi_h dP_{h\perp}^2} = & \frac{\alpha^2}{xy Q^2} \frac{y^2}{2(1-\varepsilon)} \left(1 + \frac{\gamma^2}{2x}\right) \left\{ \begin{aligned} & F_{UU,T} + \varepsilon F_{UU,L} + \sqrt{2\varepsilon(1+\varepsilon)} \cos\phi_h F_{UU}^{\cos\phi_h} \\ & + \varepsilon \cos(2\phi_h) F_{UU}^{\cos 2\phi_h} + \lambda_e \sqrt{2\varepsilon(1-\varepsilon)} \sin\phi_h F_{LU}^{\sin\phi_h} \end{aligned} \right. \quad \text{unpol target} \\
 & + S_{\parallel} \left[\sqrt{2\varepsilon(1+\varepsilon)} \sin\phi_h F_{UL}^{\sin\phi_h} + \varepsilon \sin(2\phi_h) F_{UL}^{\sin 2\phi_h} \right] + S_{\parallel} \lambda_e \left[\sqrt{1-\varepsilon^2} F_{LL} + \sqrt{2\varepsilon(1-\varepsilon)} \cos\phi_h F_{LL}^{\cos\phi_h} \right] \quad \rightarrow \text{pol target} \\
 & + |S_{\perp}| \left[\begin{aligned} & \sin(\phi_h - \phi_S) \left(F_{UT,T}^{\sin(\phi_h - \phi_S)} + \varepsilon F_{UT,L}^{\sin(\phi_h - \phi_S)} \right) \\ & + \varepsilon \sin(\phi_h + \phi_S) F_{UT}^{\sin(\phi_h + \phi_S)} + \varepsilon \sin(3\phi_h - \phi_S) F_{UT}^{\sin(3\phi_h - \phi_S)} \\ & + \sqrt{2\varepsilon(1+\varepsilon)} \sin\phi_S F_{UT}^{\sin\phi_S} + \sqrt{2\varepsilon(1+\varepsilon)} \sin(2\phi_h - \phi_S) F_{UT}^{\sin(2\phi_h - \phi_S)} \end{aligned} \right] \quad \uparrow \text{pol target} \\
 & + |S_{\perp}| \lambda_e \left[\begin{aligned} & \sqrt{1-\varepsilon^2} \cos(\phi_h - \phi_S) F_{LT}^{\cos(\phi_h - \phi_S)} + \sqrt{2\varepsilon(1-\varepsilon)} \cos\phi_S F_{LT}^{\cos\phi_S} \\ & + \sqrt{2\varepsilon(1-\varepsilon)} \cos(2\phi_h - \phi_S) F_{LT}^{\cos(2\phi_h - \phi_S)} \end{aligned} \right] \left. \right\}, \\
 \end{aligned}$$

18 structure functions

Semi-Inclusive Deep Inelastic Scattering

$$\begin{aligned}
 \frac{d\sigma}{dx dy d\psi dz d\phi_h dP_{h\perp}^2} = & \frac{\alpha^2}{xy Q^2} \frac{y^2}{2(1-\varepsilon)} \left(1 + \frac{\gamma^2}{2x}\right) \left\{ F_{UU,T} + \varepsilon F_{UU,L} + \sqrt{2\varepsilon(1+\varepsilon)} \cos\phi_h F_{UU}^{\cos\phi_h} \right. \\
 & + \varepsilon \cos(2\phi_h) F_{UU}^{\cos 2\phi_h} + \lambda_e \sqrt{2\varepsilon(1-\varepsilon)} \sin\phi_h F_{LU}^{\sin\phi_h} \\
 & + S_{\parallel} \left[\sqrt{2\varepsilon(1+\varepsilon)} \sin\phi_h F_{UL}^{\sin\phi_h} + \varepsilon \sin(2\phi_h) F_{UL}^{\sin 2\phi_h} \right] + S_{\parallel} \lambda_e \left[\sqrt{1-\varepsilon^2} F_{LL} + \sqrt{2\varepsilon(1-\varepsilon)} \cos\phi_h F_{LL}^{\cos\phi_h} \right] \\
 & + |S_{\perp}| \left[\sin(\phi_h - \phi_S) \left(F_{UT,T}^{\sin(\phi_h - \phi_S)} + \varepsilon F_{UT,L}^{\sin(\phi_h - \phi_S)} \right) \right. \\
 & + \varepsilon \sin(\phi_h + \phi_S) F_{UT}^{\sin(\phi_h + \phi_S)} + \varepsilon \sin(3\phi_h - \phi_S) F_{UT}^{\sin(3\phi_h - \phi_S)} \\
 & + \sqrt{2\varepsilon(1+\varepsilon)} \sin\phi_S F_{UT}^{\sin\phi_S} + \sqrt{2\varepsilon(1+\varepsilon)} \sin(2\phi_h - \phi_S) F_{UT}^{\sin(2\phi_h - \phi_S)} \left. \right] \\
 & + |S_{\perp}| \lambda_e \left[\sqrt{1-\varepsilon^2} \cos(\phi_h - \phi_S) F_{LT}^{\cos(\phi_h - \phi_S)} + \sqrt{2\varepsilon(1-\varepsilon)} \cos\phi_S F_{LT}^{\cos\phi_S} \right. \\
 & \left. + \sqrt{2\varepsilon(1-\varepsilon)} \cos(2\phi_h - \phi_S) F_{LT}^{\cos(2\phi_h - \phi_S)} \right] \left. \right\},
 \end{aligned}$$

14 independent azimuthal modulations

Semi-Inclusive Deep Inelastic Scattering

$$\begin{aligned}
 \frac{d\sigma}{dx dy d\psi dz d\phi_h dP_{h\perp}^2} = & \frac{\alpha^2}{xy Q^2} \frac{y^2}{2(1-\varepsilon)} \left(1 + \frac{\gamma^2}{2x}\right) \left\{ F_{UU,T} + \varepsilon F_{UU,L} + \sqrt{2\varepsilon(1+\varepsilon)} \cos\phi_h F_{UU}^{\cos\phi_h} \right. \\
 & + \varepsilon \cos(2\phi_h) F_{UU}^{\cos 2\phi_h} + \lambda_e \sqrt{2\varepsilon(1-\varepsilon)} \sin\phi_h F_{LU}^{\sin\phi_h} \\
 & + S_{\parallel} \left[\sqrt{2\varepsilon(1+\varepsilon)} \sin\phi_h F_{UL}^{\sin\phi_h} + \varepsilon \sin(2\phi_h) F_{UL}^{\sin 2\phi_h} \right] + S_{\parallel} \lambda_e \left[\sqrt{1-\varepsilon^2} \right. \\
 & + |S_{\perp}| \left[\sin(\phi_h - \phi_S) \left(F_{UT,T}^{\sin(\phi_h - \phi_S)} + \varepsilon F_{UT,L}^{\sin(\phi_h - \phi_S)} \right) \right. \\
 & + \varepsilon \sin(\phi_h + \phi_S) F_{UT}^{\sin(\phi_h + \phi_S)} + \varepsilon \sin(3\phi_h - \phi_S) F_{UT}^{\sin(3\phi_h - \phi_S)} \\
 & + \sqrt{2\varepsilon(1+\varepsilon)} \sin\phi_S F_{UT}^{\sin\phi_S} + \sqrt{2\varepsilon(1+\varepsilon)} \sin(2\phi_h - \phi_S) F_{UT}^{\sin(2\phi_h - \phi_S)} \\
 & \left. \left. + |S_{\perp}| \lambda_e \left[\sqrt{1-\varepsilon^2} \cos(\phi_h - \phi_S) F_{LT}^{\cos(\phi_h - \phi_S)} + \sqrt{2\varepsilon(1-\varepsilon)} \cos\phi_S F_{LT}^{\cos\phi_S} \right. \right. \right. \\
 & \left. \left. \left. + \sqrt{2\varepsilon(1-\varepsilon)} \cos(2\phi_h - \phi_S) F_{LT}^{\cos(2\phi_h - \phi_S)} \right] \right\},
 \end{aligned}$$

14 independent azimuthal modulations

amplitudes of the modulations
→ TMD PDFs

Semi-Inclusive Deep Inelastic Scattering

$$\begin{aligned}
 \frac{d\sigma}{dx dy d\psi dz d\phi_h dP_{h\perp}^2} = & \frac{\alpha^2}{xy Q^2} \frac{y^2}{2(1-\varepsilon)} \left(1 + \frac{\gamma^2}{2x}\right) \left\{ F_{UU,T} + \varepsilon F_{UU,L} + \sqrt{2\varepsilon(1+\varepsilon)} \cos\phi_h F_{UU}^{\cos\phi_h} \right. \\
 & + \varepsilon \cos(2\phi_h) F_{UU}^{\cos 2\phi_h} + \lambda_e \sqrt{2\varepsilon(1-\varepsilon)} \sin\phi_h F_{LU}^{\sin\phi_h} \\
 & + S_{\parallel} \left[\sqrt{2\varepsilon(1+\varepsilon)} \sin\phi_h F_{UL}^{\sin\phi_h} + \varepsilon \sin(2\phi_h) F_{UL}^{\sin 2\phi_h} \right] + S_{\parallel} \lambda_e \left[\sqrt{1-\varepsilon^2} \right. \\
 & + |S_{\perp}| \left[\sin(\phi_h - \phi_S) \left(F_{UT,T}^{\sin(\phi_h - \phi_S)} + \varepsilon \sin(\phi_h + \phi_S) F_{UT}^{\sin(\phi_h + \phi_S)} + \varepsilon \right) \right. \\
 & + \sqrt{2\varepsilon(1+\varepsilon)} \sin\phi_S F_{UT}^{\sin\phi_S} + \sqrt{2\varepsilon(1-\varepsilon)} \cos(\phi_h - \phi_S) F_{LT}^{\cos(\phi_h - \phi_S)} + \sqrt{2\varepsilon(1-\varepsilon)} \cos\phi_S F_{LT}^{\cos\phi_S} \\
 & \left. \left. + \sqrt{2\varepsilon(1-\varepsilon)} \cos(2\phi_h - \phi_S) F_{LT}^{\cos(2\phi_h - \phi_S)} \right] \right\},
 \end{aligned}$$

14 independent azimuthal modulations

**amplitudes of the modulations
→ TMD PDFs**

SIDIS

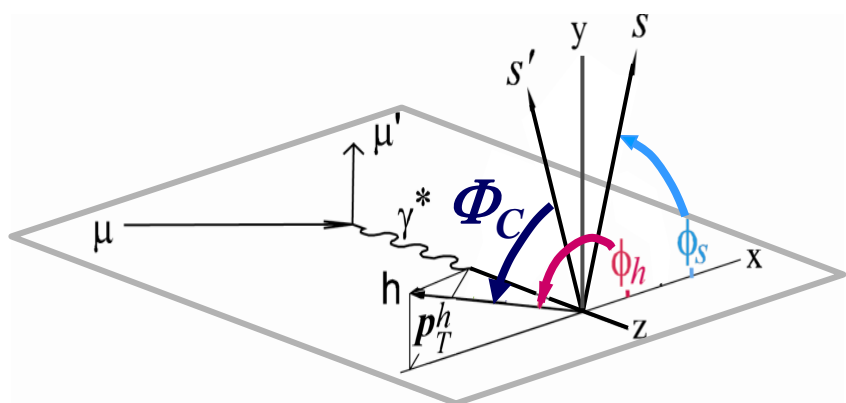
- allows to disentangle the effects related to the different TMD PDFs and to access all of them
 - by identifying the final state hadrons and using different targets allows for flavour separation
- *very powerful tool*

all the amplitudes (AA) have been measured in COMPASS

Collins asymmetry

amplitude of the $\sin \Phi_C$ modulation
in the azimuthal distribution
of the final state hadrons

$$N_h^\pm(\Phi_C) = N_h^0 \left[1 \pm P_T \cdot D_{NN} \cdot \mathbf{A}_{Coll} \cdot \sin \Phi_C \right]$$



$$\Phi_C = \phi_h + \phi_s - \pi$$

transversity

“Collins FF”

$$\mathbf{A}_{Coll} \approx \frac{\sum_q e_q^2 h_{1q} \otimes H_{1q}^{\perp h}}{\sum_q e_q^2 f_{1q} \otimes D_{1q}^h}$$

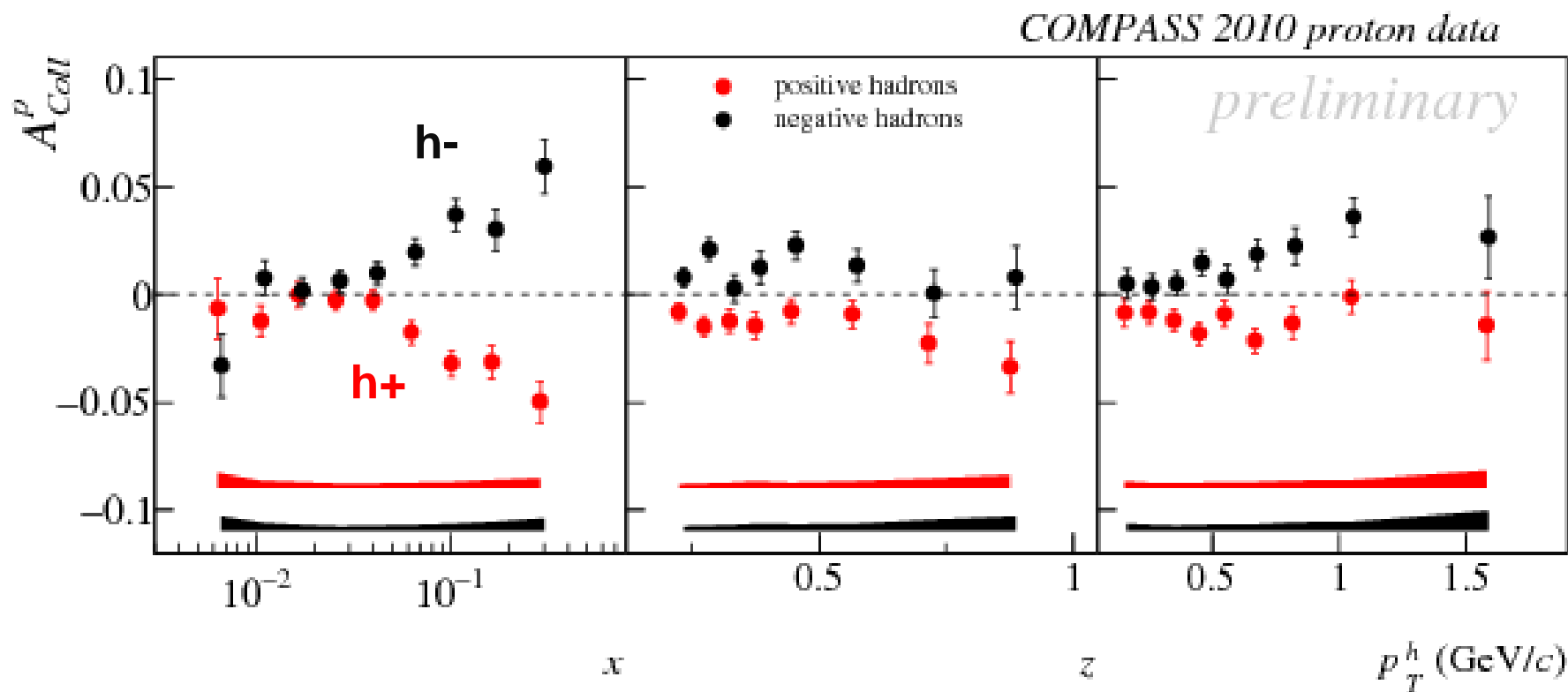
BELLE,
BaBar

today the most promising way
to access transversity, together
with 2h asymmetry

Collins asymmetry on **proton**

charged hadrons

2010 data

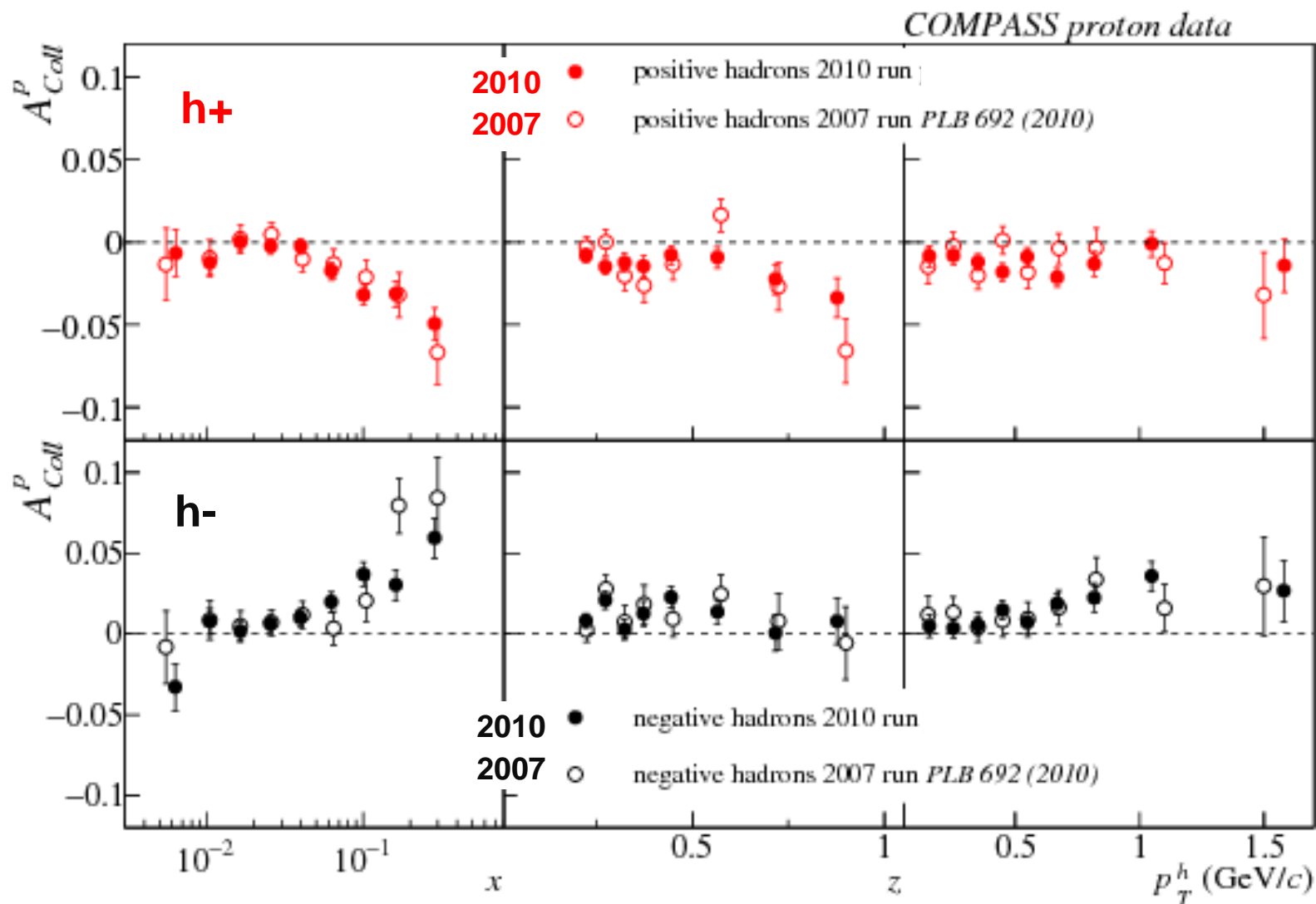


$$\sigma_{\text{syst}} \sim 0.5 \sigma_{\text{stat}}$$

Collins asymmetry on **proton**

charged hadrons

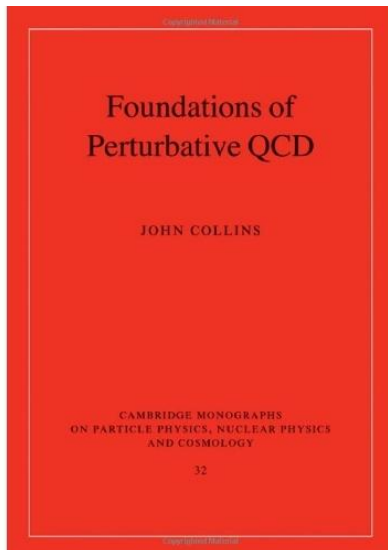
2010 vs 2007 data



TMD factorization

Observables = Universal TMDs × pQCD Coefficients

$$\begin{aligned} W^{\mu\nu} &= \\ &\sum_f |\mathcal{H}_f(Q; \mu)^2|^{\mu\nu} \int d^2\mathbf{k}_{1T} d^2\mathbf{k}_{2T} \delta^{(2)}(\mathbf{k}_{1T} + \mathbf{q}_T - \mathbf{k}_{2T}) \\ &\quad \times F_{f/p}(x, \mathbf{k}_{1T}; \mu; \zeta_F) D_{h/f}(z, z\mathbf{k}_{2T}; \mu; \zeta_D) \\ &= \sum_f |\mathcal{H}_f(Q; \mu)^2|^{\mu\nu} \int \frac{d^2\mathbf{b}_T}{(2\pi)^2} e^{-i\mathbf{q}_T \cdot \mathbf{b}_T} \\ &\quad \times \tilde{F}_{f/p}(x, \mathbf{b}_T; \mu; \zeta_F) \tilde{D}_{h/f}(z, \mathbf{b}_T; \mu; \zeta_D) \end{aligned}$$



Collins, Soper
Ji, Ma, Yuan

The Quark Contribution to the Nucleon Spin

$$\Delta\Sigma = \Delta u + \Delta d + \Delta s$$

in polarised DIS one measures

$$g_1(x) = \frac{1}{2} \sum_q e_q^2 \cdot \Delta q(x) \qquad \Gamma_1 = \int_0^1 g_1(x) dx$$

using complementary information from
the **WEAK DECAY CONSTANTS** of the **BARYONS**

$$\Delta u - \Delta d = F + D = 1.257 \pm 0.003$$

$$\Delta u + \Delta d - 2 \Delta s = 3F - D = \sqrt{3} \cdot [0.34 \pm 0.02]$$

one can get Δu , Δd , Δs and then $\Delta\Sigma$