

Stopping and the K/π horn

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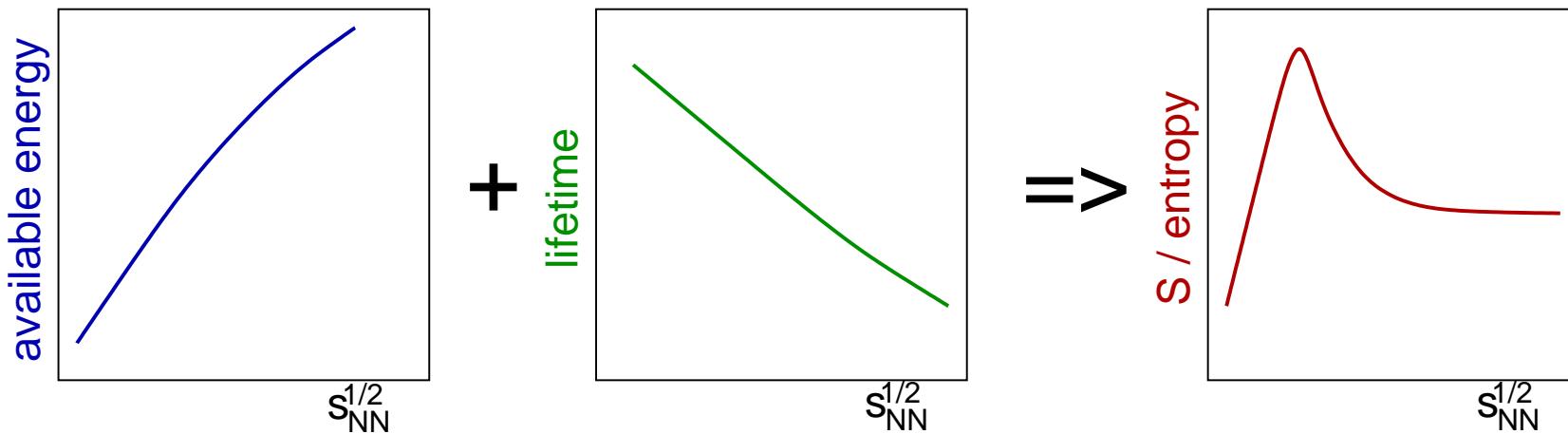
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1 Introduction: strangeness production in nuclear collisions

Two handles for the strangeness production:

- energy (density)
- total time

A cartoon of the working hypothesis:



- we try hadronic scenario; intrinsically assume non-equilibrium

2 The model

- want to calculate **ratios** of yields → look at **densities** of species
- study **evolution of the (kaon, pion, strangeness, . . .) densities**

$$\frac{dn_K}{d\tau} = \frac{d}{d\tau} \frac{N_K}{V} = -\frac{N_K}{V} \frac{1}{V} \frac{dV}{d\tau} + \frac{1}{V} \frac{dN_K}{d\tau}$$

$$\frac{dn_K}{d\tau} = n_K \left(-\frac{1}{V} \frac{dV}{d\tau} \right) + \sum_{ij} \langle v\sigma_{ij}^+ \rangle n_i n_j - \sum_i \langle v\sigma_{Ki}^- \rangle n_K n_i$$

expansion rate

production rate

annihilation rate

ansatz for this

calculate from known cross-sections

and evolved densities

Ansatz for the expansion

$$\varepsilon(\tau) = \begin{cases} \varepsilon_0(1 - a\tau - b\tau^2) & \tau < \tau_s \quad \text{acceleration} \\ \frac{\beta}{(\tau - \tau_0)^\alpha} & \tau > \tau_s \quad \text{power-law expansion} \end{cases}$$
$$n_{B,I}(\tau) = \begin{cases} n_{0;B,I}(1 - a\tau - b\tau^2)^\delta & \tau < \tau_s \quad \text{acceleration} \\ \frac{\gamma}{(\tau - \tau_0)^{\alpha\delta}} & \tau > \tau_s \quad \text{power-law expansion} \end{cases}$$

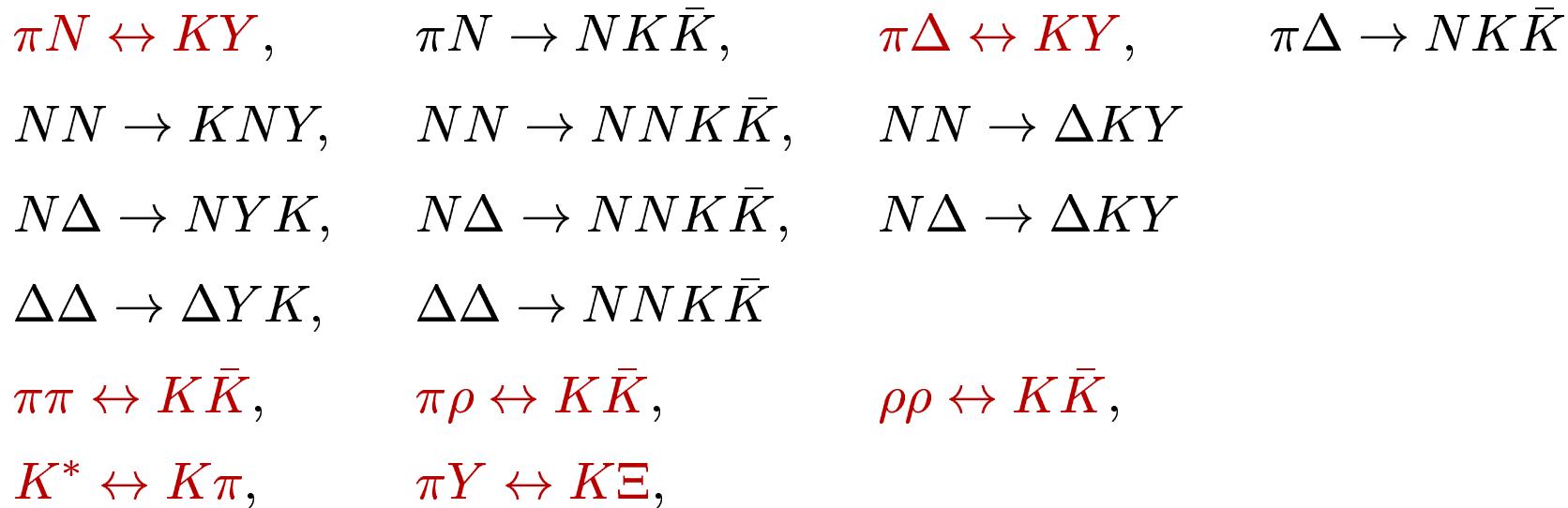
- explore a range of values for the model parameters
- at the end power-law scaling suggested by HBT
- this is a parametrisation “between Landau and Bjorken”

Production and annihilation

Calculation of densities:

- explicit **kinetic calculation**: K^+ , K^0 , K^{*+} , K^{*0}
- chemical equilibrium: non-strange species
- relative chemical equilibrium: $S < 0$ sector (\bar{K} , Λ , Σ , Ξ , Ω)
- no antibaryons assumed at these energies

Implemented K -production (and annihilation) rates:



How to control K^+ , K^- , and Λ production

- density of $S > 0$ species (K^+ , K^0 , K^{*+} , K^{*0}) is controlled by **time** and also temperature
- density of $S < 0$ species must **balance** strangeness such that total strangeness of the system remains 0
- density of K^- **relative to** that of Λ is given by temperature and chemical potentials (relative equilibrium)

Final and initial conditions

- use fit results to chemical FO [Becattini et al., PRC **69** (2004) 024905]:
 $T, \mu_B, \mu_3, \mu_S, \gamma_S \rightarrow \text{get } \varepsilon_{FO}, n_{B,FO}, n_{3,FO}$
- choose initial ε_0 and total lifetime; parametrise evolution from initial state to freeze-out by Becattini et al.
⇒ this helps to aim the parametrisation for the correct final state
- initial strangeness content estimated from pp (pn, nn) collisions

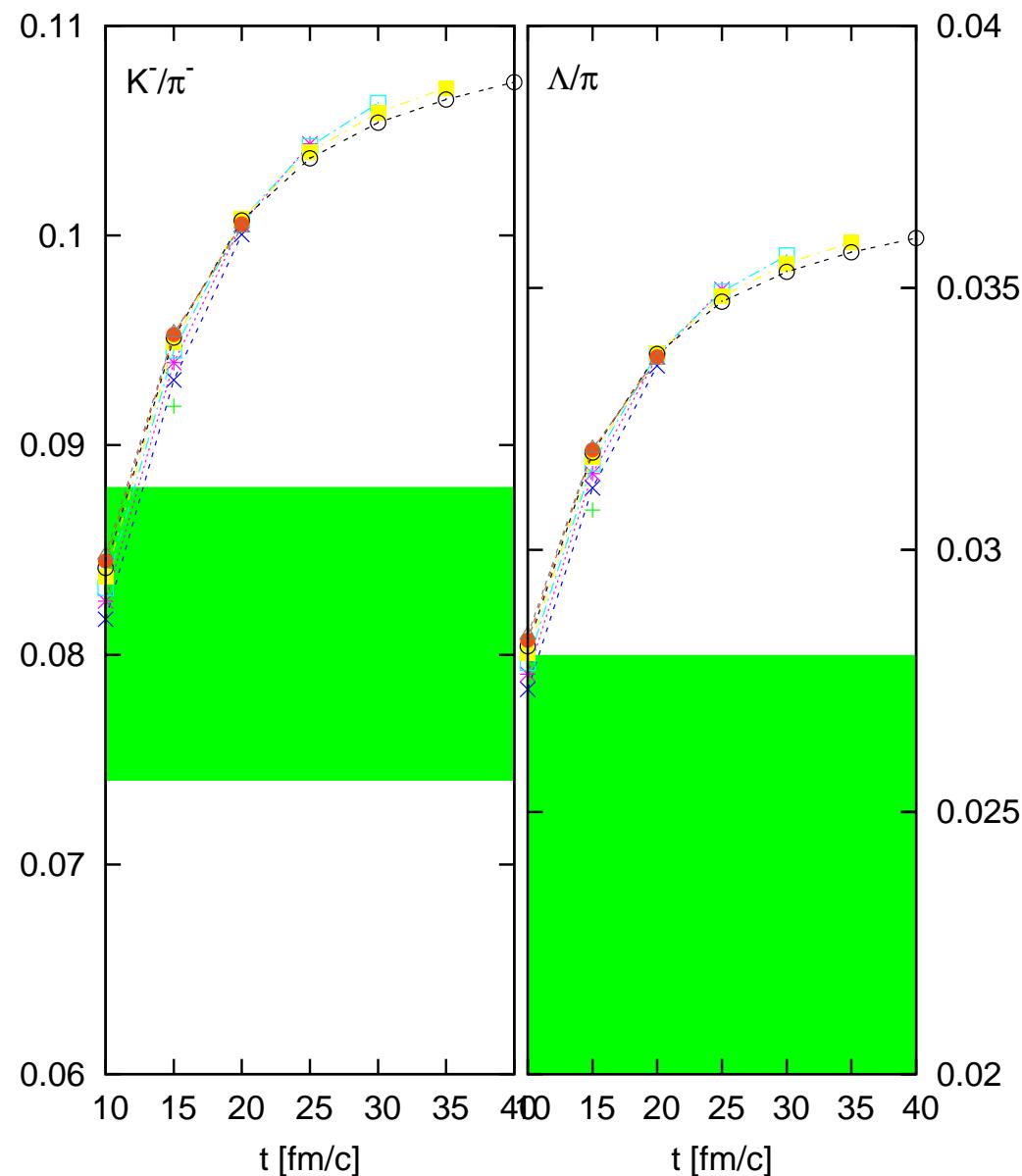
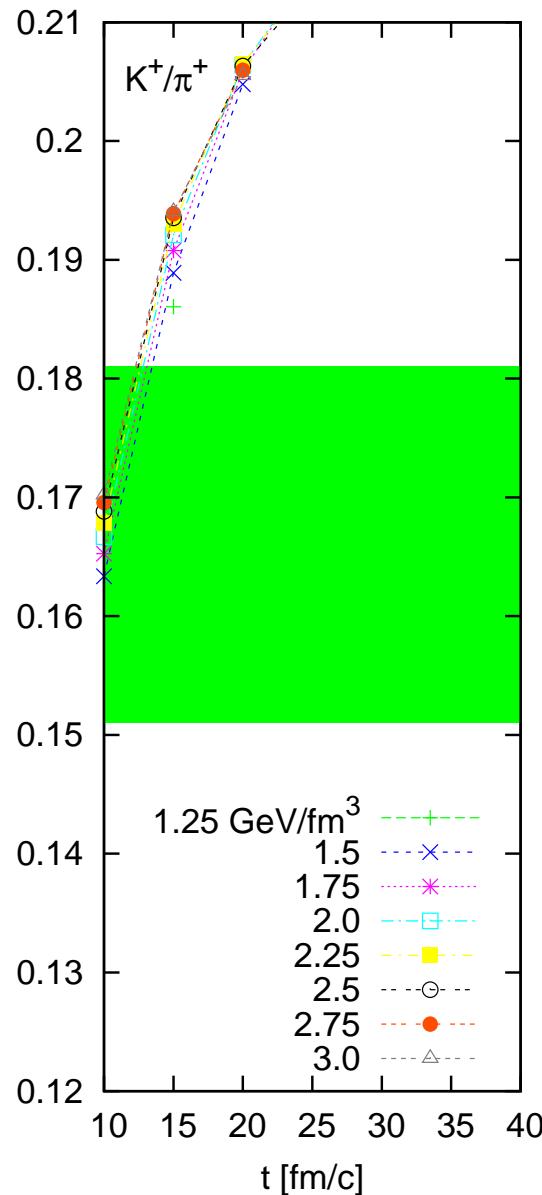
3 Results

- ratios of K^+/π^+ , K^-/π^- , Λ/π as a function of time and initial energy density as they result from the various parametrised evolution scenarios

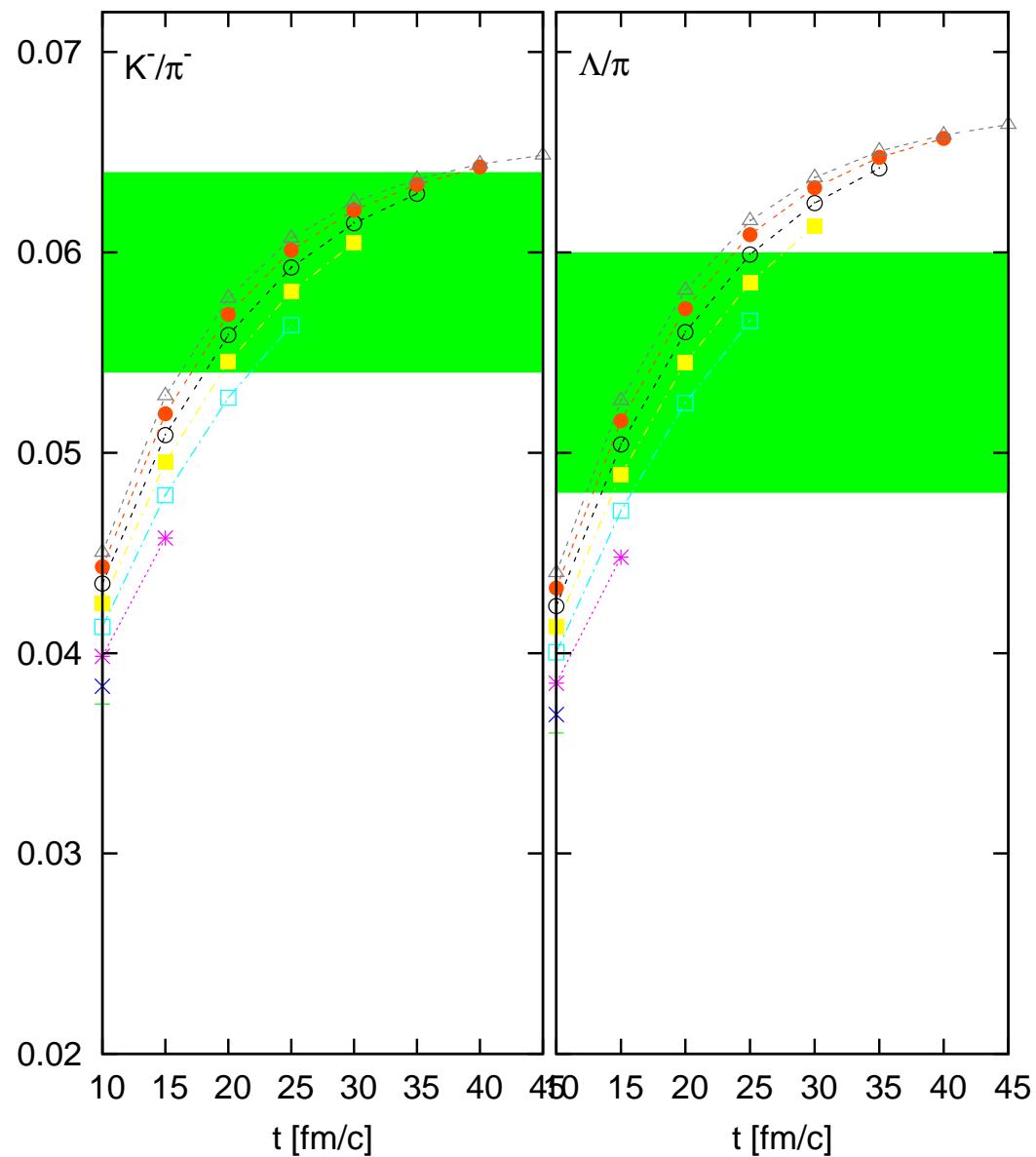
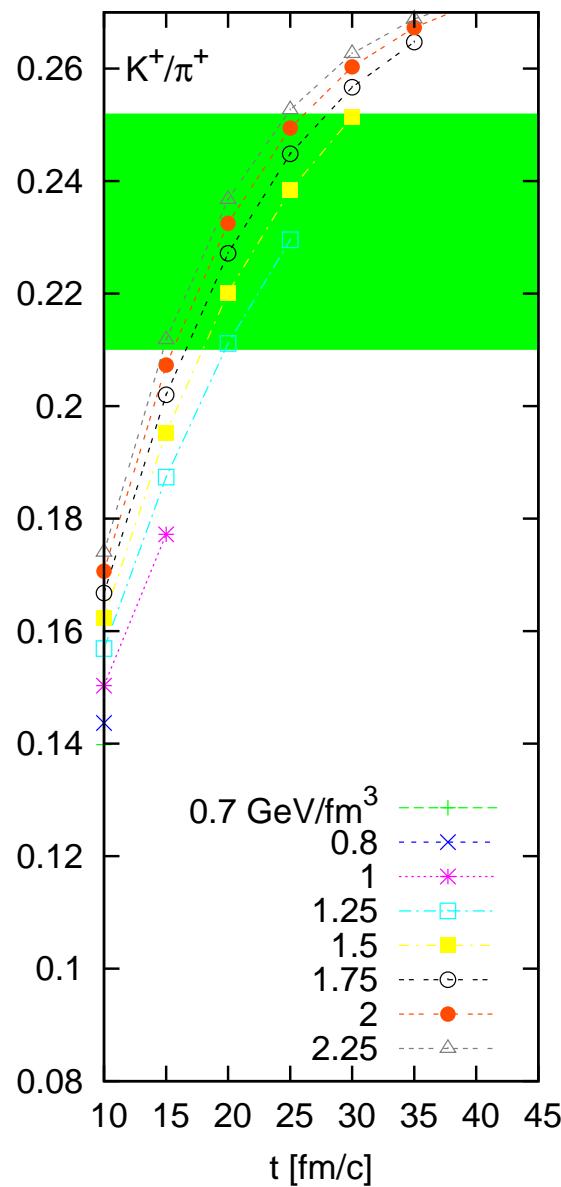
observe:

- strong dependence on the **total lifetime**
- weaker dependence on (initial) energy density (particularly at higher collision energies)
- dependence on total lifetime is plotted
- different curves correspond to different initial energy densities
- green bands show measured values with errorbars

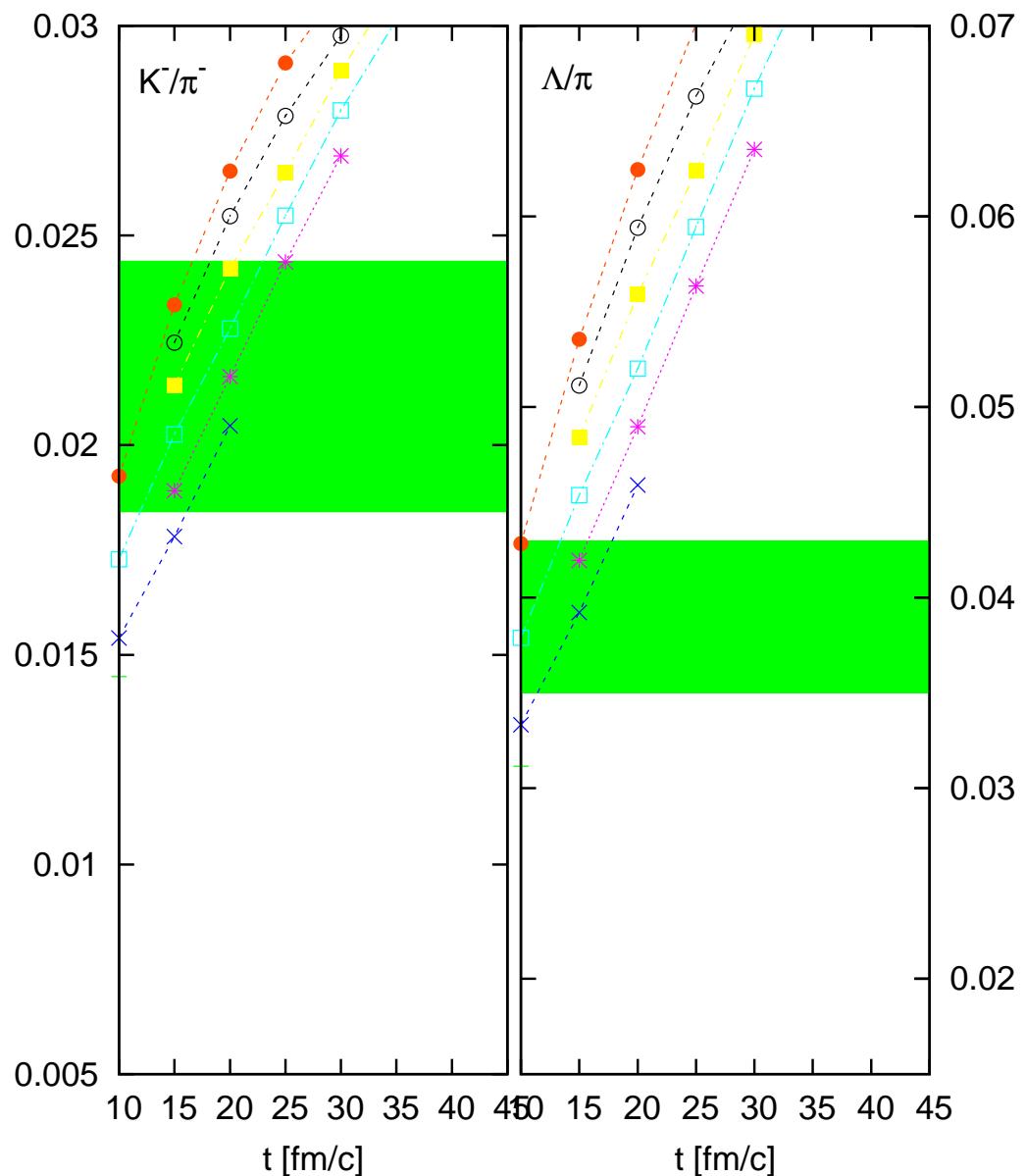
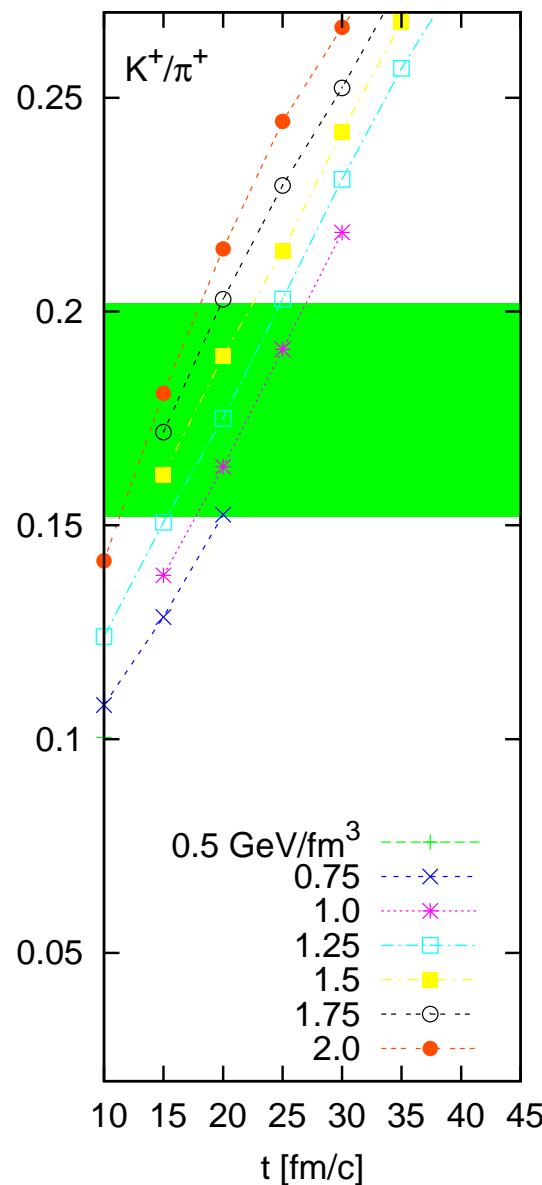
158 AGeV



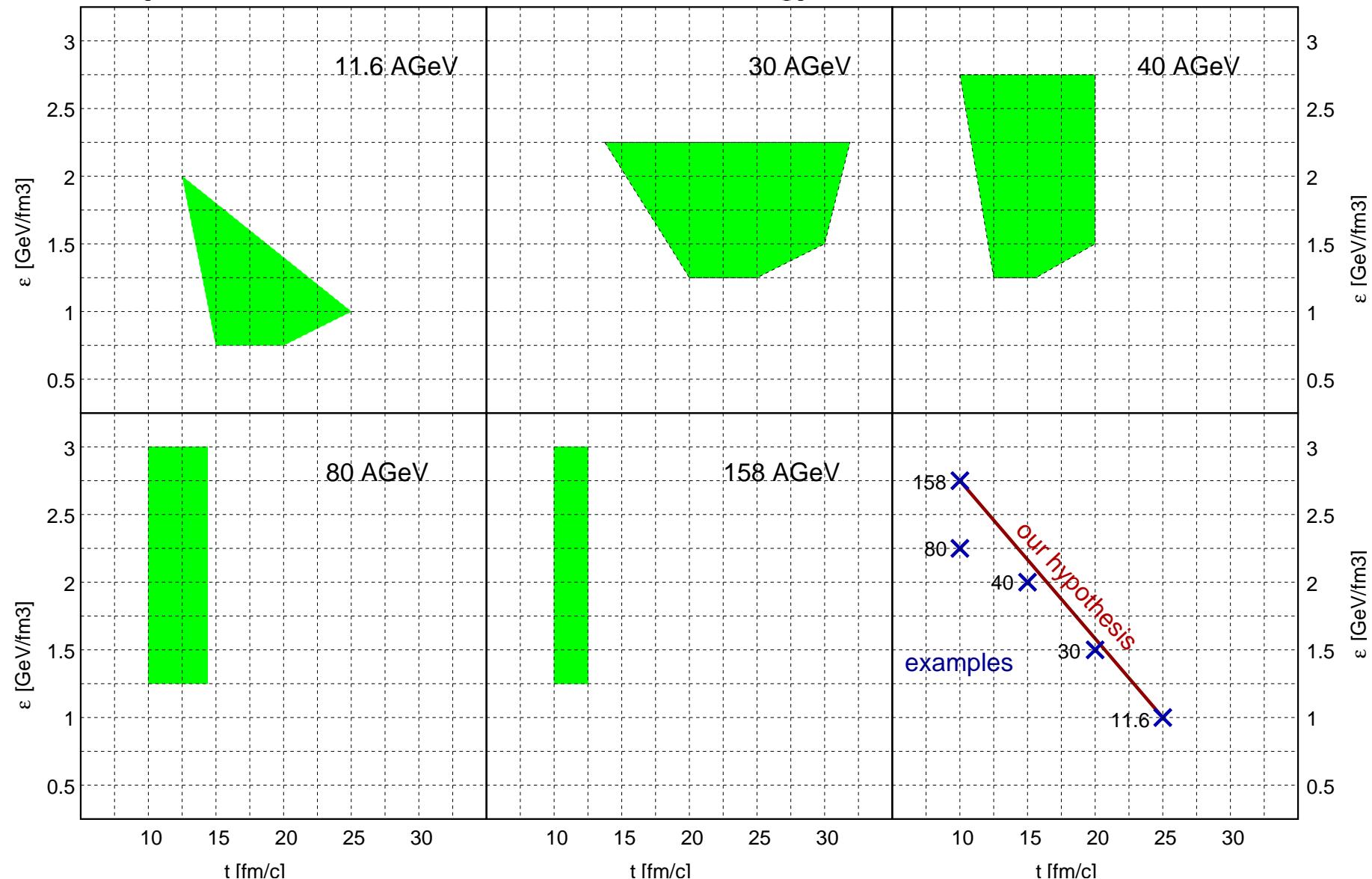
30 AGeV



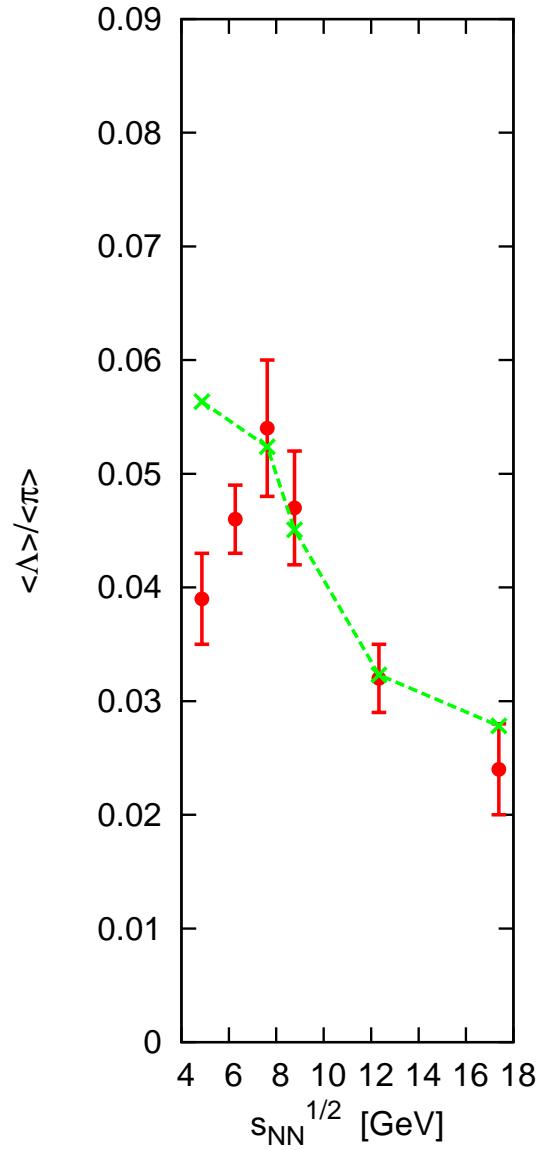
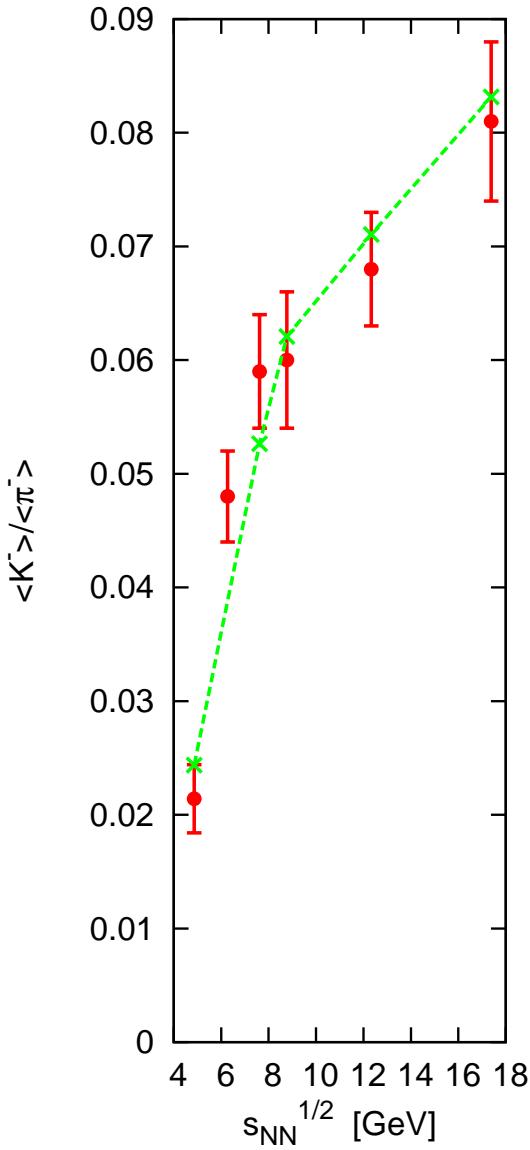
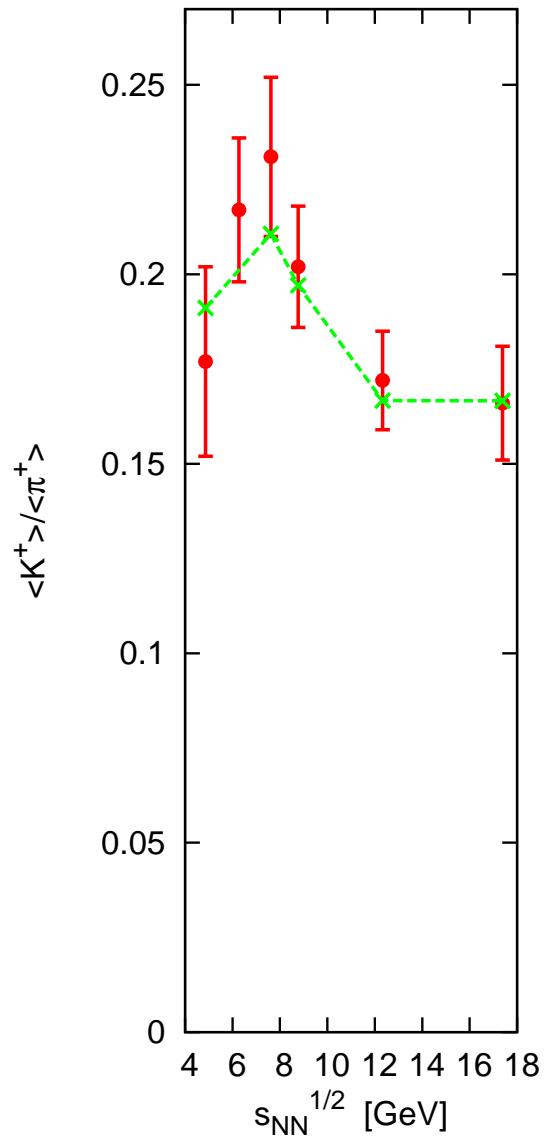
11.6 AGeV



Summary of allowed lifetimes and initial energy densities



Comparison to data



4 Conclusions

- $\sqrt{s_{NN}}$ dependence of strangeness production can be interpreted in terms of decreasing stopping with growing collision energy.
(We may need to re-tune the evolution scenario for AGS energies.)
- this conjecture should be cross-checked with
 - careful analysis of kinetic and chemical freeze-out (spectra, HBT, abundancies)
 - dilepton spectra
 - ...