

Yields and fluctuations in statistical models

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- Definition
- Equilibrium ($T = 170\text{MeV}, \gamma_{q,s} = 1$) vs
explosive ($T = 140\text{MeV}, \gamma_{q,s} > 1$) vs
Continuous emission+resonance modification vs
More complicated models (dynamics)
- Yields and fluctuations in SHM
- Sensitive probes with yields and fluctuations
- Analysis of existing (130 and 200 GeV RHIC data)
- Do it yourself: SHARE
<http://www.physics.arizona.edu/~torrieri/SHARE/share.html>

The statistical model:

$$N = \int \mathcal{M} \prod_i \frac{d^3 \vec{p}_i}{E_i} \delta_E \delta_Q$$

$\mathcal{M} \rightarrow \text{constant}$ (dynamics $\rightarrow$ phase space)

$$P_N = \frac{\Omega_N}{\sum_n \Omega_n} \quad \Omega = \int \prod_i \frac{d^3 \vec{p}_i}{E_i} \delta_E \delta_Q$$

Observables:

$$\langle N \rangle, \quad \omega = \frac{\langle N^2 \rangle - \langle N \rangle^2}{\langle N \rangle}, \quad \text{higher cumulants}$$

calculable through **partition function**

Several ways of defining $\delta_{E,Q} \rightarrow$ **Ensembles**.

Ensembles , or how to deal with conservation laws
 $\lim_{V \rightarrow \infty}^{N/V = \text{const}} \langle N \rangle$ same in $\forall$ ensembles. not ω

Micro-canonical : EbyE conservation

$$\delta_E \delta_Q = \delta \left(\sum_i E_i - E_T \right) \delta \left(\sum_i Q_i - Q_T \right) \quad \omega_E = \omega_Q = 0$$

Canonical : Energy conserved on average
 Appropriate for system in equilibrium with bath

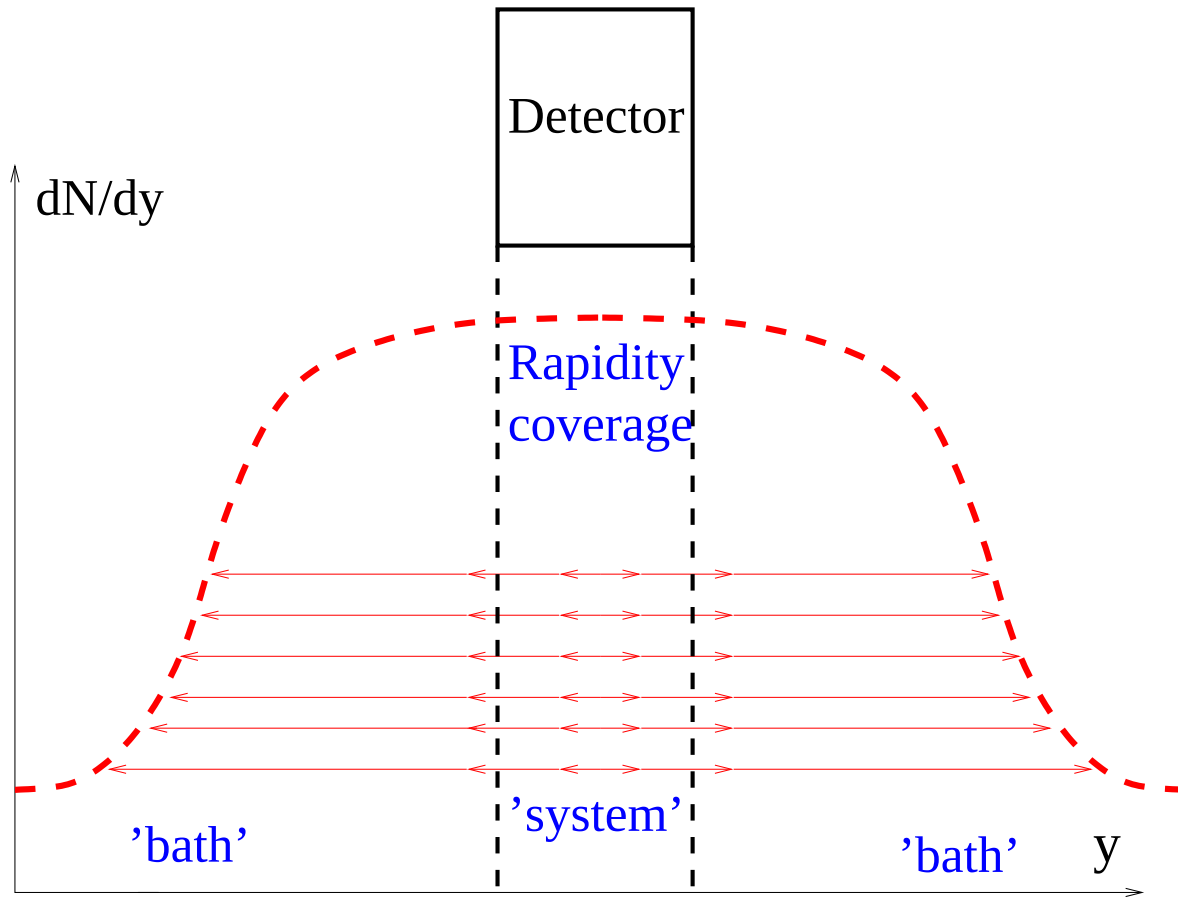
$$\delta_E \rightarrow \delta (E_T - \langle E \rangle) \quad \omega_E \sim 1$$

Grand Canonical : Charge conserved on average

$$\delta_Q \rightarrow \delta (Q_T - \langle Q \rangle) \quad \omega_E \sim \omega_Q \sim 1$$

Appropriate for detector sampling part of a fluid

Freeze-out from ideal fluid at mid-rapidity



Boost invariance: Rapidity $\Leftrightarrow$ configuration space

- Mid-rapidity $\Leftrightarrow$ system
- Peripheral regions $\Leftrightarrow$ bath

$\Rightarrow$ Grand Canonical ensemble needs to be used!

Cleymans, Redlich, PRC 60, 054908 (1999):

$$\left[\frac{dN}{dy} \right]_{b.i.} \sim \langle N \rangle_{4\pi} \quad \left[\frac{d(\Delta N)^2}{dy} \right]_{b.i.} \sim (\Delta N)_{4\pi}^2$$

- All details of flow and freeze-out integrate out
- Up to Normalization, $\langle N \rangle, \omega$ calculable from Grand Canonical T, λ_i

$$\left. \begin{array}{l} \text{Ideal hydro} \\ \text{Fast freezeout} \end{array} \right\} \text{Statistical model fits well} \\ \langle N \rangle \quad \underline{\text{AND}} \quad \langle N^2 \rangle - \langle N \rangle^2$$

Clusters etc. $\Rightarrow$ deviation from hydro.

So lets see how the statistical model does!
But which one?

Grand canonical statistical hadronization

All particles described in terms of T and $\lambda_{q,s,I3}$.

Detailed balance: $\lambda_{\bar{q}} = \lambda_q^{-1}$ Integral can be done in rest-frame wrt flow using Bessel function K_2

$$N_i = V' \sum_{n=1}^{\infty} (\mp 1)^{n+1} \frac{\lambda_i^n}{n} m_i^2 T K_2 \left(\frac{nm_i}{T} \right)$$

$$\Delta N_i^2 = V' \sum_{n=1}^{\infty} (\mp 1)^{n+1} \frac{\lambda_i^n}{n} \binom{2+n-1}{n} m_i^2 T K_2 \left(\frac{nm_i}{T} \right)$$

$$V' = V(2J_i + 1) \frac{4\pi}{(2\pi)^3}$$

Chemical potentials for conserved quantities

$$\lambda_i = \lambda_u^{u-\bar{u}} \lambda_d^{u-\bar{u}} \lambda_s^{s-\bar{s}}$$

Non-equilibrium through phase space occupancies

$$\lambda_i \rightarrow \lambda_i^{\text{eq}} \gamma_u^{u+\bar{u}} \gamma_d^{u+\bar{u}} \gamma_s^{s+\bar{s}} \quad \gamma^{\text{eq}} = 1$$

Resonance feed-down

$$N_i = N_i^{direct} + \sum_j b_{j \rightarrow i} N_j$$

$$\Delta N_i^2 = \Delta N_i^2 + \sum_j [b_{j \rightarrow i} (1 - b_{j \rightarrow i}) N_j + b_{j \rightarrow i}^2 \Delta N_j^2]$$

Widths

Decay $M \rightarrow m_1, m_2, \dots$ width Γ_i , total width Γ_T
relative angular momentum l , threshold mass M_{th}

$$N(M, \lambda) \rightarrow \frac{\int_{M_{Th}}^{\infty} \rho(m) N(m, T, \lambda) dm}{\int_{M_{Th}}^{\infty} \rho(m) dm}$$

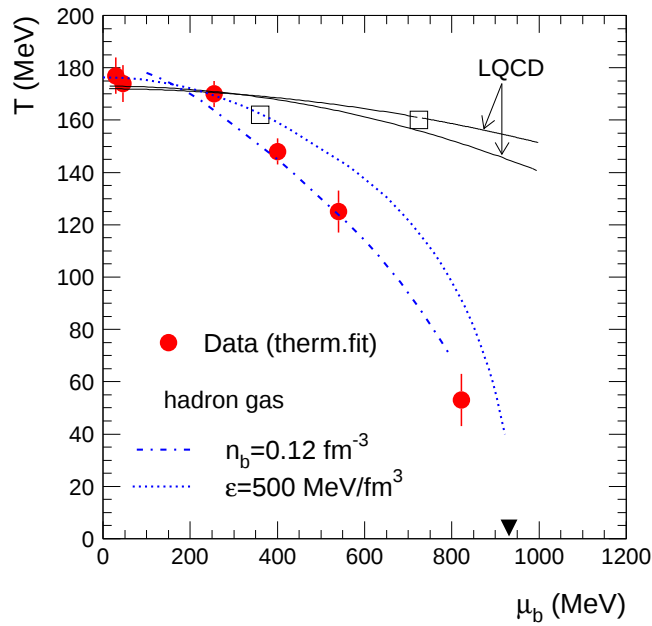
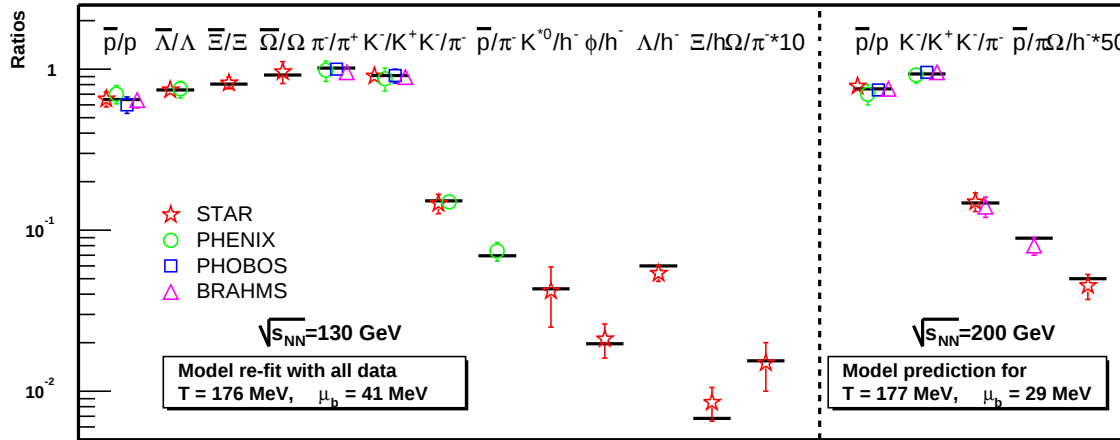
Where

$$\rho(m) = \frac{\Gamma_T \Gamma(m)}{(M - m)^2 + \Gamma^2(m)/4}$$

$$\Gamma(m) = \Gamma_i \left[1 - \left(\frac{M_{Th}}{m} \right) \right]^{l/2}$$

Model I: Equilibrium statistical mechanics

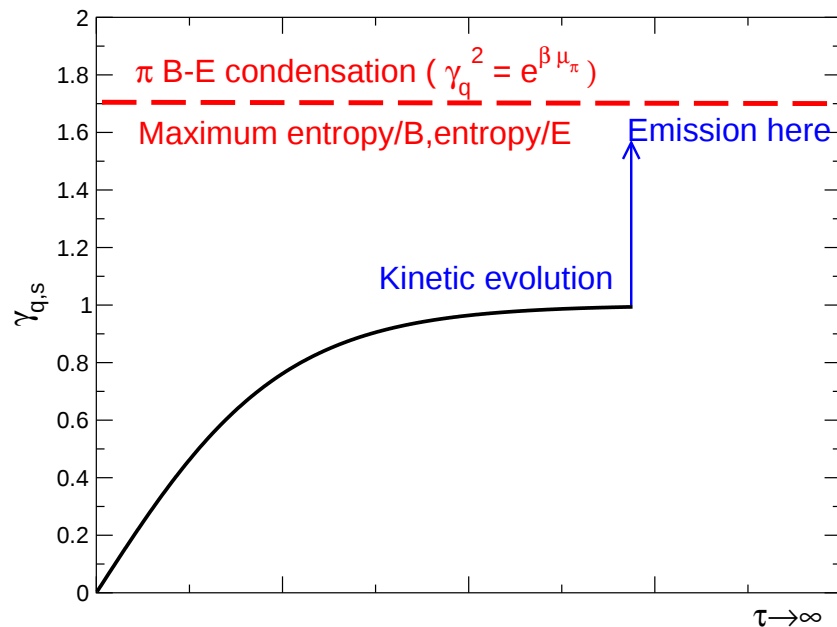
(Braun-Munzinger, Magestro, Florkowski, Broniowski, Redlich,...)



$$\frac{dT}{d\sqrt{s}} > 0 \quad \frac{d\mu_B}{d\sqrt{s}} < 0 \quad \frac{E}{V} \sim 1 \text{ GeV}$$

Resonances not well described (Rescattering?)

Alternatively... Supercooling+oversaturation
(chemical Non-equilibrium but statistical emission)



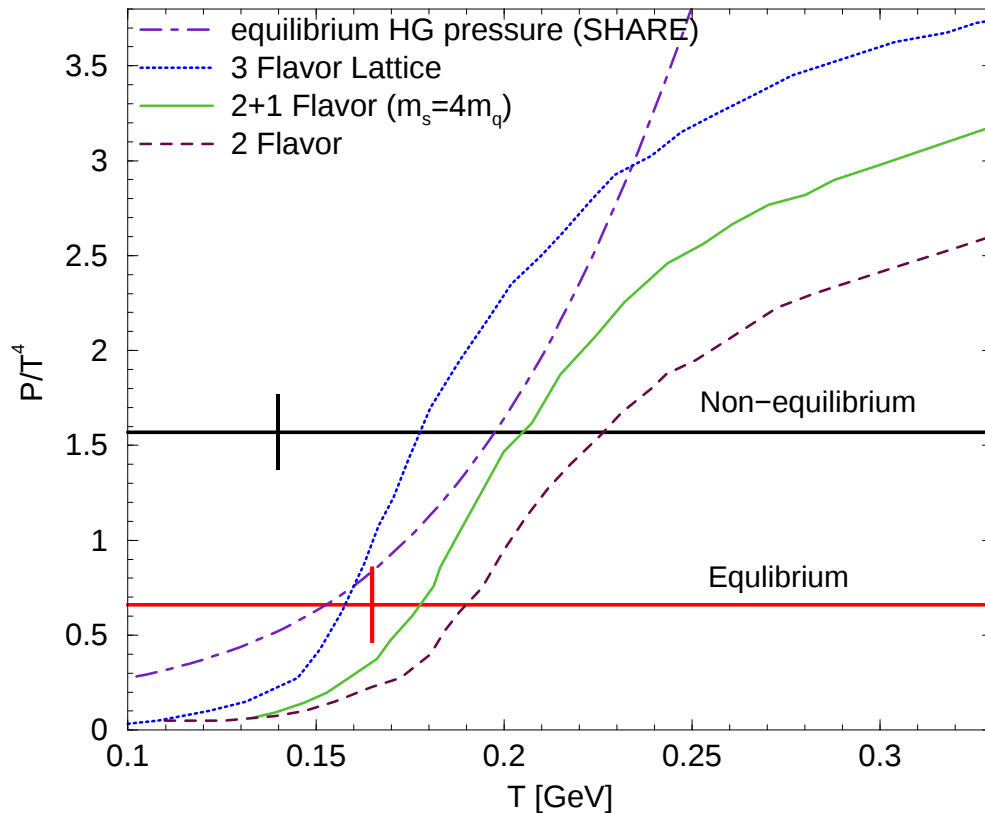
inaccessible by kinetic evolution
(if put “in a box”, this system would heat)
but accessible in a **fast** phase transition
from a **high entropy** phase $\gamma_q > 1, \gamma_s/\gamma_q > 1$.

- **Small $\lambda_{eq} \rightarrow \gamma$ dominates** $\Leftrightarrow$ Statistical coalescence!
- Entropy survives coalescence by $\gamma > 1$

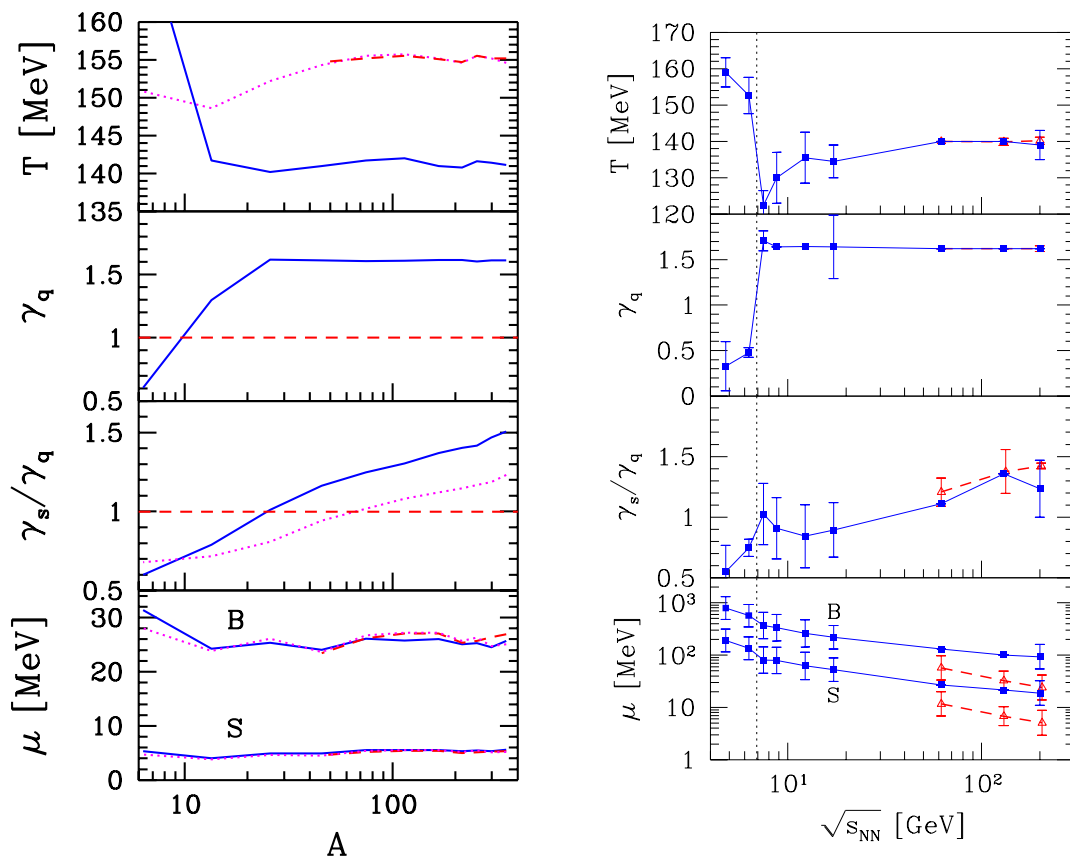
J. Rafelski, J Letessier, PRL 85:4695-4698,2000:
Explosive hadronization from supercooled QGP

$$P_{vacuum} = P_{QGP} \quad S_{HG} = S_{QGP} \quad V \sim \frac{2}{3} V_{equilibrium}$$

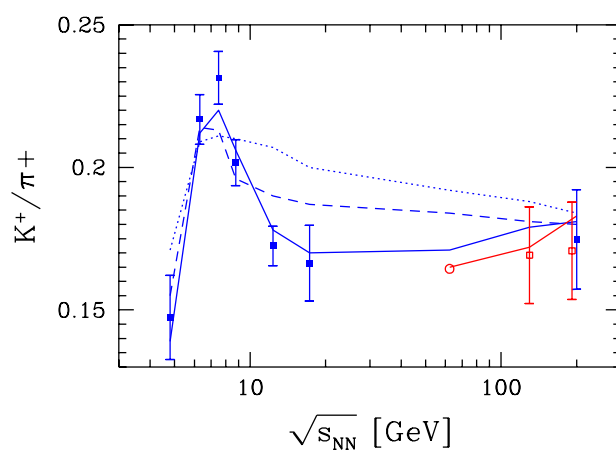
$$T = 140 \text{ MeV}, \quad \gamma_q \sim 0.9 e^{m_\pi/2T} \sim 1.6$$



$\gamma_q > 1, T \rightarrow \sim 140 \text{ MeV} @ \text{critical } \sqrt{s}, N_{part}$



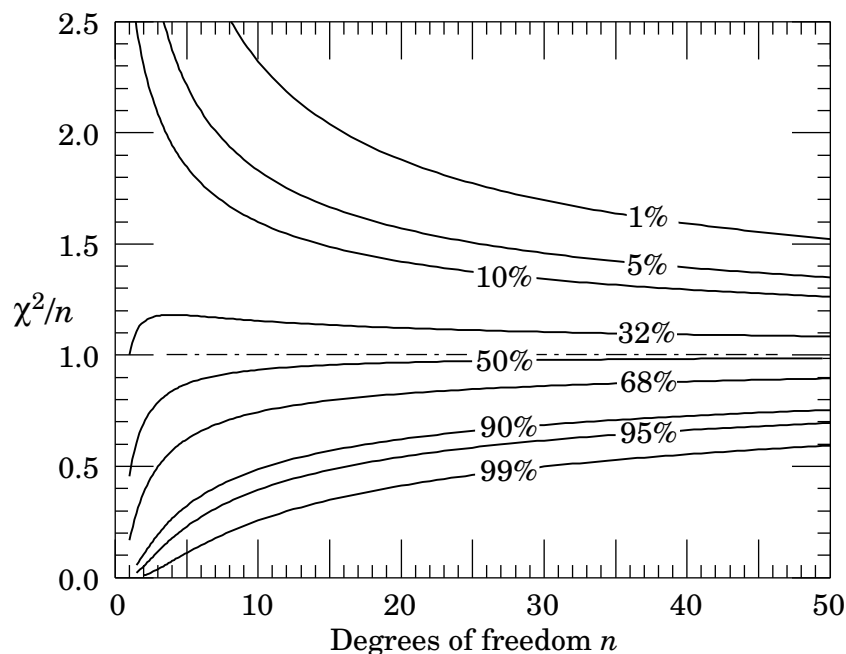
“Horn” explained by $\sqrt{s}$ dependance of T, γ_q



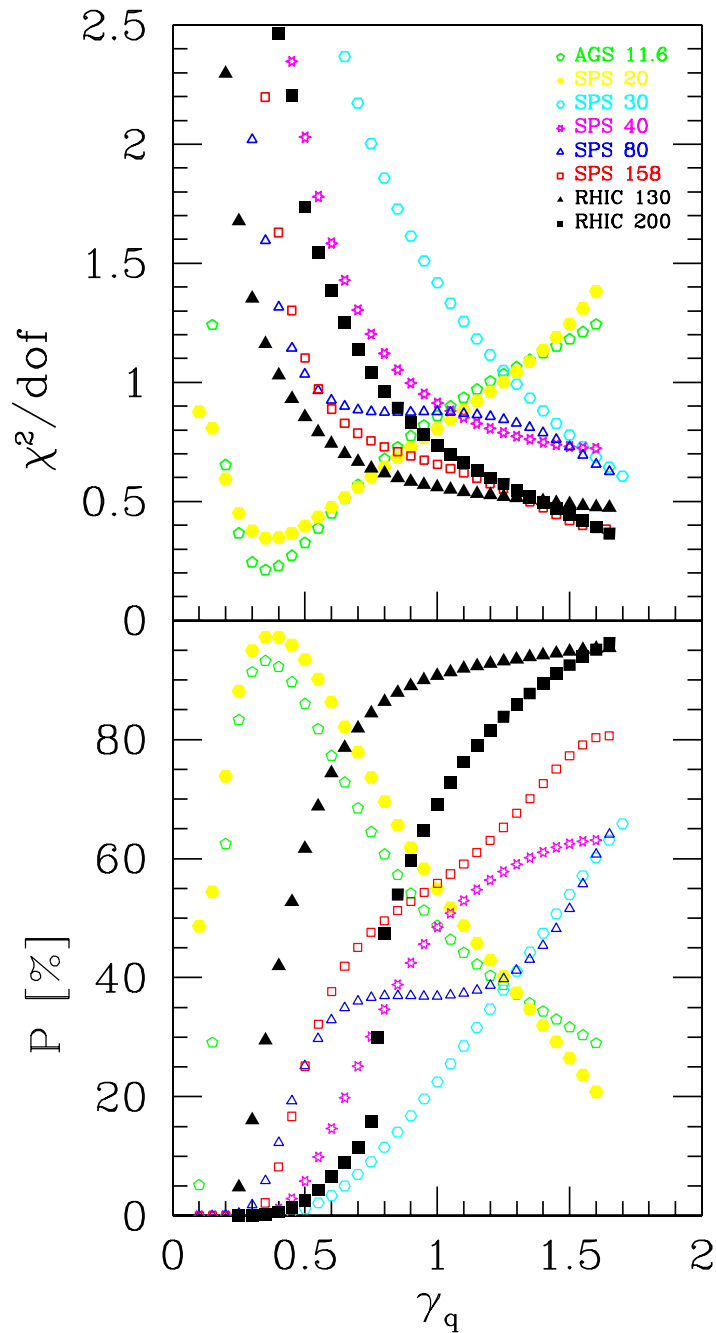
Smoking gun 4 deconfinement?? Perhaps...

Equilibrium and non-equilibrium models have a different number of parameters.

Comparison standard: **Statistical significance**



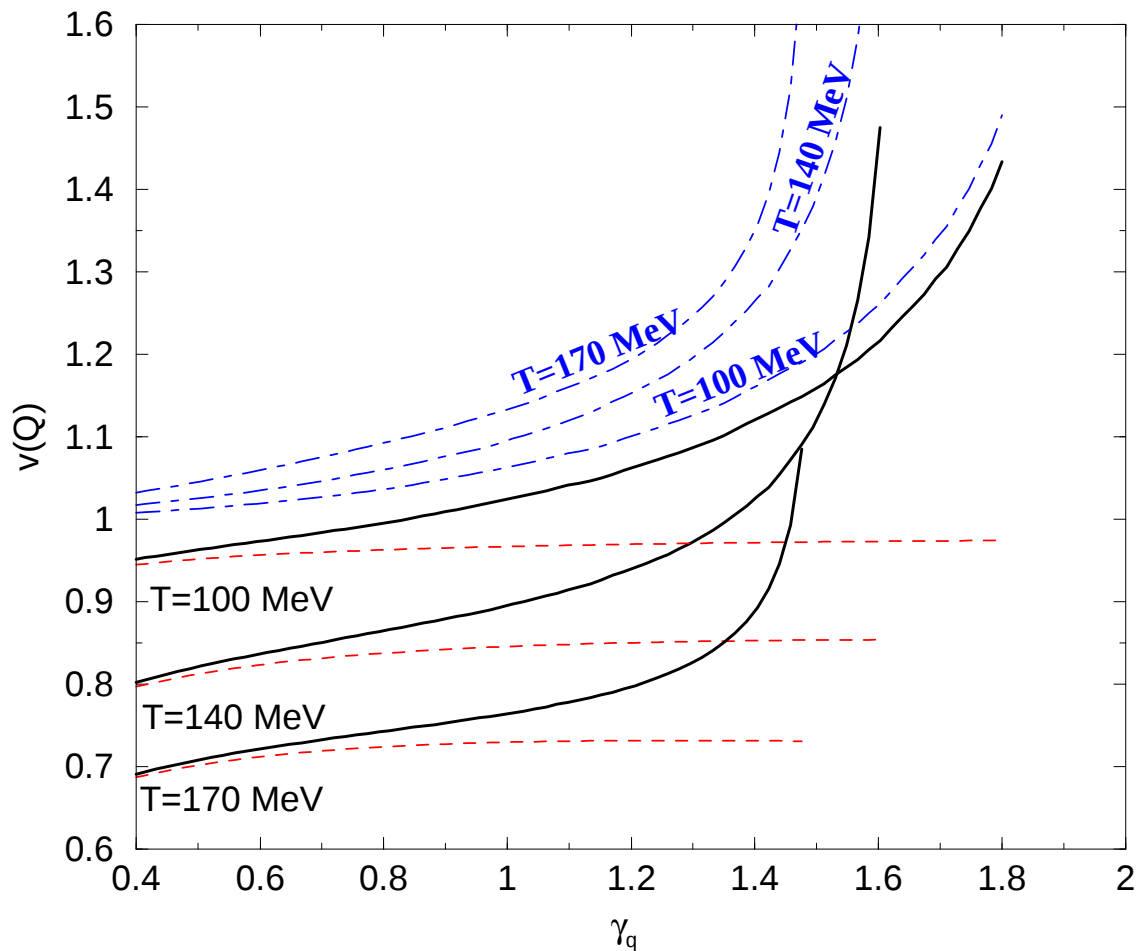
- **Statistical significance**, the probability of getting χ^2 with n DoF given that “your model is true”, is a quantitative measure of your fit’s goodness
- With few DoF, “nice” looking graphs can have a very small statistical significance.
- It is said that you can fit an elephant with enough parameters. Maybe so, but if you are honest, you won’t get a good statistical significance.



Maximum for SPS and RHIC is at $\gamma_q > 1$, but equilibrium not ruled out!

Need further data capable of determining γ_q .
fluctuations!

Yields and fluctuations: Non-equilibrium



T increase $\Rightarrow \pi$ Fluctuations decrease because of enhanced resonance production

over-saturation ($\gamma_q > 1$) $\Rightarrow \pi$ Fluctuations increase because of BE corrections

$\gamma_q > 1$, unlike resonances (detectable, rescattered, in-medium modified,...) affects **fluctuations** rather than **correlations**

Yields and fluctuations: Reinteraction (or not)

Consider $Y^* \rightarrow Y\pi$

$\sigma_{Y/\pi}$ probes correlation of Y and π from Y^*
at chemical freeze-out.

(further rescattering/regeneration does not change the correlation.

Y^*/Y **yield** probes Y^* at thermal freeze-out (after all rescattering.

So...

- If can fit **stable particles** and **resonances** and **fluctuations** in same fit $\rightarrow$ **no reinteraction**
- If **Stable particles** + **Fluctuations** fit gives wrong value for **resonances** $\rightarrow$ magnitude of reinteraction

Suitable:

Yields Independent of $\gamma_s, \lambda_s, \text{volume}$

- Λ/K^- , better
 - Λ corrected for $\Xi, \Omega \rightarrow \Lambda$
 - K^- corrected for $\phi \rightarrow K^- K^+$
- Ξ/ϕ

Fluctuations $v(Q), \sigma_{\pi^+/\pi^-}$

$$v(Q) = \frac{\langle Q^2 \rangle - \langle Q \rangle^2}{N_{ch}}$$

$$\sigma_{N_1/N_2} = \frac{\omega_{N_1}}{\langle N_1 \rangle} + \frac{\omega_{N_2}}{\langle N_2 \rangle} - 2 \frac{\langle N_1 N_2 \rangle - \langle N_1 \rangle \langle N_2 \rangle}{\langle N_1 \rangle \langle N_2 \rangle}$$

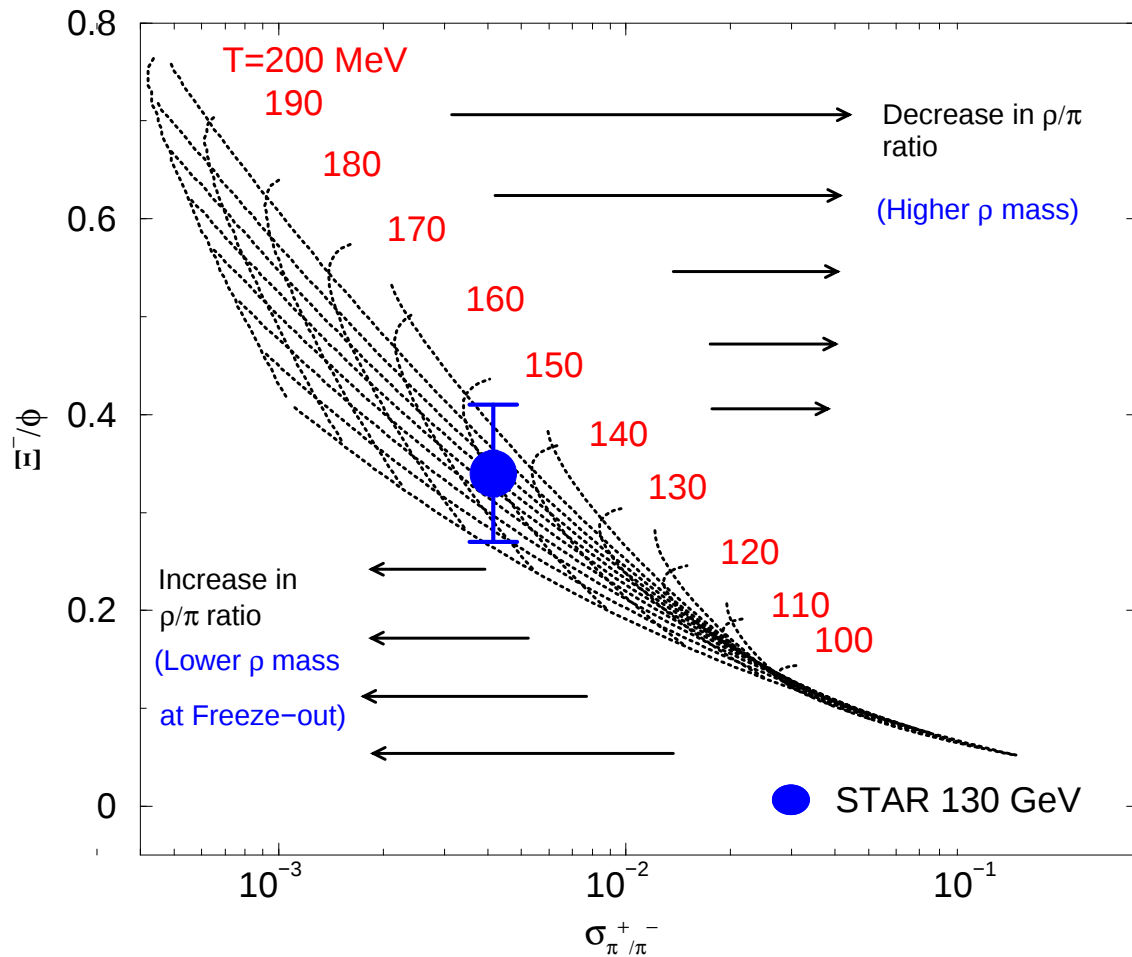
σ_{π^+/π^-} can also be used to probe ρ, f_0 mass modification due to $\pi^+ - \pi^-$ correlations

undetectable or rescattered resonances also contribute to fluctuations!!!

Volume fluctuations are not well understood, and show up in all $\langle N^2 \rangle - \langle N \rangle^2$. Avoid them choosing observables such as

- $(\Delta Q)^2$. $\frac{\langle Q \rangle}{V}$ small, so is $\Delta V \frac{\langle Q \rangle}{V}$
(Jeon, Koch)
- Fluctuations of ratios (Jeon, Koch)
Volume fluctuations irrelevant to 1st order
- For most other data-points can fit ΔV , $(\Delta N)^2 = V(\Delta \rho)^2 + [\Delta V \langle N \rangle]^2$

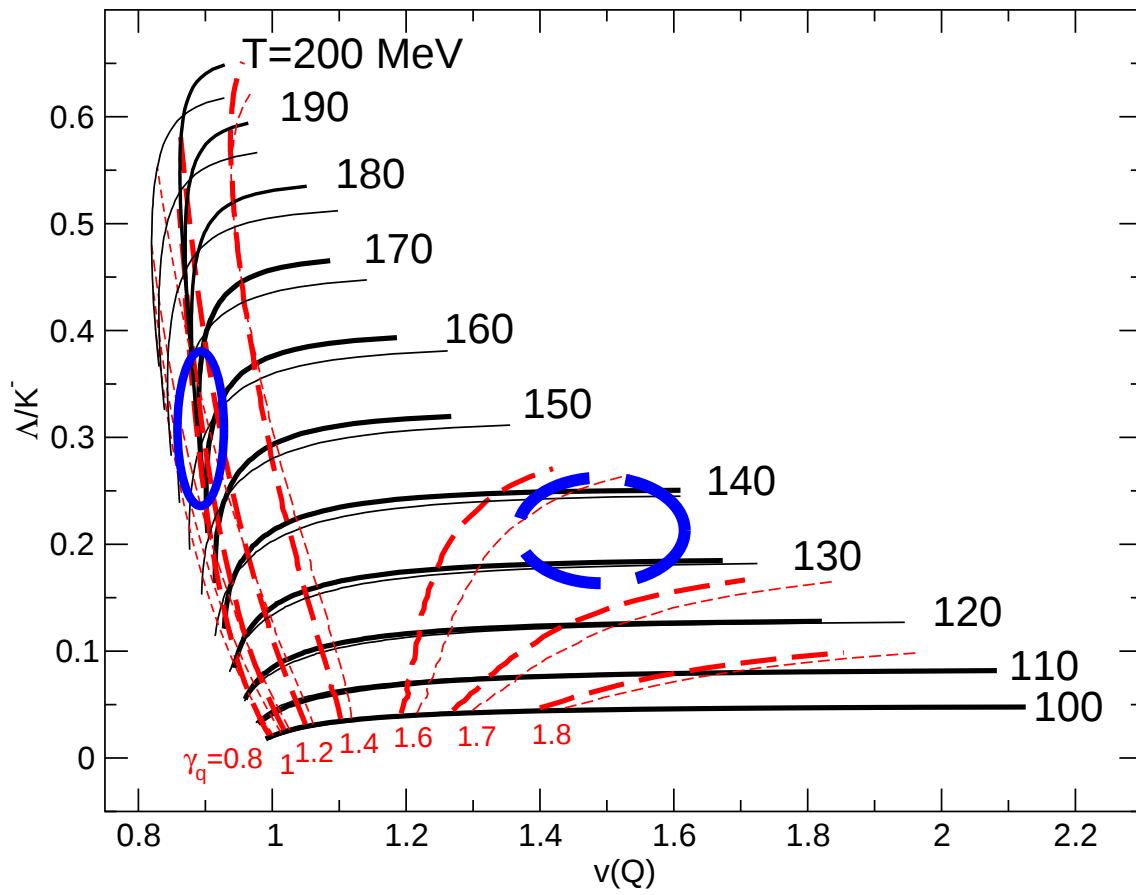
σ_{π^+/π^-} vs Ξ^-/ϕ



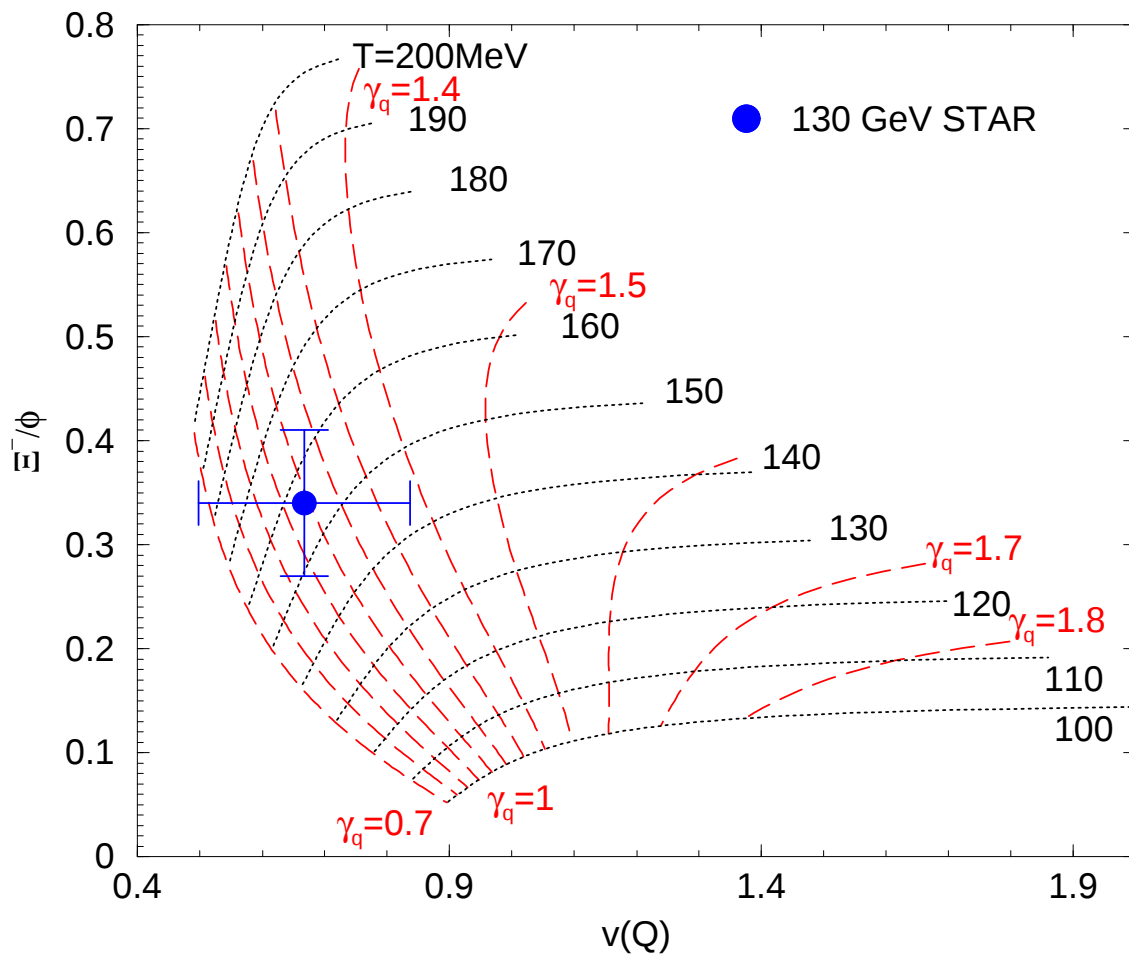
- σ_{π^+/π^-} probes T independently of γ_q
- Allowed region narrow $\rightarrow$ mass modification probe

130 GeV within band! Good agreement with non-equilibrium No freeze-out ρ, σ modification needed

$v(Q)$ vs Λ/K^-

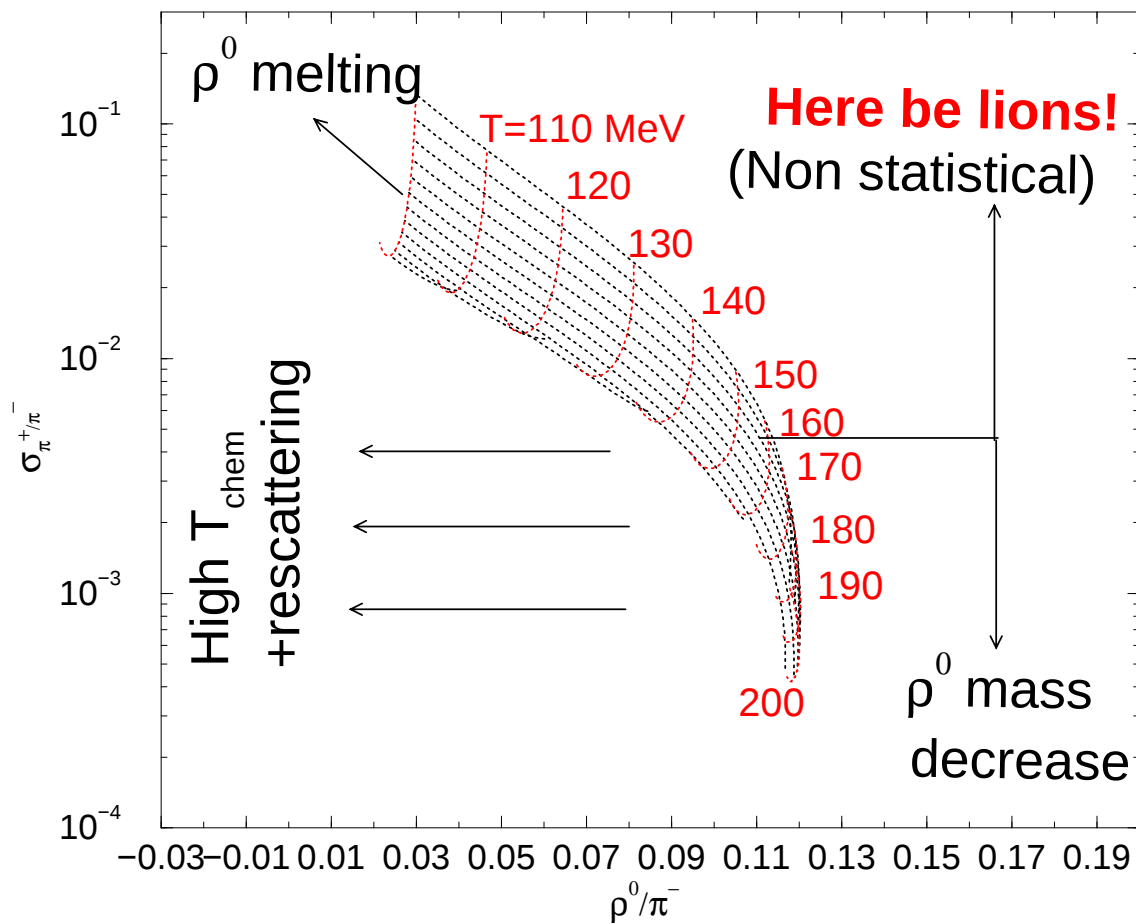


$v(Q)$ vs Ξ^-/ϕ

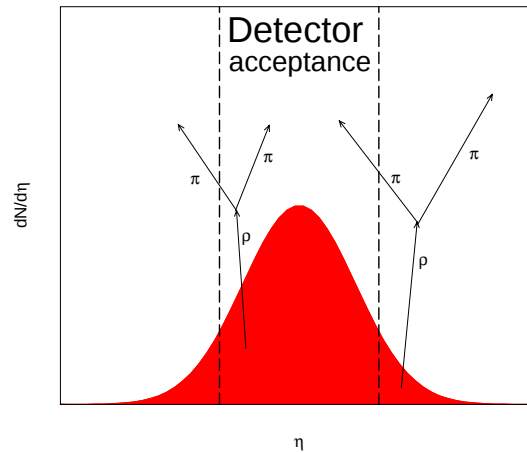


σ_{π^+/π^-} vs ρ^0/π^-

Probes (lack of?) reinteraction and mass
modification separately



These diagrams are made with static fluctuations
susceptible to detector response effects



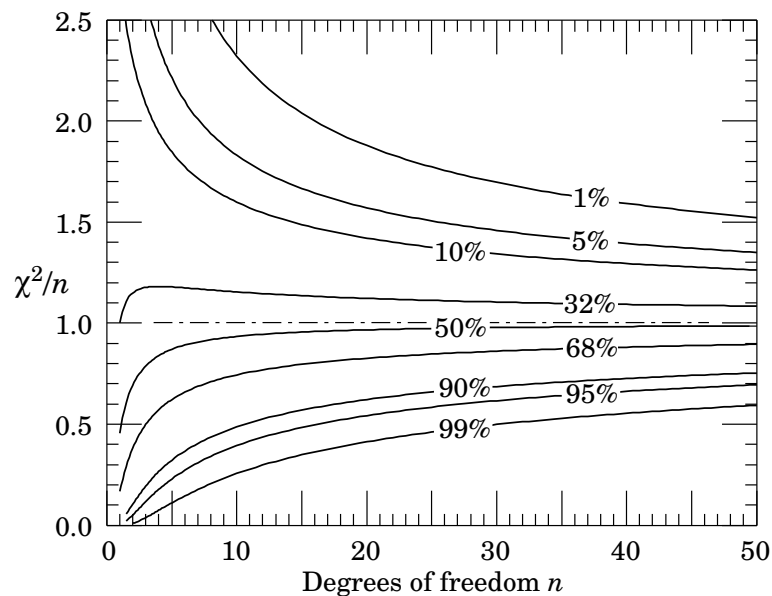
use dynamical fluctuations $\sigma_{dyn} = \sigma - \sigma_{stat}$ Where
 $\sigma_{stat} \sim \frac{1}{\langle N_1 \rangle} + \frac{1}{\langle N_2 \rangle}$ obtained by mixed event
technique

σ_{dyn} robust against detector acceptance but needs
more parameters (“volume”) to be described $\Rightarrow$ no
diagrams. Can use it in fit, including one/more yields
at same centrality as σ_{dyn} .

For large acceptance, can hope fitting both σ_{dyn} and
 $\sigma, \omega, v(Q)$

Fit exp. yields, ratios, $\omega, \sigma, \langle s \rangle = 0, \frac{Q}{B} = 0.4$ for

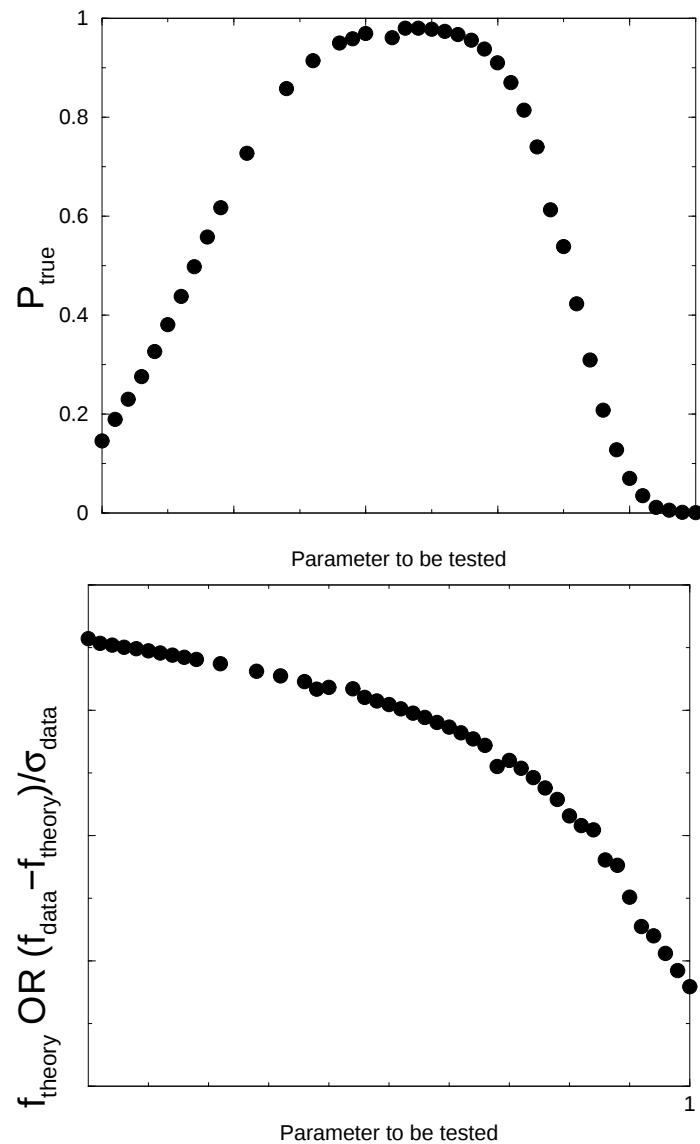
- Equilibrium parameters T, λ_{q,s,I_3}
- Non-equilibrium parameters γ_{q,s,I_3}
- System “volume” dV/dy



Different models have a different DoF $\rightarrow$ Use Statistical significance (P_{true}) to judge fit quality

Non-trivial correlations/data-point sensitivity can be analyzed by Profiles in statistical significance

All other parameters at their best fit value for point in abscissa



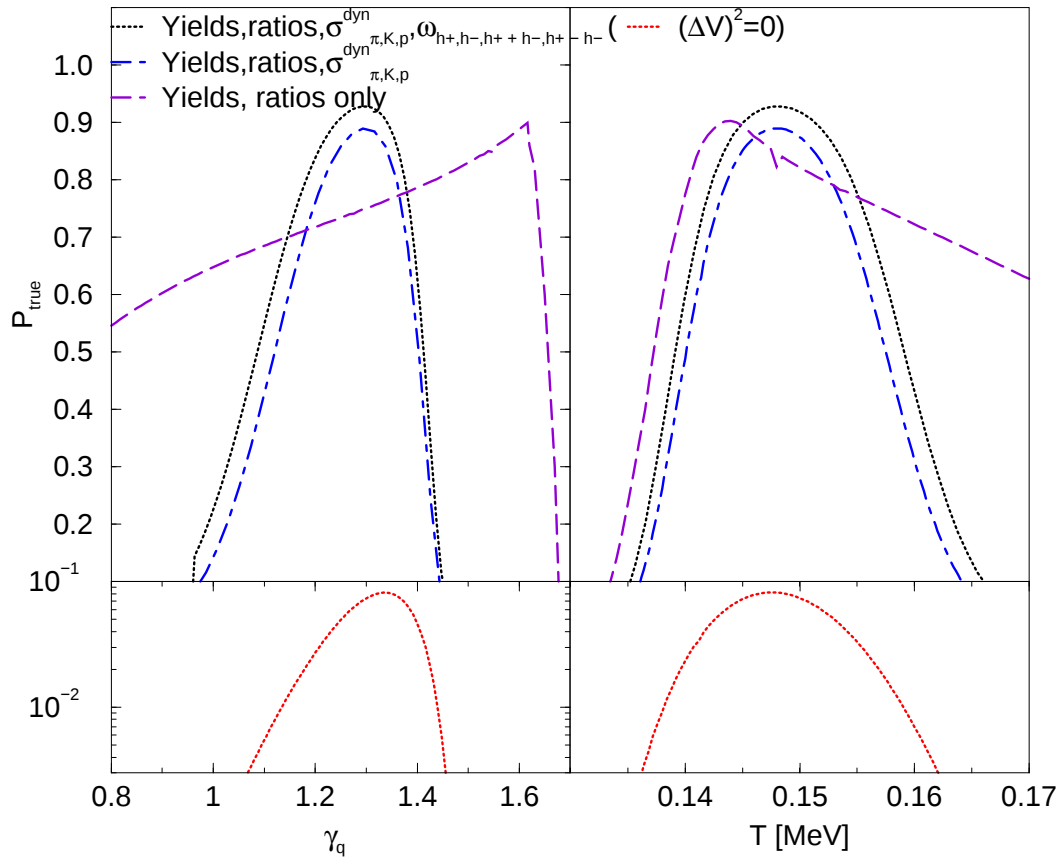
Fits at 130 GeV

fluctuations

- $\sigma_{\pi^+/\pi^-}^{dyn}, \sigma_{K^+/K^-}^{dyn}, \sigma_{\bar{p}/p}^{dyn}$: STAR
PRC 68, 044905 (2003), nucl-ex/0307007
Insensitive to Volume fluctuations but require $\langle V \rangle$ (6% centrality)
- $\omega_{h^+}, \omega_{h^-}, \omega_{h^+-h^-}, \omega_{h^-+h^+}$ STAR
Nucl. Phys. A 698 (2002) 611
Need $\Delta V/V$, independent of centrality.
No error bar, assume Ad Hoc 10%

Yields and Ratios

- $h^{charged}$, PHOBOS
PRL (2003) , nucl-ex/0201005 (6% centrality)
- $\pi, K, p, \Lambda, \Xi, \phi, K^*$ and antiparticles, STAR,
various (5% centrality)



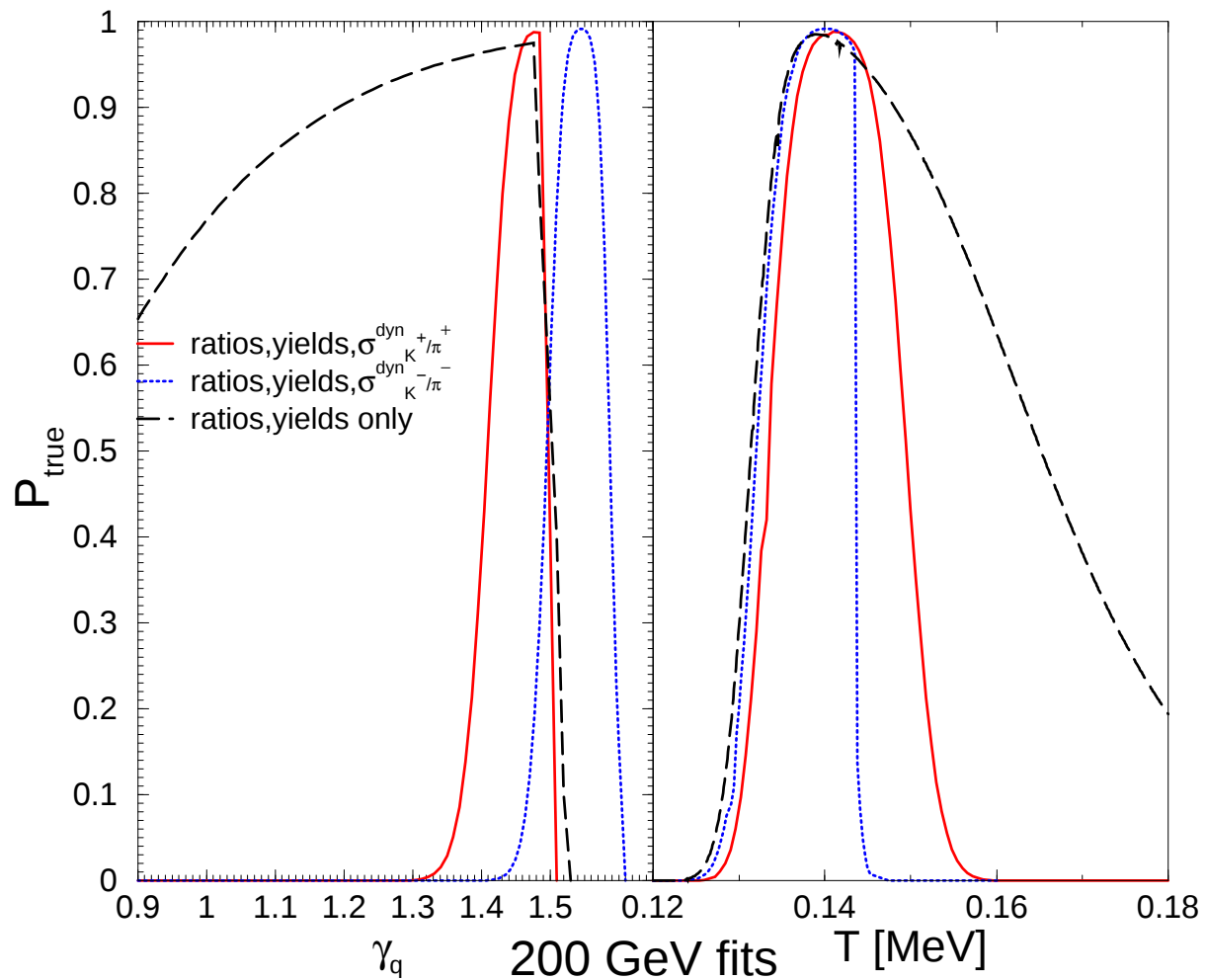
T, γ_q not as well determined as it could be

- As we saw earlier, $\sigma_{\pi,K,p}^{dyn}$ are not very good constraints on γ_q
- **Volume fluctuation** essential for fitting $\omega_{h+}, \omega_{h-}, \omega_{h+\pm h-}$ (if $\Delta V = 0$, $P_{true} < 0.1$ due to disagreement with these data points), but correlates with normalization, γ_q .

Fits at 200 GeV

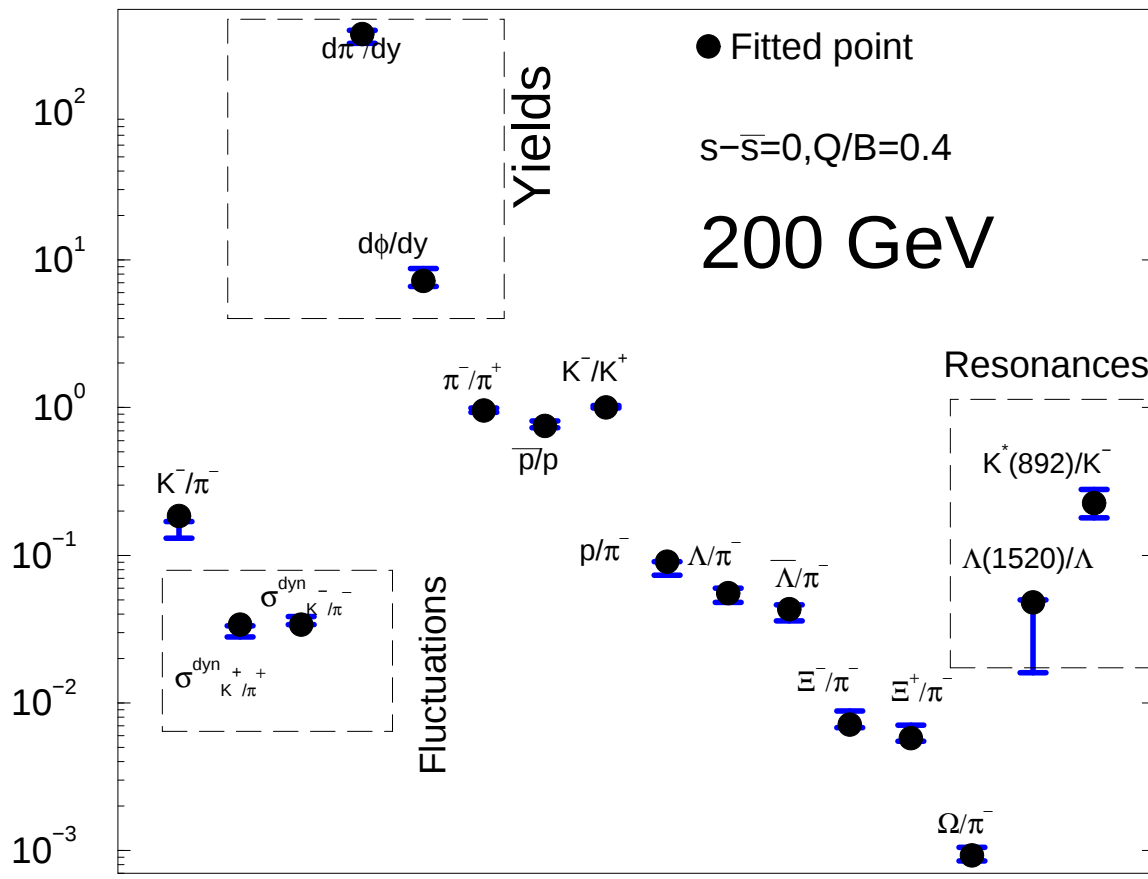
- $\sigma_{K/\pi}^{dyn}$: Supriya Das et al [STAR]
nucl-ex/0503023
- Ratios: O. Barannikova et al [STAR]
nucl-ex/0403014

NB: All preliminary



With fluctuations, T, γ_q determined
 ($K/\pi+$ its fluctuation is very sensitive to γ_q).

- firmly in $\gamma_q > 1, T \sim 140 \text{ MeV}$,
- T, γ_q consistent with 130 GeV considering
all preliminary



Some datapoints fail but no exp. systematic errors yet!

σ_{K^+/π^+} vs σ_{K^-/π^-} :	too different
$\bar{\Omega}/\Omega > 1$	<u>not</u> thermal: Bleicher, Liu, Aichelin

If these fitted, best fit unchanged but $P_{true} < 0.1$

Some tentative conclusions

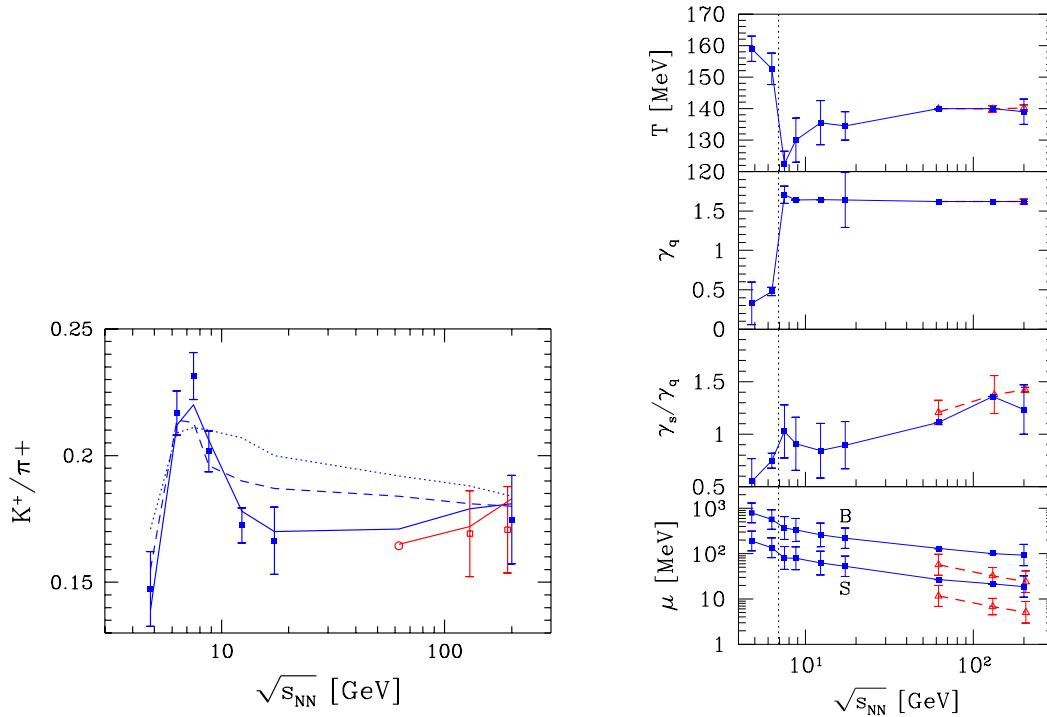
- Ξ/ϕ vs $v(Q)$, σ_{π^+/π^-} fits SHM well, **no** indication of ρ, σ modification
- Yields, ratios, resonances and fluctuations **fit reasonably well** provided $\gamma_q > 1$
- very unlikely equilibrium SHM can do it
- $\Lambda(1520)$, K^* well accounted for in non-equilibrium. $\sigma_{K/\pi}$ also fits well. **Consistent with sudden freeze-out!**

Do it yourself: SHARE

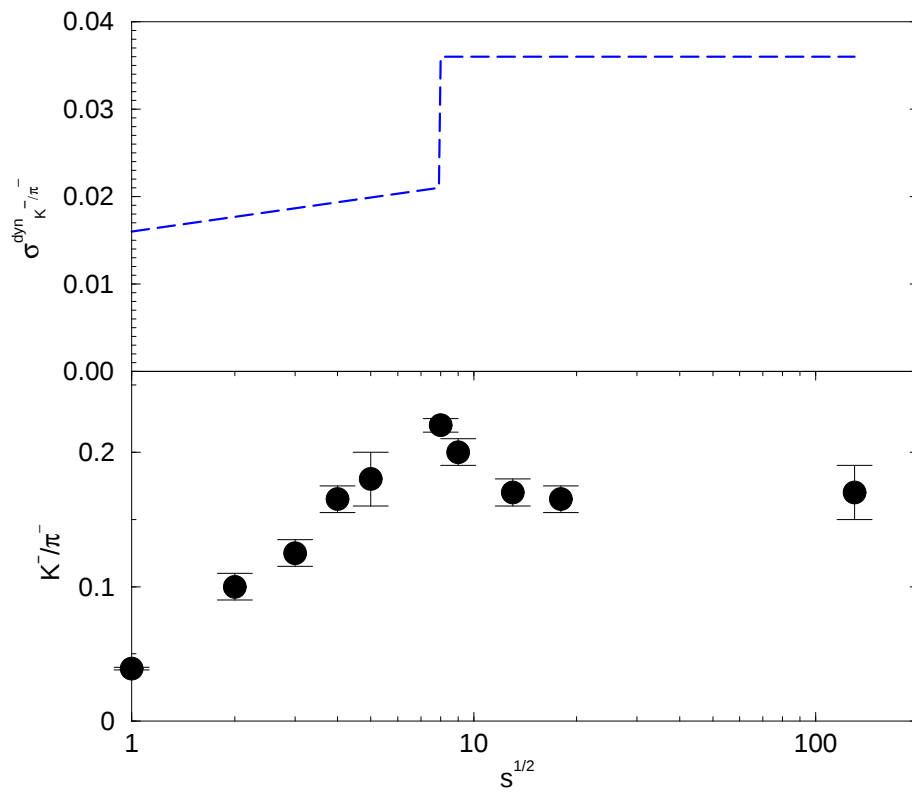
<http://www.physics.arizona.edu/~torrieri/SHARE/share.html>

A prediction...

It was recently shown ([nucl-th/0504028](#)) that the "horn" in the K^-/π^- yield



Can be explained by the onset of γ_q (sudden freeze-out from a Partonic plasma) at the "horn" energy.



Fluctuations in K^-/π^- should also rise sharply at the horn point.